

The Scaling of Forced, Collisionless Reconnection in a Simple 3D System

Brian Sullivan, University of New Hampshire
Barrett Rogers, Dartmouth College
22 Nov 2008

Objectives

- Observe scaling of the reconnection process in a simple 3D system
- Compare observed scaling to that in an analogous 2D system

1st : 2D Scaling Study

- 2D 2-fluid model with finite e^- inertia

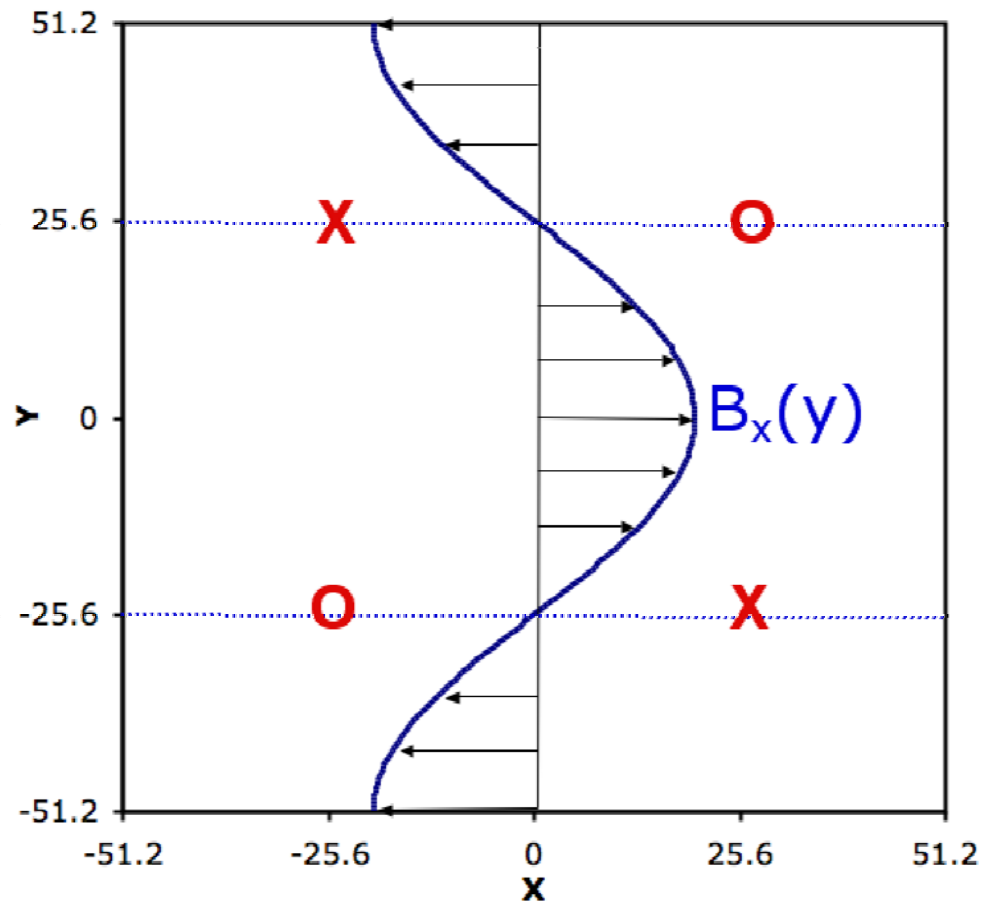
- $\frac{m_e}{m_i} = \frac{1}{25}$

- $L_x = L_y = 102.4 c / \omega_{pi}$

- Double Tearing mode equilibrium

- Periodic B.C.'s in both x and y

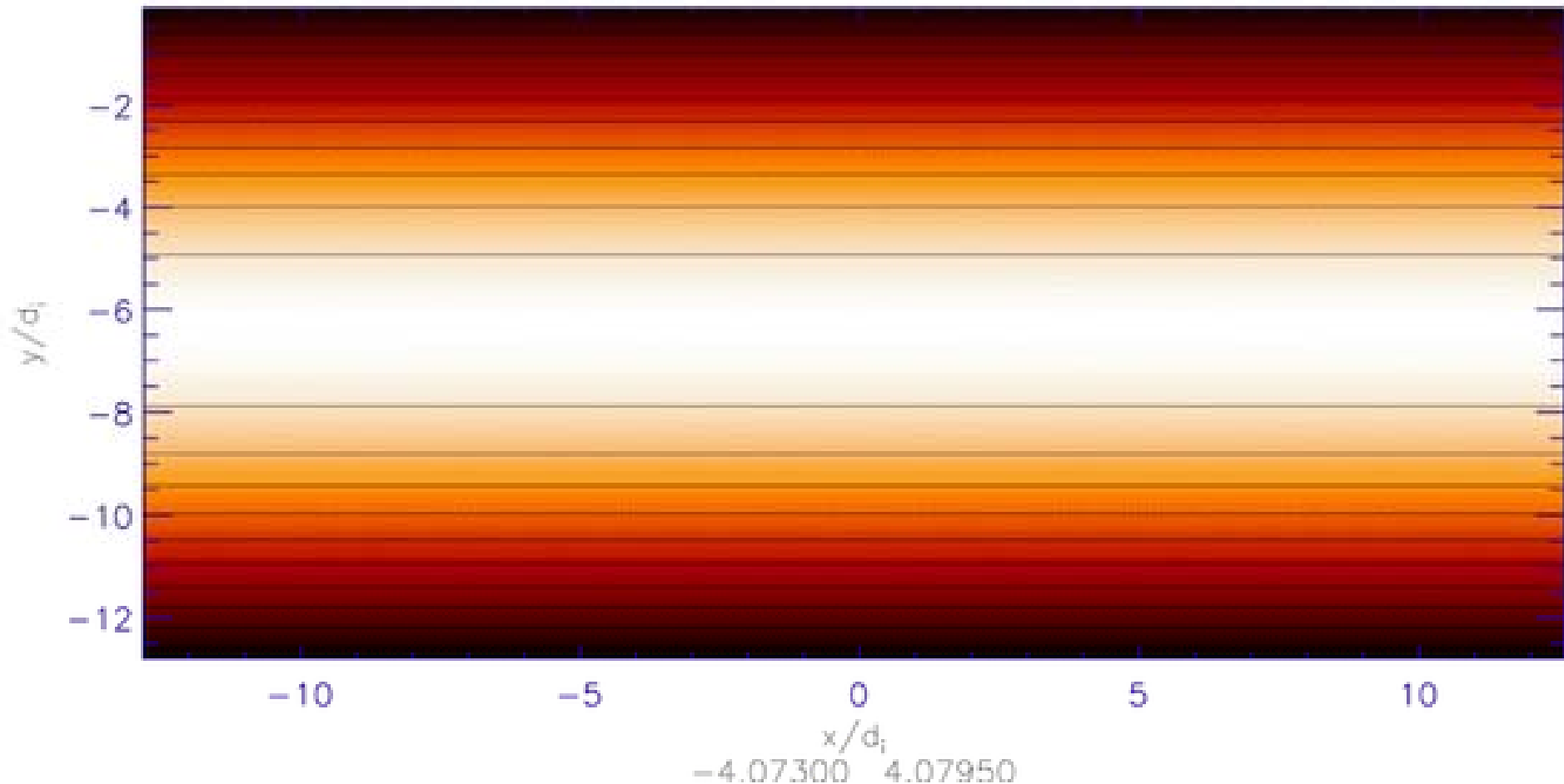
- Stable in absence of Forcing ($k_x \delta_J \geq 1$)



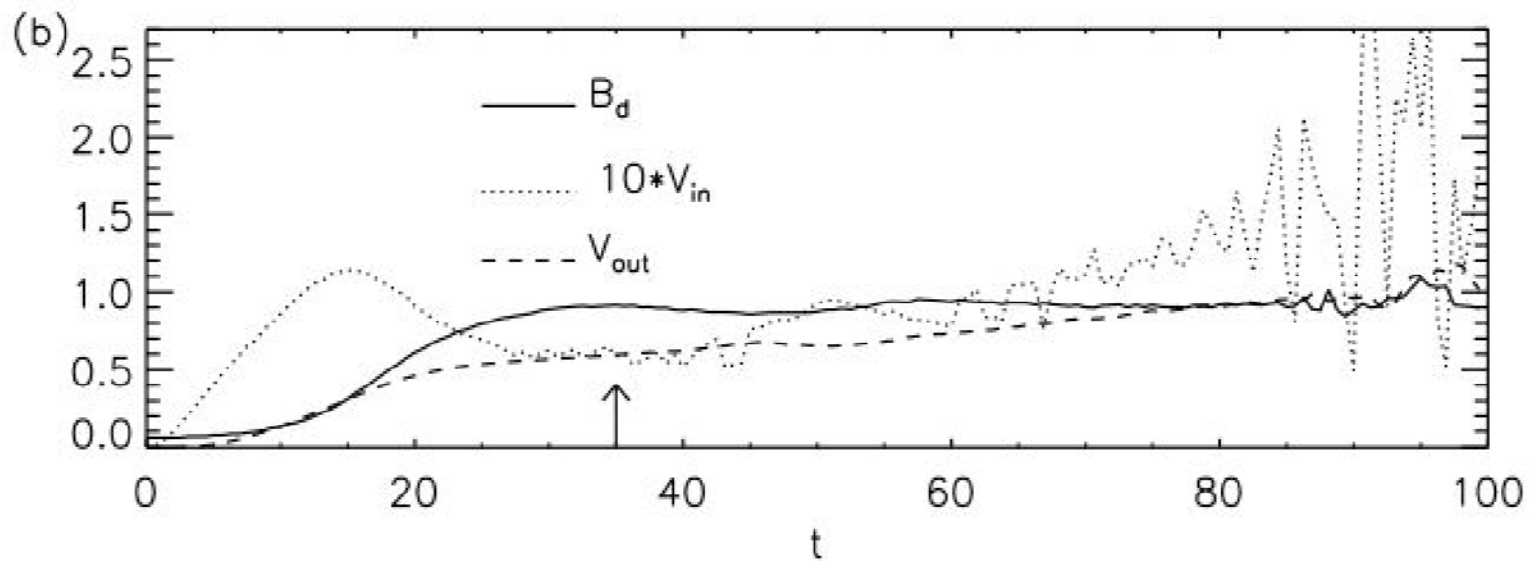
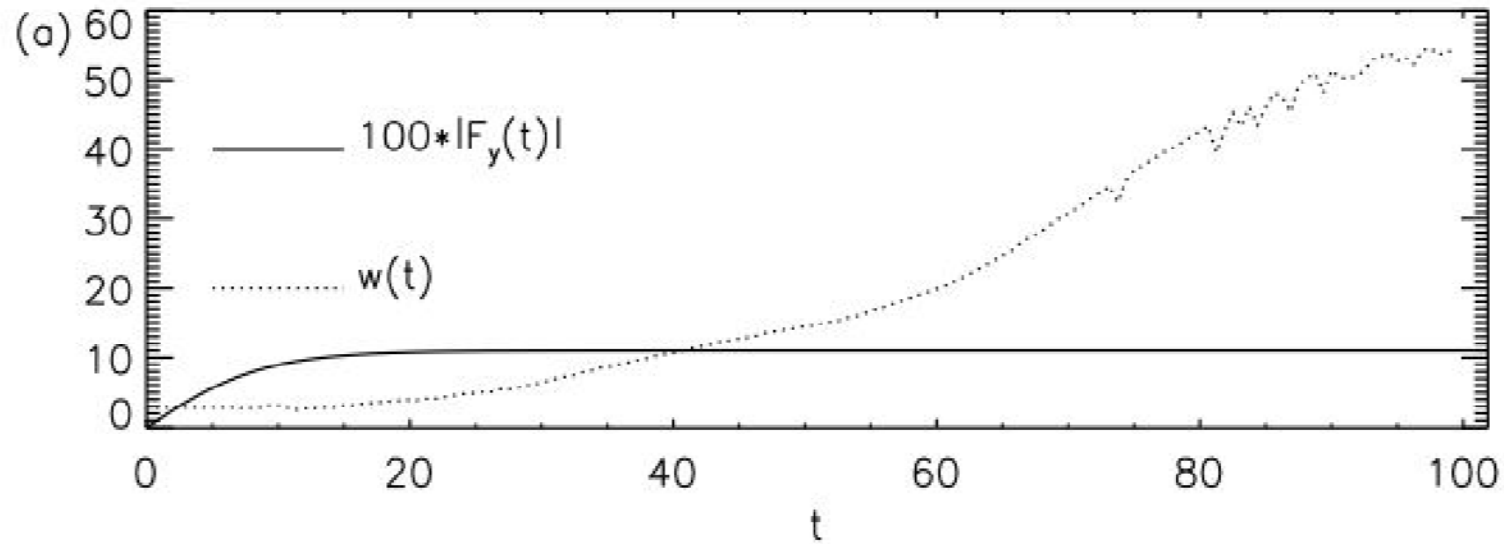
This system is stable, but if you force it...

