

Lecture # 7

We have established for  
a cylinder, of period  
 $L = 2\pi R_0$ , where there is  
a "toroidal" field

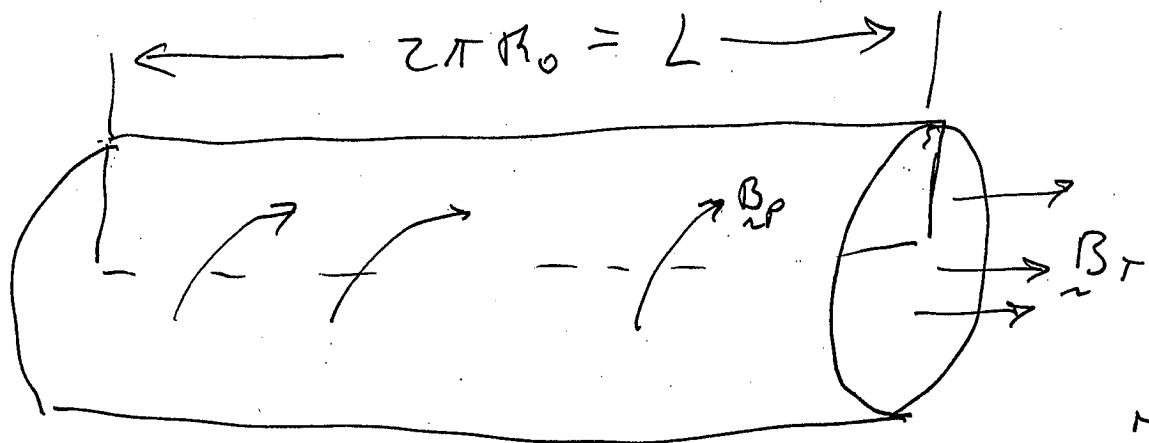
$$\begin{aligned} \underline{B}_z \hat{z} &= \underline{\nabla} \times A_\theta \hat{\theta} = \underline{\nabla} \times r A_\theta \underline{\nabla} \theta \\ &= \underline{\nabla} (A_\theta r) \times \underline{\nabla} \theta \end{aligned}$$

$$\begin{aligned} \underline{B}_\theta \hat{\theta} &= \underline{\nabla} \times A_z \hat{z} = \underline{\nabla} \times A_z R_0 \underline{\nabla} \phi \\ R_0 \underline{\nabla} \phi &= \underline{\nabla} z \end{aligned}$$

One easily shows

$$2\pi r A_\theta \equiv \text{Enclosed toroidal flux} \\ \text{within radius } r \equiv \Psi_T \equiv 2\pi \psi_T$$

$$\begin{aligned} 2\pi R_0 A_z &\equiv \text{Enclosed poloidal flux} \\ \text{within radius } r \\ &\equiv \Psi_P \equiv 2\pi \psi_P \end{aligned}$$



$$\Phi_p = \int B_\theta dA = L \int_0^R B_\theta dr = 2\pi R_0 \int_0^R B_\theta dr$$

$$\Phi_T = \int B_z dA = 2\pi \int_0^R dr r B_z$$

Magnetic field can be written as

$$\vec{B} = \nabla \psi_T \times \nabla \theta - \nabla \psi_p \times \nabla \phi$$

Also  $g(r) = \frac{d\psi_p(r)}{d\psi_T(r)}$

$$\therefore \nabla \psi_p = \frac{\partial \psi_p}{\partial \psi_T} \nabla \psi_T = g \nabla \psi_T$$

and we have

$$\vec{B} = \nabla \psi_p \times (\nabla (\theta g - \psi_T))$$

This is in fact a general form for  $\vec{B}$  with nested surfaces  $\psi$

# Generalized Coordinate System

Formed by intersection of three families of surfaces

$$q^1(x^1, x^2, x^3) = q^1(\underline{r})$$

$$q^2(\underline{r}), \quad q^3(\underline{r})$$

Then

$$d\underline{r} = \sum_i d q^i \frac{\partial \underline{r}}{\partial q^i} = d q^i \frac{\partial \underline{r}}{\partial q^i}$$

and

$$d q^i = \frac{\partial q^i}{\partial \underline{r}} \cdot d \underline{r} = \frac{\partial q^i}{\partial \underline{r}} \cdot \frac{\partial \underline{r}}{\partial q^j} d q^j$$

$$\text{And we have } \frac{\partial q^i}{\partial \underline{r}} \cdot \frac{\partial \underline{r}}{\partial q^j} = \delta^i_j = \begin{cases} 1 & \text{if } i=j \\ 0 & \text{if } i \neq j \end{cases}$$