$$\vec{J}_z(x, y)$$

t = 0.00000



Time series Data



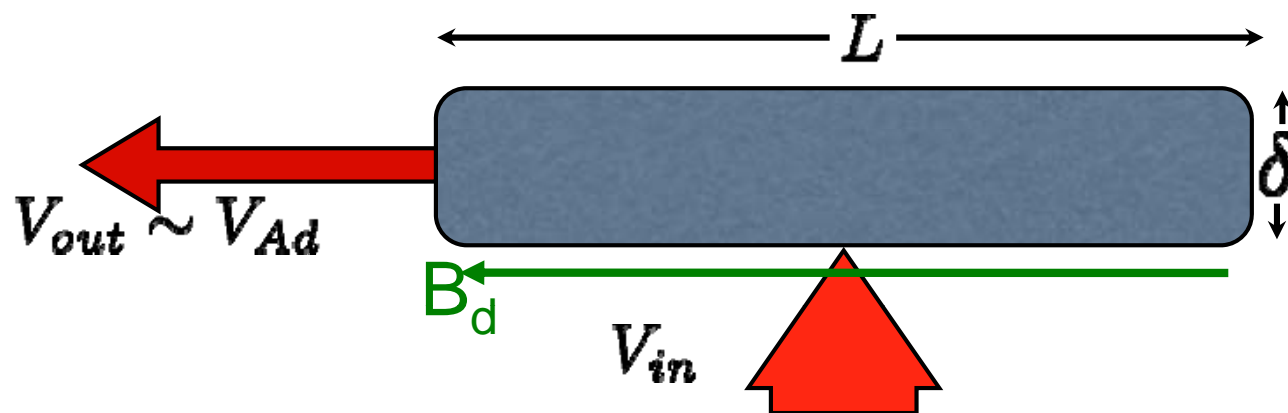
Expected R. Rate Scaling

$$R.Rate \sim E_z \sim V_{in} B_d$$

$$V_{in} \sim \left(\frac{\delta}{L} \right) V_{out}$$

$$V_{out} \sim V_{Ad} = \frac{B_d}{\sqrt{4\pi m_i n}}$$

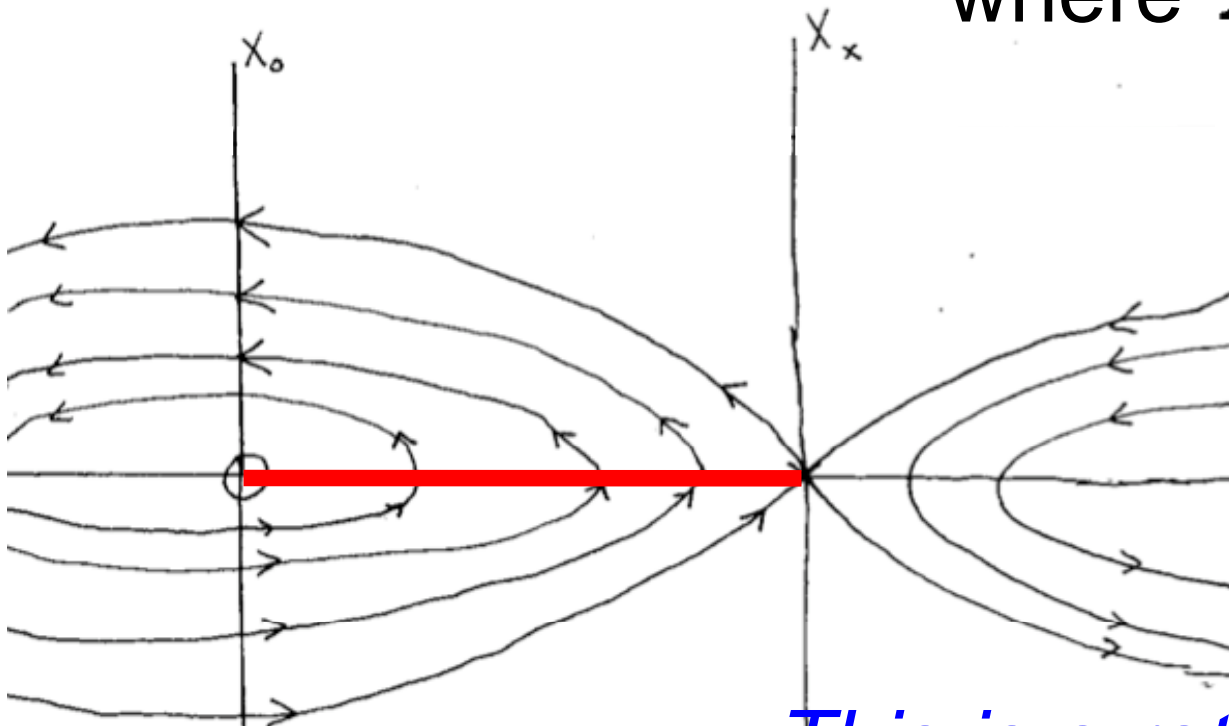
$$R.Rate \sim \frac{\delta}{L} \frac{B_d^2}{\sqrt{4\pi m_i n}}$$



Defining The Reconnection Rate

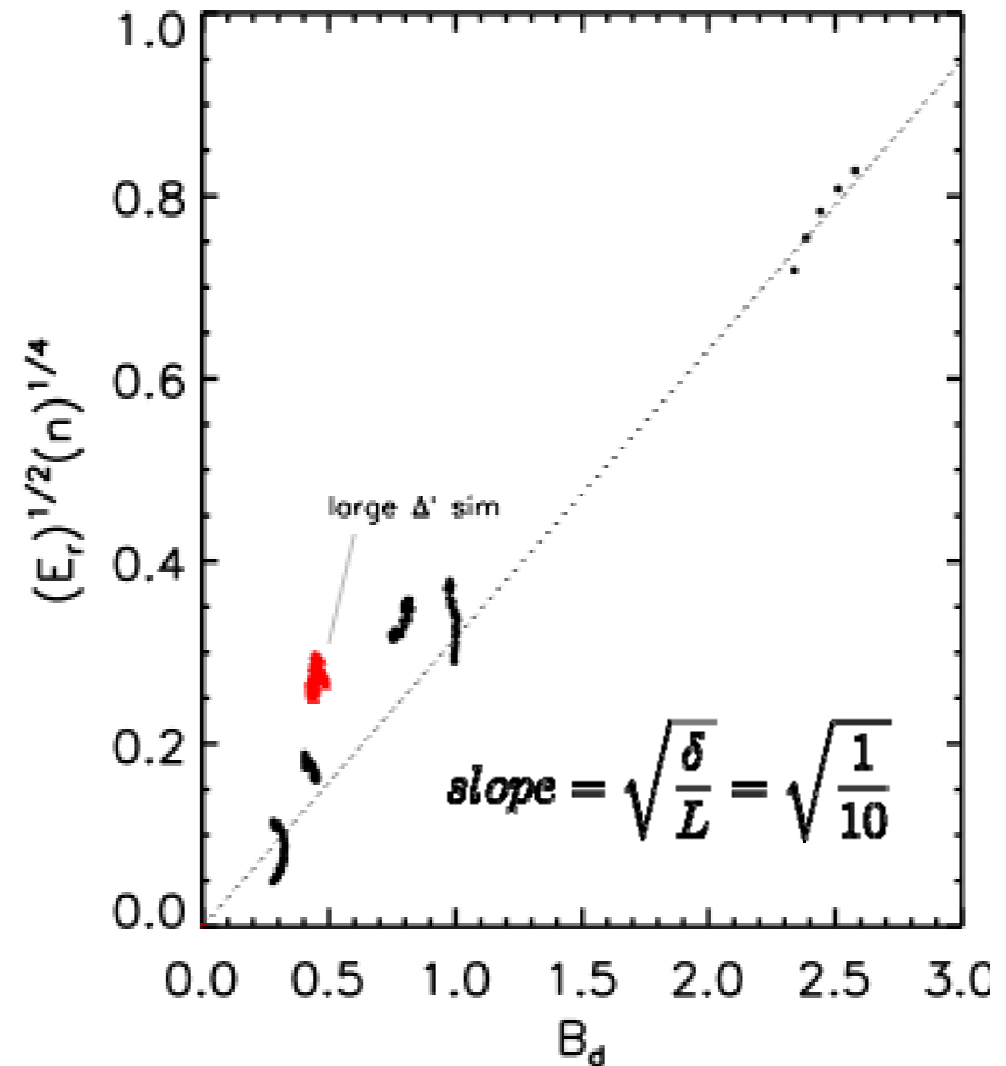
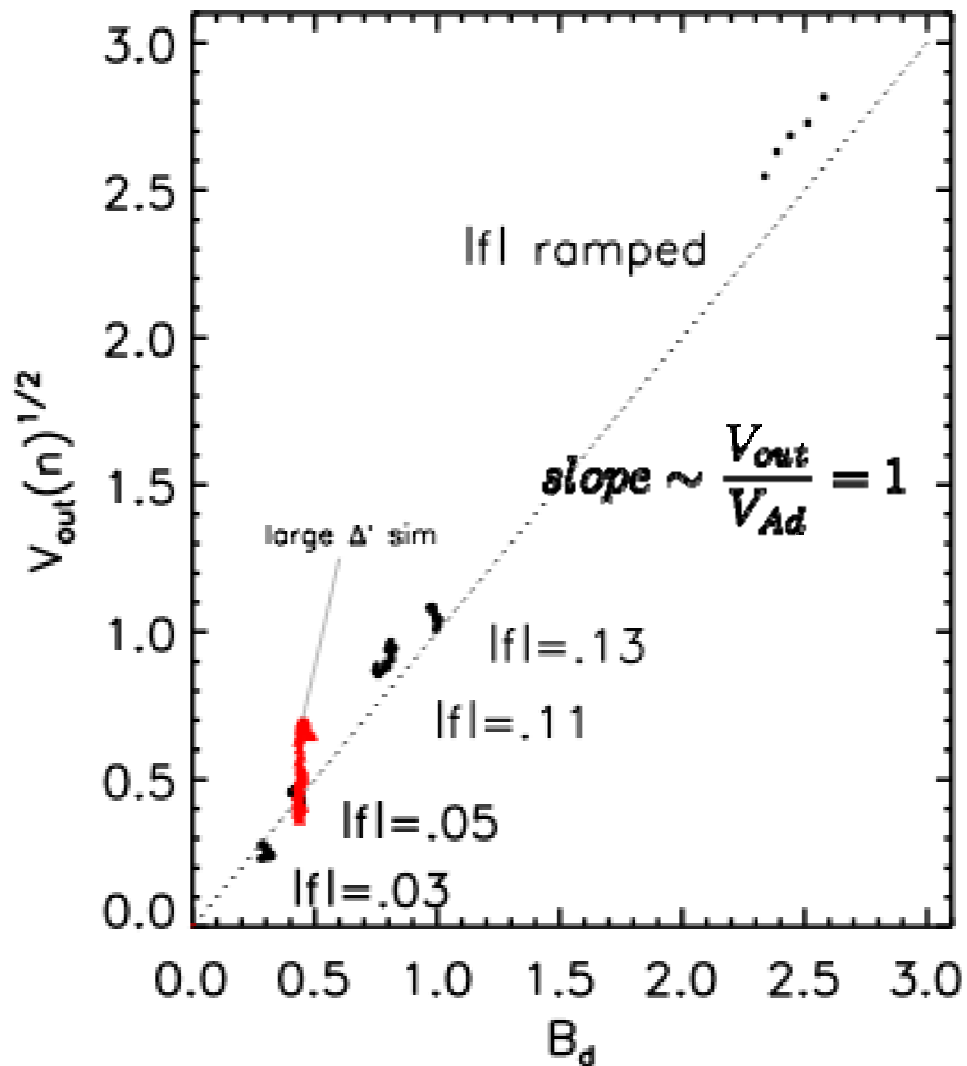
$$\begin{aligned} R_0 &\equiv \frac{\partial}{\partial t} [\psi(x_x, -L_y/4) - \psi(x_o, -L_y/4)] \\ &= \frac{\partial}{\partial t} \int_{x_o}^{x_x} B_y(x, -L_y/4) dx \end{aligned}$$

where $\vec{B} = \hat{z} \times \nabla\psi + B_z \hat{z}$
 $= (-\partial_y\psi, \partial_x\psi, B_z)$



This is a rate per unit length 6/24

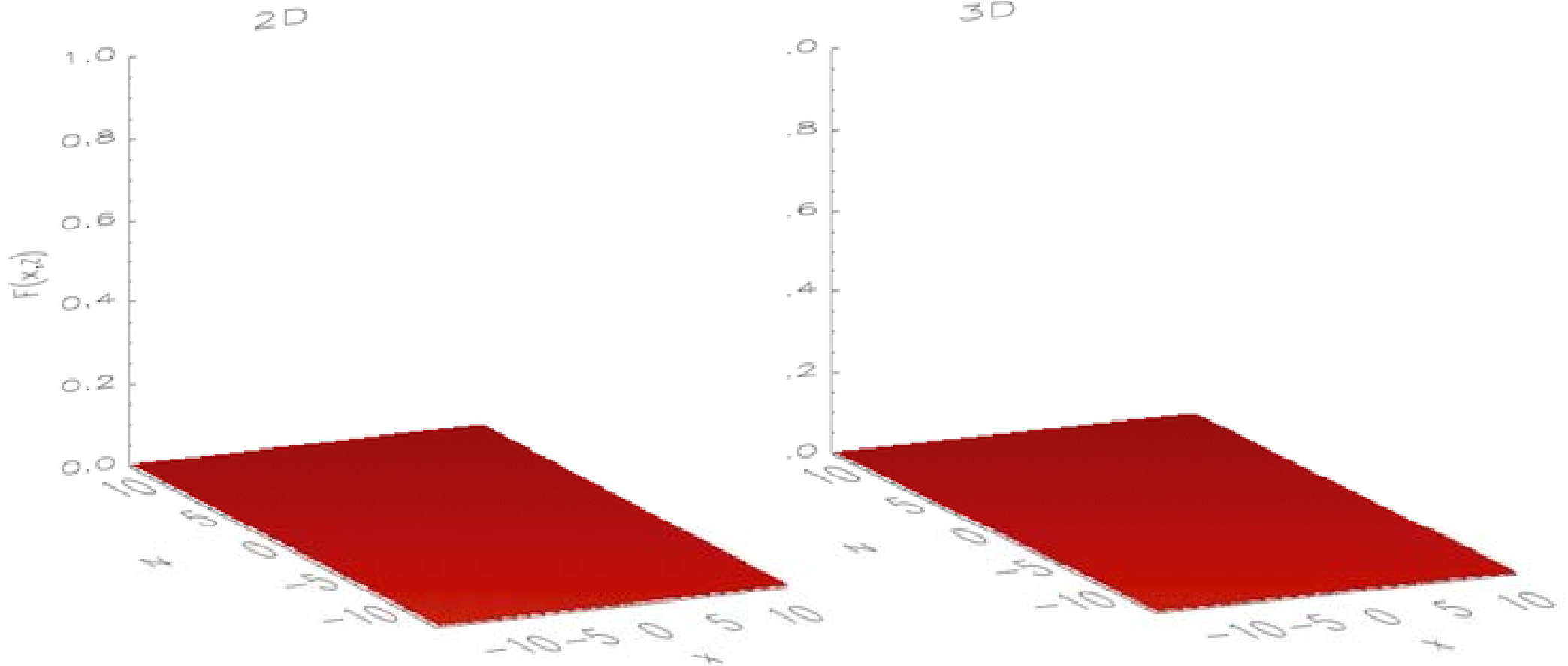
Forced 2D Data Agree with GEM Challenge Scaling



3D Study

- Same 1D initial equilibrium as in 2D study
- Smaller System size: $(25.6d_i \times 25.6d_i \times 25.6d_i)$
vs. $(102.4 d_i \times 102.4 d_i)$
- Forcing function scaled down with system
- Forcing localized in out-of-plane direction
- Comparisons made to a 2D analog

Out-Of-Plane Localization of Forcing



$$F_y(x, z)$$

Defining the 3D Reconnection Rate

In 3D we extend the 2D definition along z :