Now a vector,  $\underline{A}$ , can be expressed in terms of covariant components  $\underline{A} = A_1 \underline{\nabla} q^1 + A_2 \underline{\nabla} q^2 + A_3 \underline{\nabla} q^3$

or contravariant components

$$\underline{A} = A^1 \frac{\partial \underline{r}}{\partial q^1} + A^2 \frac{\partial \underline{r}}{\partial q^2} + A^3 \frac{\partial \underline{r}}{\partial q^3}$$

Then we note

$$A^i = \nabla q^i \cdot \underline{A}$$

which follows from the  
relation

$$\frac{\partial q^i}{\partial r} \cdot \frac{\partial r}{\partial q^j} = \delta^i_j$$

Proof

$$\nabla q^i \cdot \left[ \frac{\partial r}{\partial q_1} A^1 + \frac{\partial r}{\partial q_2} A^2 + \frac{\partial r}{\partial q_3} A^3 \right]$$

$$= \nabla q^i \cdot \frac{\partial r}{\partial q^k} A^k = \delta^i_k A^k = A^i$$

Similarly Proof  $A_i = \frac{\partial r}{\partial q^i} \cdot \underline{A}$

$$A_i = \frac{\partial r}{\partial q^i} \cdot \underline{A} = \frac{\partial r}{\partial q^i} A_j \nabla q^j$$

$$\frac{\partial r}{\partial q^i} \cdot \left[ A_1 \nabla q^1 + A_2 \nabla q^2 + A_3 \nabla q^3 \right]$$

$$= \frac{\partial r}{\partial q^i} \cdot A_j \frac{\partial q^j}{\partial r} = A_j \delta^j_i = A_i$$

Further, we observe that

$$\frac{\partial \underline{r}}{\partial q^1} = \frac{\nabla q^2 \times \nabla q^3}{\nabla q^1 \cdot (\nabla q^2 \times \nabla q^3)}$$

as we then fulfill the conditions:

$$\frac{\partial \underline{r}}{\partial q^1} \cdot \nabla q^1 = 1$$

$$\frac{\partial \underline{r}}{\partial q^1} \cdot \nabla q^{2,3} = 0$$

Similarly:

$$\frac{\partial \underline{r}}{\partial q^2} = \frac{\nabla q^3 \times \nabla q^1}{\nabla q^1 \cdot (\nabla q^2 \times \nabla q^3)}, \quad \frac{\partial \underline{r}}{\partial q^3} = \frac{\nabla q^1 \times \nabla q^2}{\nabla q^1 \cdot (\nabla q^2 \times \nabla q^3)}$$

$$\frac{\partial \underline{r}}{\partial q^i} = \frac{\epsilon_{ijk} \nabla q^j \times \nabla q^k}{\nabla q^1 \cdot (\nabla q^2 \times \nabla q^3)}$$

Also recall that

$$\nabla q^1 \cdot \nabla q^2 \times \nabla q^3 = \det \left( \frac{\partial q^i}{\partial x^k} \right) \equiv J = \text{Jacobian}$$

$$dV = dx^1 dx^2 dx^3 = dq^1 dq^2 dq^3 J^{-1}$$

One way to see this must be the case is if we form a volume element

$$dV = d\vec{r}^1 \times d\vec{r}^2 \cdot d\vec{r}^3$$

$$d\vec{r}^1 = \frac{dq^1 \nabla q^2 \times \nabla q^3}{\nabla q^1 \cdot \nabla q^2 \times \nabla q^3} = \frac{dq^1 \nabla q^2 \times \nabla q^3}{J}$$