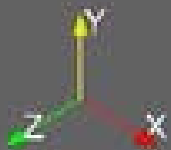
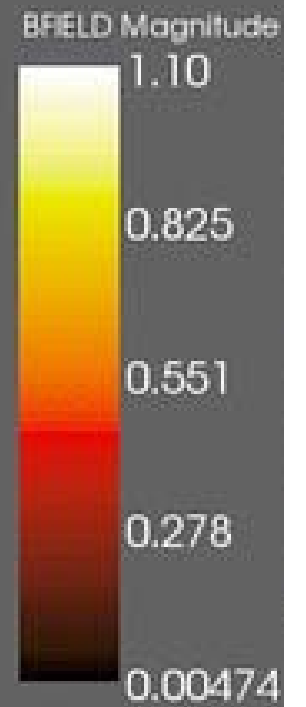
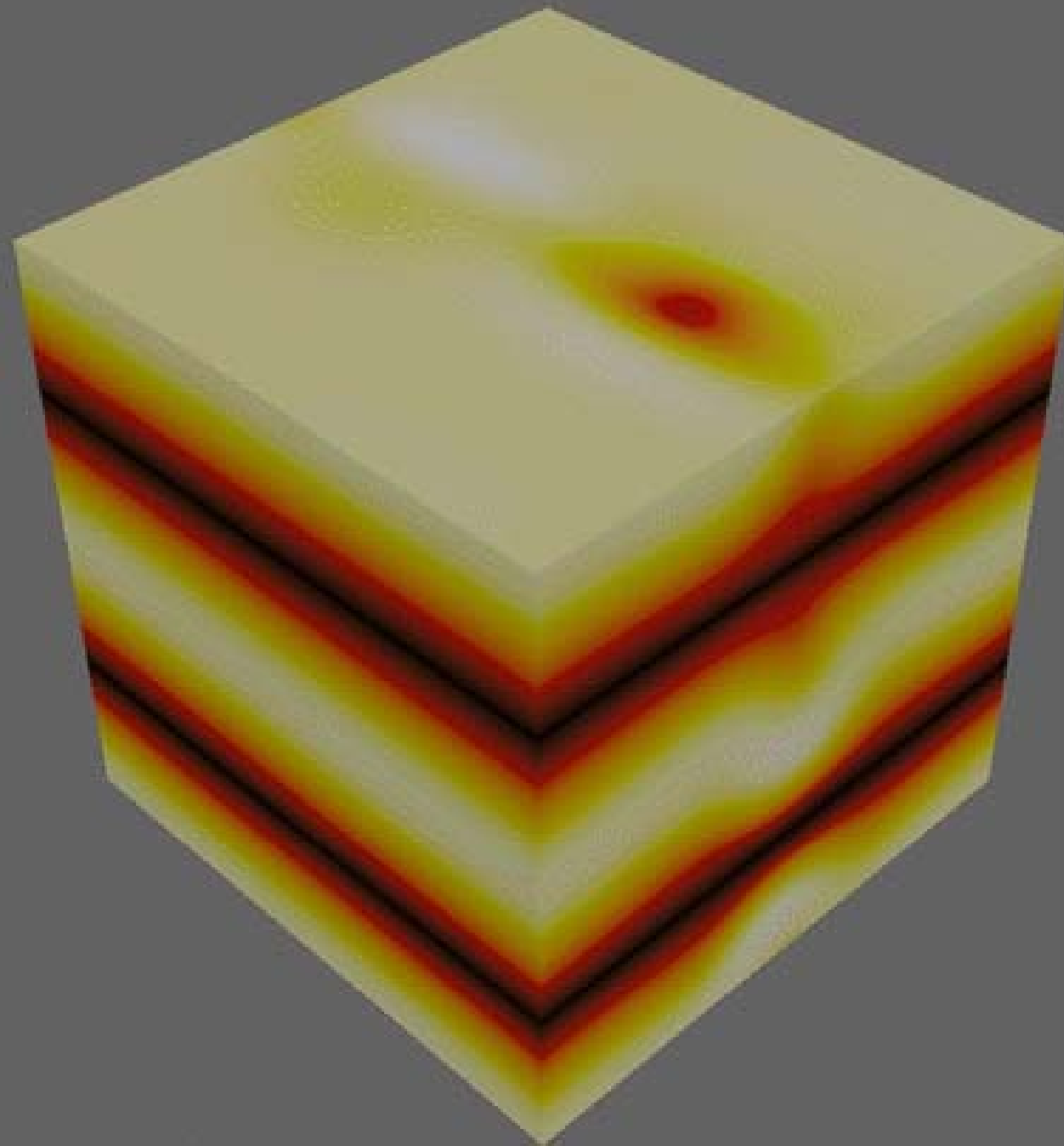
$$\Phi_{tot} = \int_A B_y(x, z) da$$

where A is the area of the magnetic island

$$R_{tot} = \frac{\partial}{\partial t} \Phi_{tot} = \frac{\partial}{\partial t} \int_0^{L_z} (\Phi_{2D}) dz$$

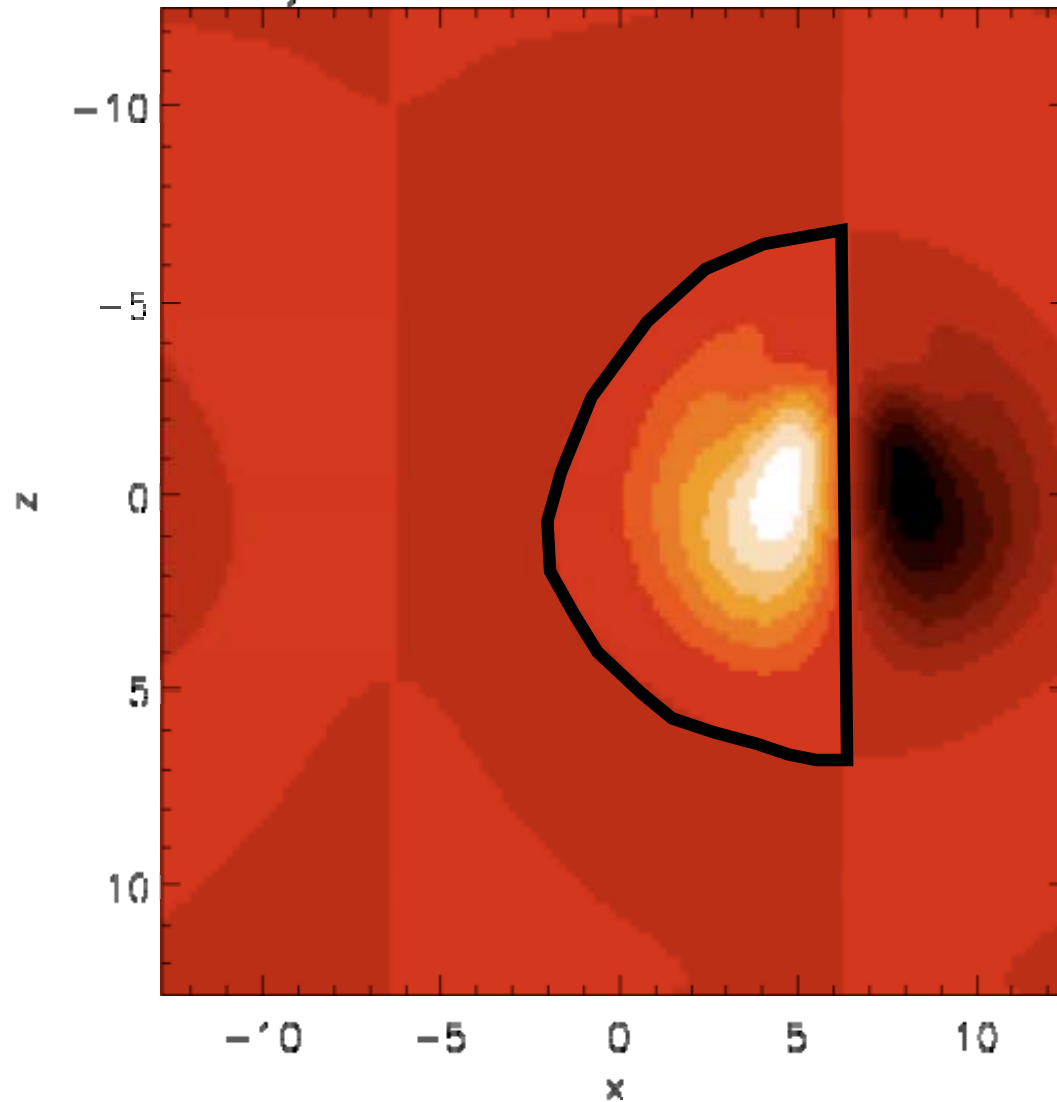
Careful: This is the total rate, not per unit length.

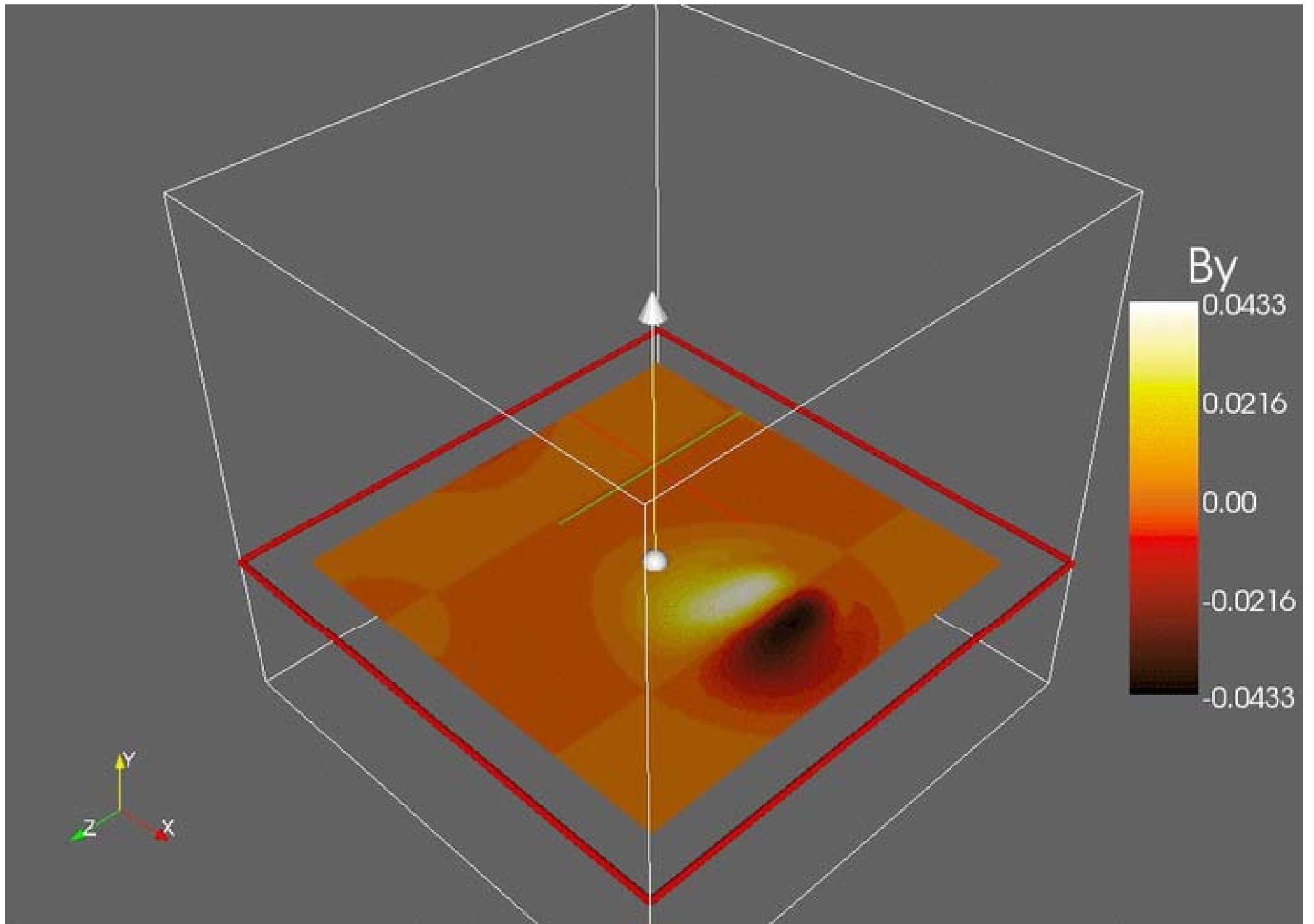
Where is the reconnected flux?

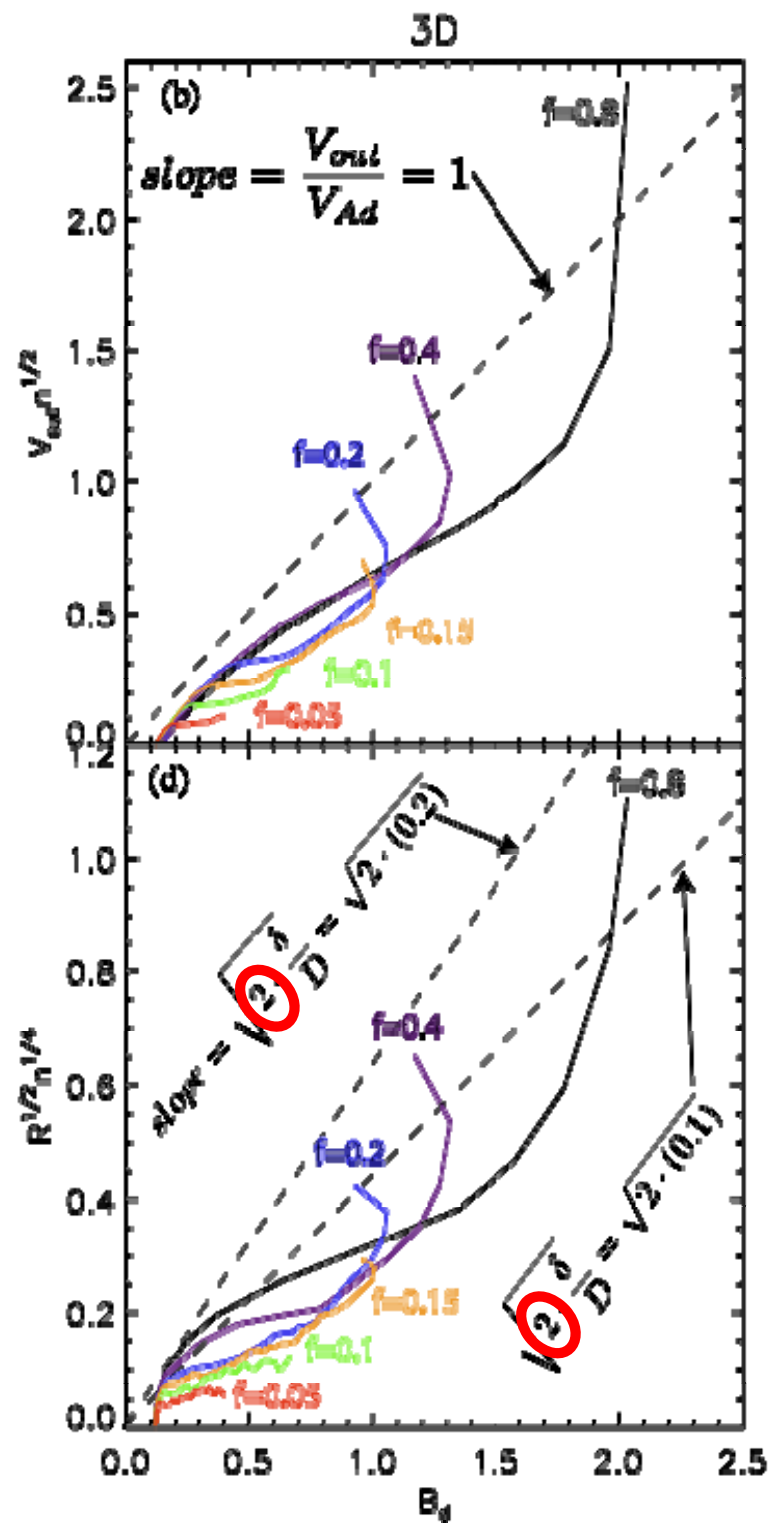
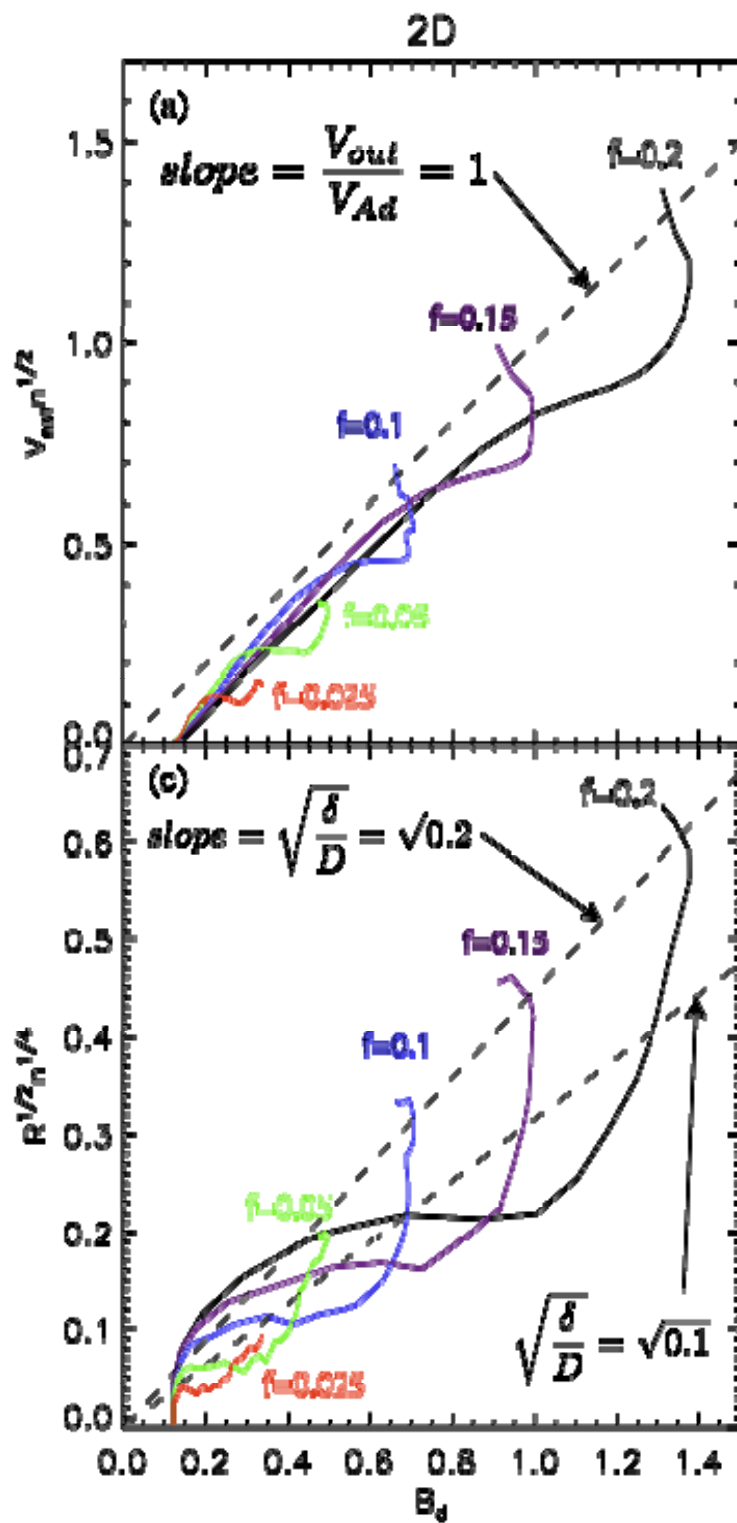