$$d\vec{r}^2 = dq^2 \nabla q^3 \times \nabla q^1 / J$$

$$d\vec{r}^3 = dq^3 (\nabla q^1 \times \nabla q^2) / J$$

Then

$$dV = d\vec{r}^1 \cdot d\vec{r}^2 \times d\vec{r}^3 =$$

$$= \frac{dq^1 dq^2 dq^3}{J^3} \left[ (\nabla q^2 \times \nabla q^3) \cdot [(\nabla q^3 \times \nabla q^1) \times (\nabla q^1 \times \nabla q^2)] \right]$$

$$[(\nabla q^2 \times \nabla q^3) \cdot \nabla q^1] (\nabla q^3 \times \nabla q^1) \cdot \nabla q^2$$

$$dV = \frac{dq^1 dq^2 dq^3}{J} = \frac{dq^1 dq^2 dq^3}{\det \left( \frac{\partial q^i}{\partial x^j} \right)}$$

$$J = \begin{vmatrix} \frac{\partial q^1}{\partial x^1} & \frac{\partial q^2}{\partial x^1} & \frac{\partial q^3}{\partial x^1} \\ \frac{\partial q^1}{\partial x^2} & \frac{\partial q^2}{\partial x^2} & \frac{\partial q^3}{\partial x^2} \\ \frac{\partial q^1}{\partial x^3} & \frac{\partial q^2}{\partial x^3} & \frac{\partial q^3}{\partial x^3} \end{vmatrix} = \frac{\partial q^1}{\partial x^1} \left( \frac{\partial q^2}{\partial x^2} \times \frac{\partial q^3}{\partial x^2} \right)$$

We also note that

$$\nabla_{\sim} \cdot \underset{\sim}{A} = \int \frac{\partial}{\partial q^i} \left( \frac{A_k}{f} \right)$$

because

$$\nabla_{\sim} = \frac{\partial q^1}{\partial q^1} + \frac{\partial q^2}{\partial q^2} + \dots$$

$$\nabla_{\sim} \cdot \left[ \frac{A_k}{f} \frac{\nabla q^2 \times \nabla q^3}{f} + \dots \right]$$

$$\text{Since } \nabla_{\sim} \cdot (\nabla q^2 \times \nabla q^3) = 0$$

$$\nabla q^2 \times \nabla q^3 = \frac{\partial q^1}{\partial q^1} \left( \frac{A_k}{f} \right) + \dots$$

$$= \int \frac{\partial}{\partial q^i} \left( \frac{A_k}{f} \right)$$

$$\nabla_{\sim} q^i \cdot \nabla_{\sim} \times \underset{\sim}{A} = \int \epsilon^{ijk} \frac{\partial}{\partial q^j} (A_k)$$

follows from

$$\nabla_{\sim} q^i \cdot \nabla_{\sim} \times [A_1 \nabla_{\sim} q^1 + A_2 \nabla_{\sim} q^2 + A_3 \nabla_{\sim} q^3]$$

$$\text{Using: } \nabla_{\sim} \times \nabla_{\sim} q^i = 0$$



# Metric Tensor

$$d\vec{r} = dq^i \frac{\partial \vec{r}}{\partial q^i}$$

$$d\vec{r} \cdot d\vec{r} = dq^i dq^j \frac{\partial \vec{r}}{\partial q^i} \cdot \frac{\partial \vec{r}}{\partial q^j} = g_{ij} dq^i dq^j$$

$$g_{ij} = \frac{\partial \vec{r}}{\partial q^i} \cdot \frac{\partial \vec{r}}{\partial q^j}$$

$$g \equiv \det g_{ij} = \det \left( \frac{\partial x^k}{\partial q^i} \quad \frac{\partial x^k}{\partial q^j} \right)$$

$$= \det \left( \frac{\partial x^k}{\partial q^i} \right) \det \left( \frac{\partial x^k}{\partial q^j} \right)$$

$$= \frac{1}{J} \cdot \frac{1}{J} = \frac{1}{J^2}$$

$$\frac{dq^1 dq^2 dq^3}{J} = \sqrt{g} dq^1 dq^2 dq^3 = dV$$