Shape of the Island

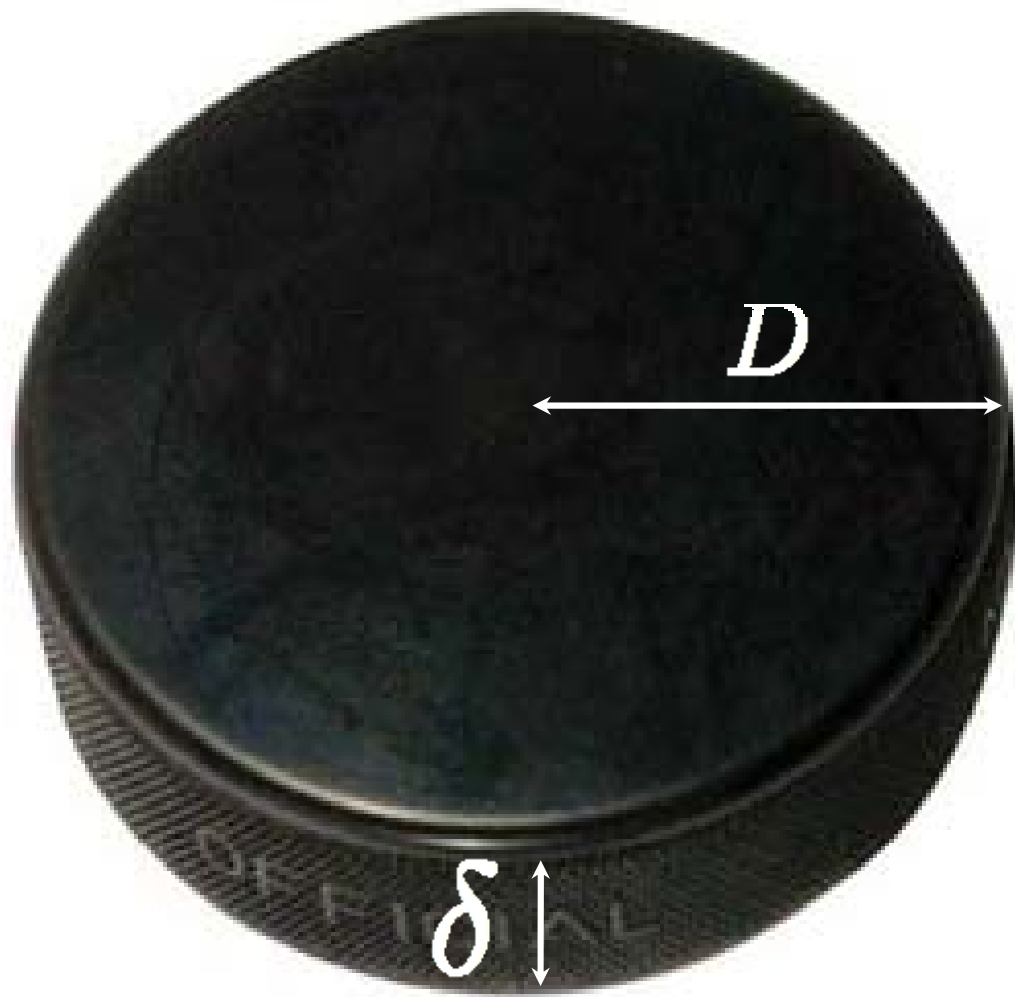
$B_y(x,z)$ at $y = L/4$, $t = 12.188$



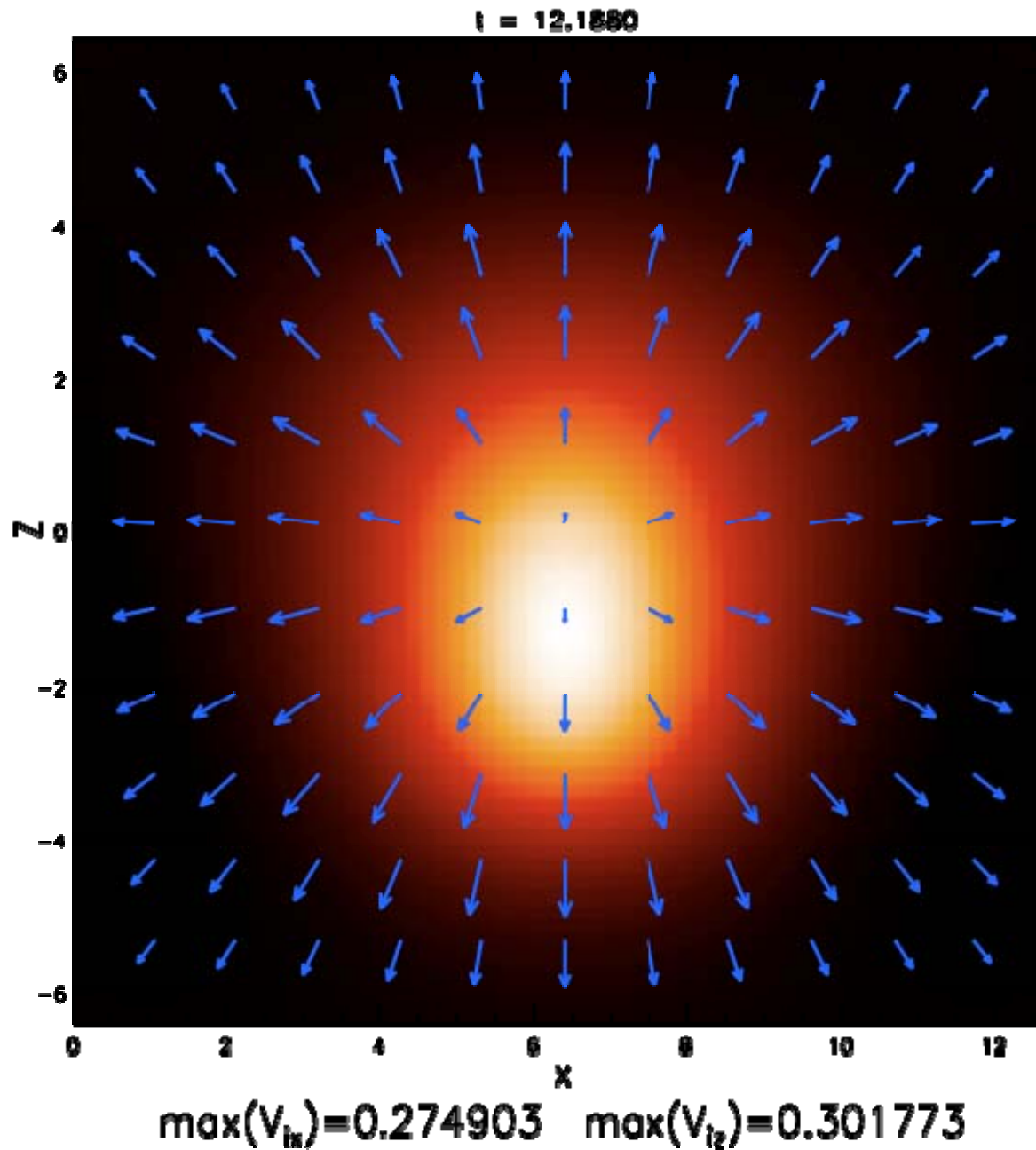




Why is inflow enhanced in 3D?



$$\frac{V_{in}}{V_{out}} \sim \frac{A_{out}}{A_{in}} = \frac{2\pi D\delta}{\pi D^2} = \frac{2\delta}{D}$$

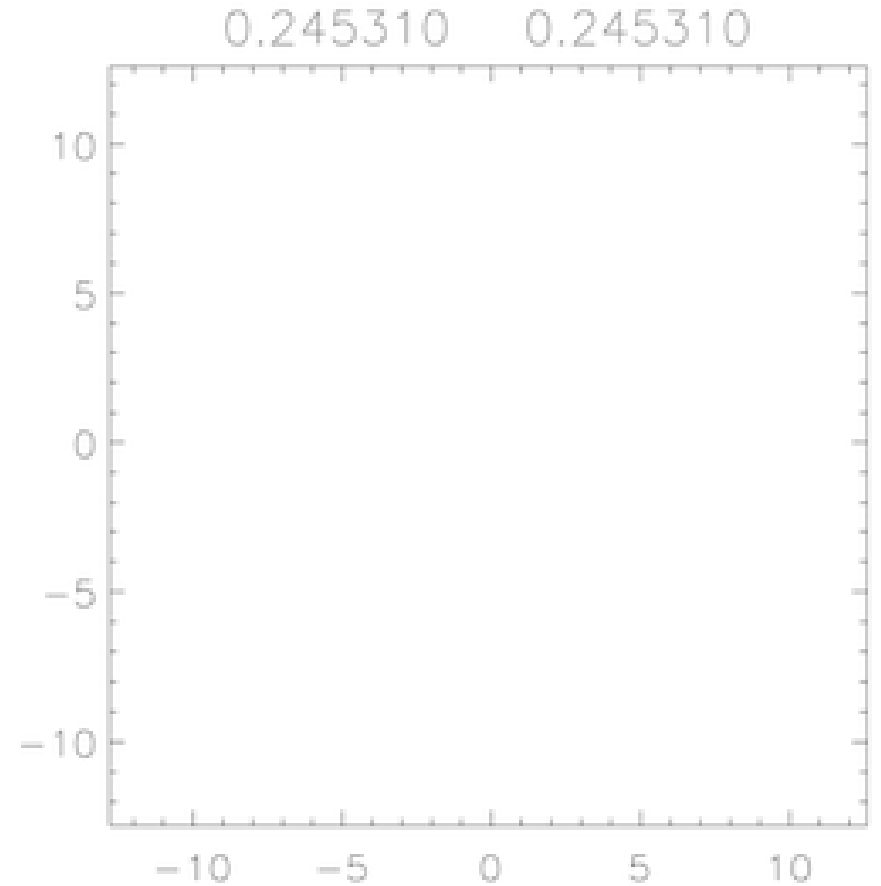
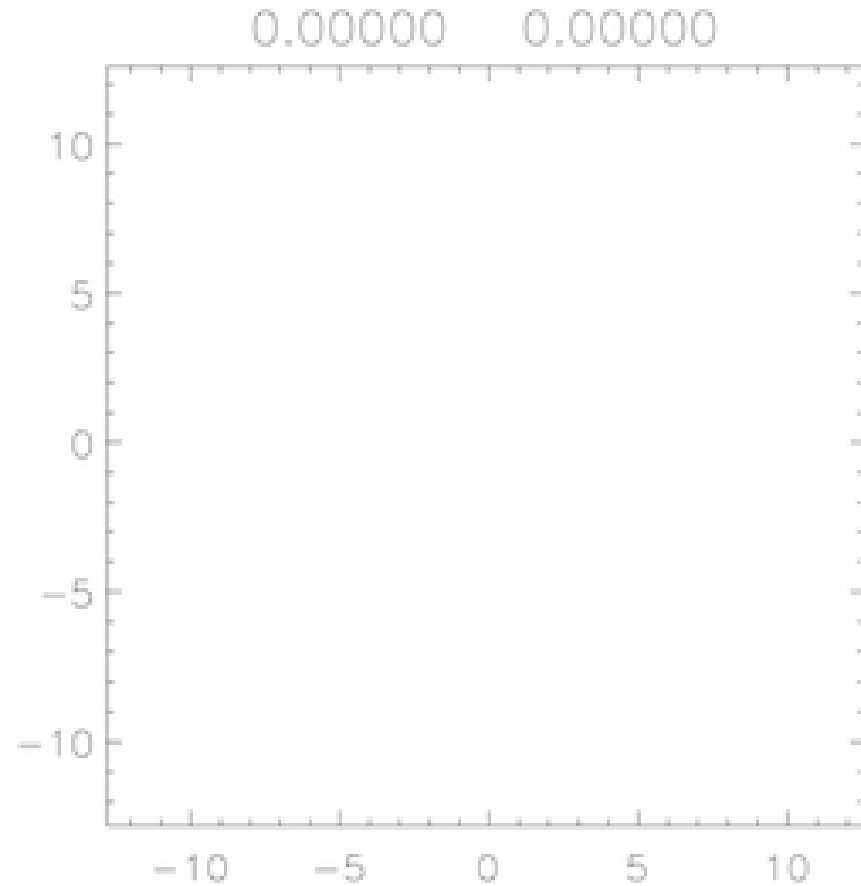


- Blue Arrows: $\vec{V}_i(x, z)$
- Greyscale : $J_z(x, z)$

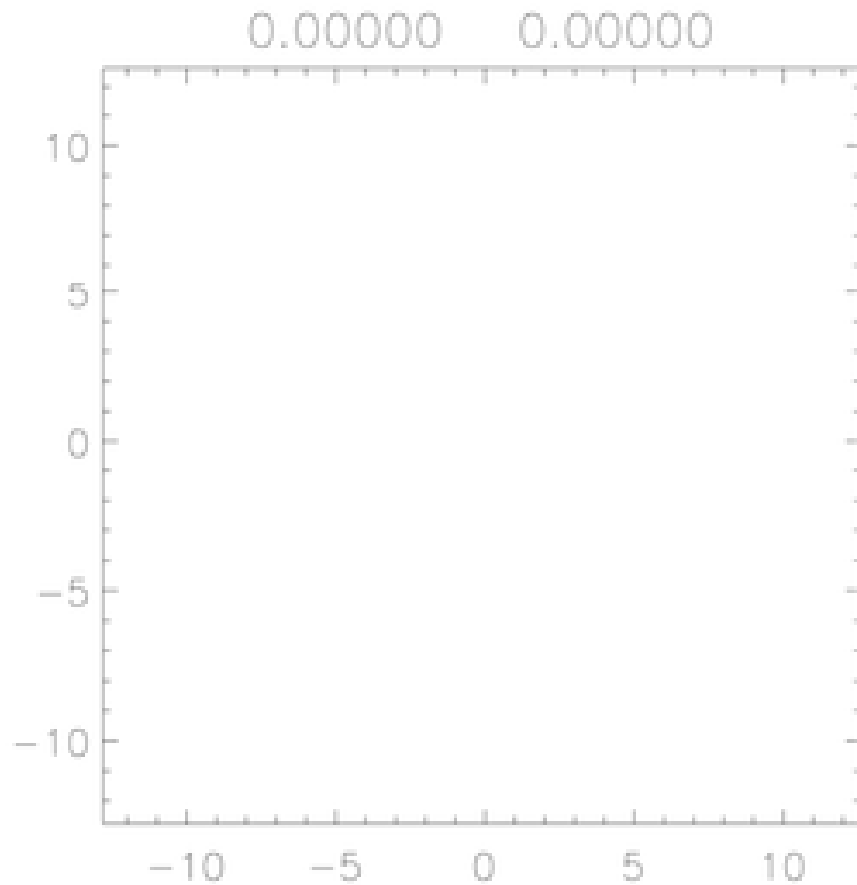
Forced Reconnection in
both 2D and 3D exhibit
very similar scaling

*However, there are some qualitative
differences in the dynamics...*

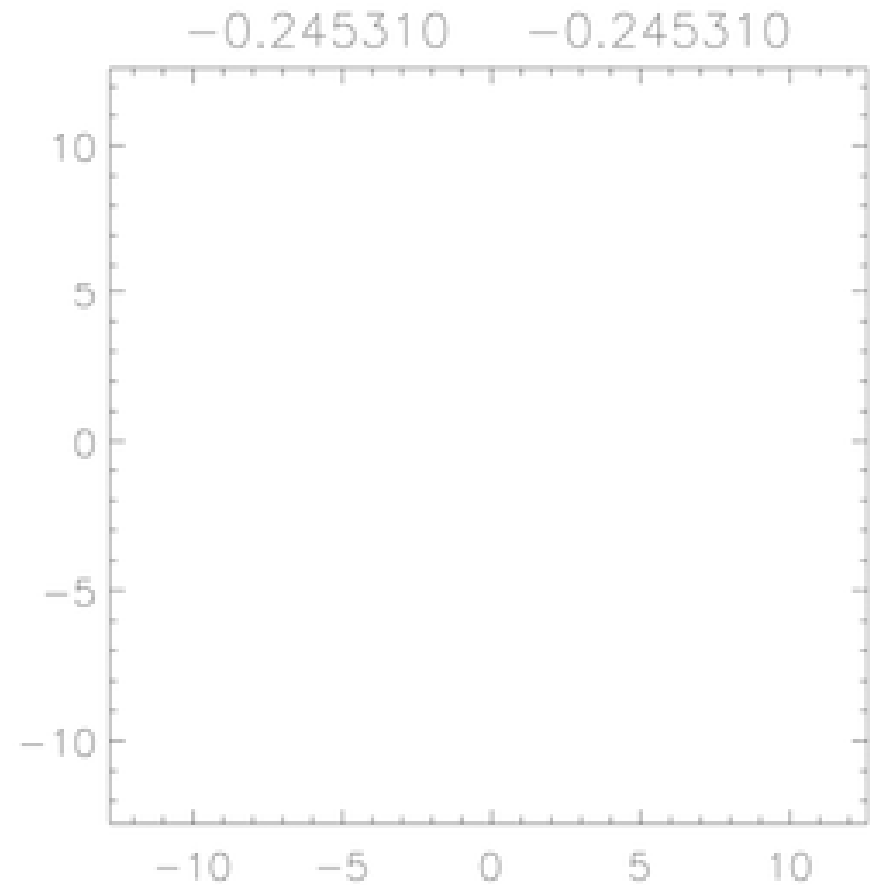
J_{iz} and J_z in 3D



The electrons carry the Magnetic Island in the $-z$ direction

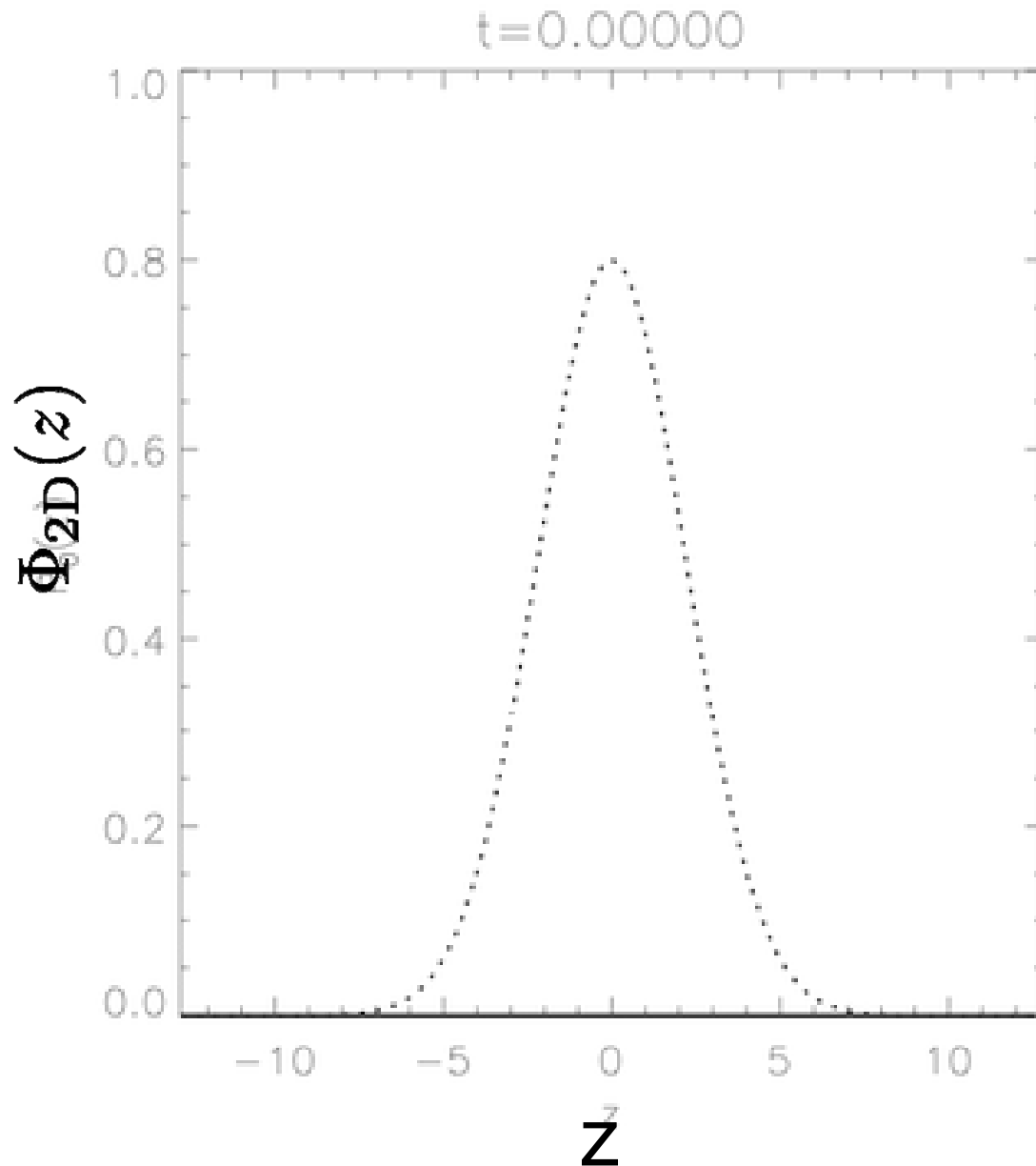


$B_y(x, z)$



$J_{ez}(x, z)$

Movement of Reconnected Flux along z

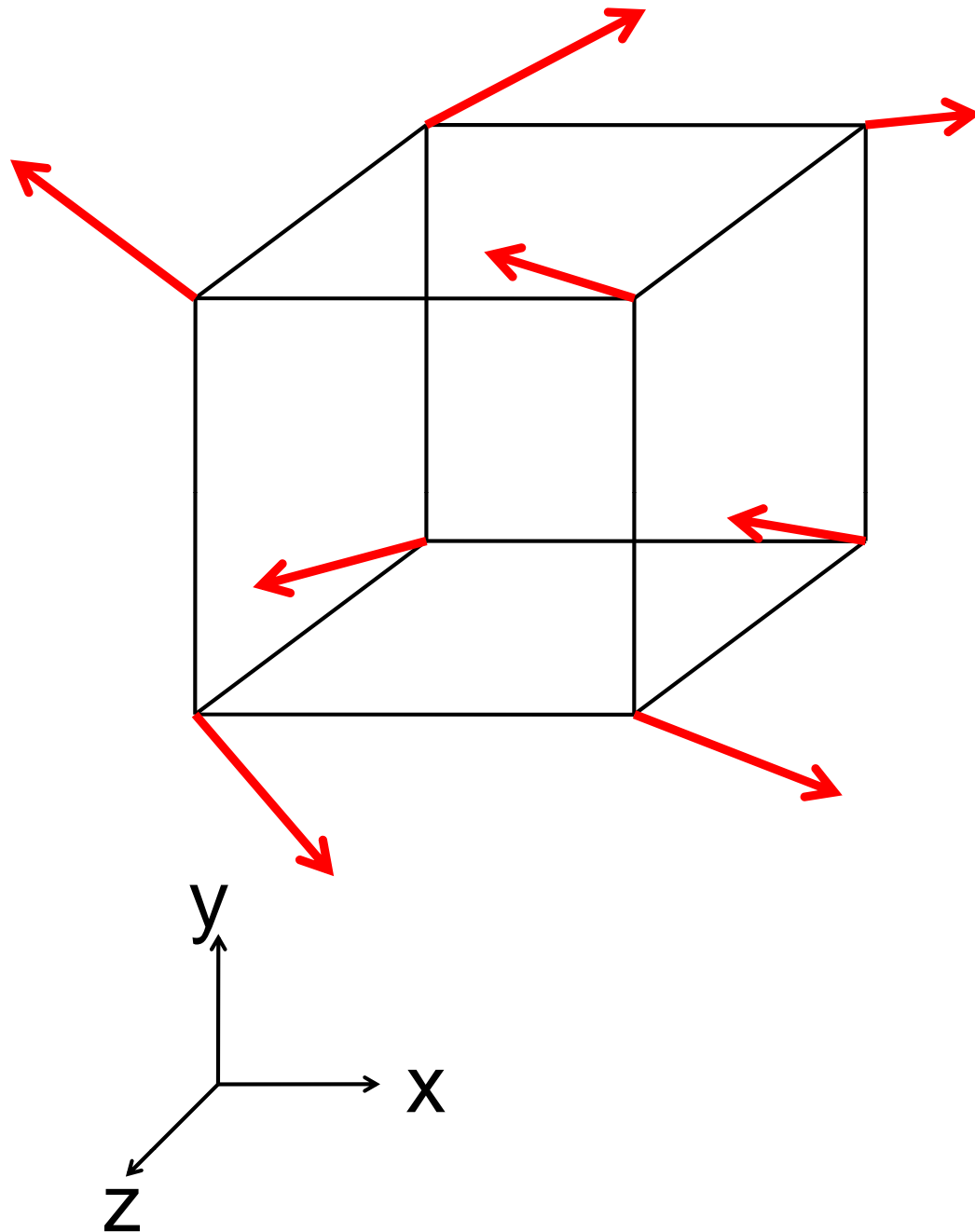


Dotted curve: Forcing $F(z)$

Solid curve: Flux $\Phi_{2D}(z)$

Where are the nulls of B?

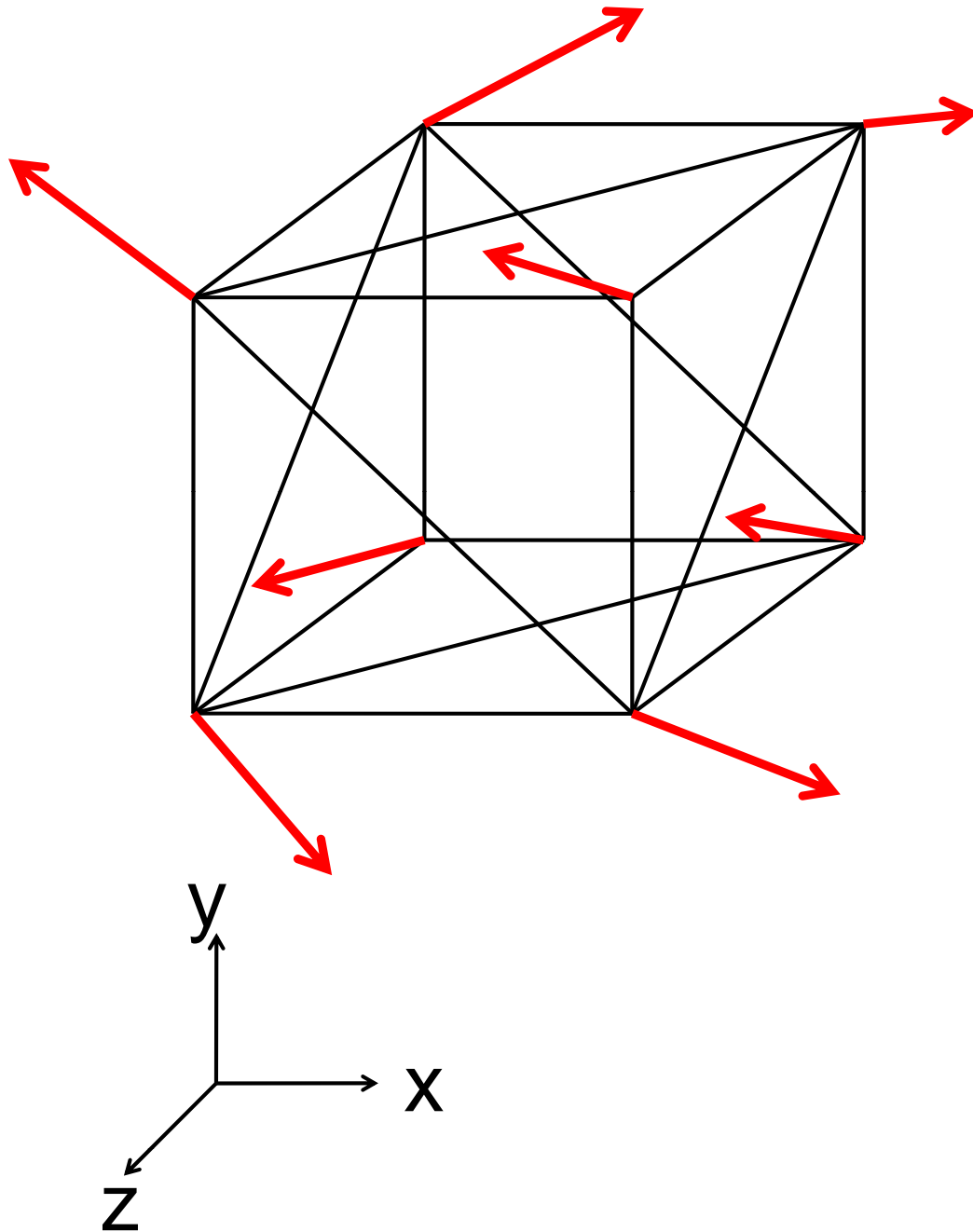
- Nulls are located using an Algorithm by John Greene (*J. Comp. Phys.* 1992)
- Algorithm implemented by John Dorelli
- Here's how it works...

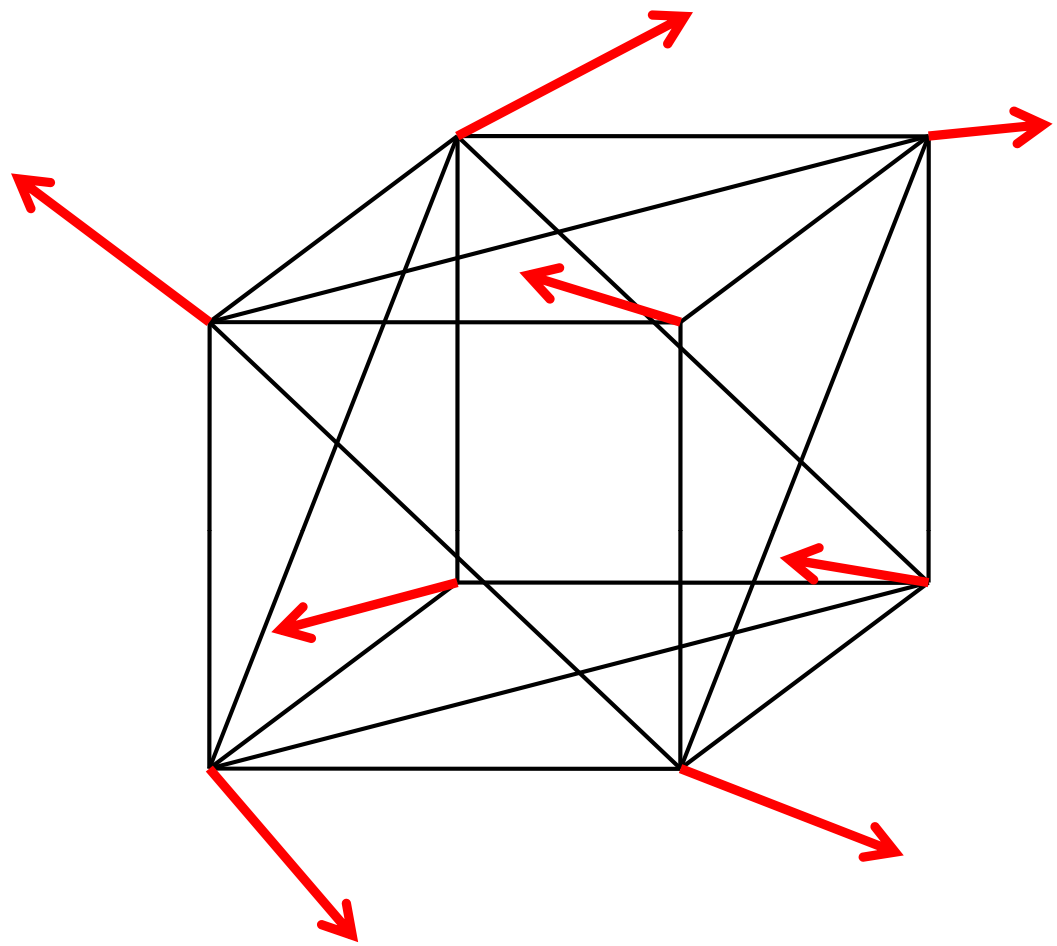


Magnetic field
vectors at
vertices of one
grid cell

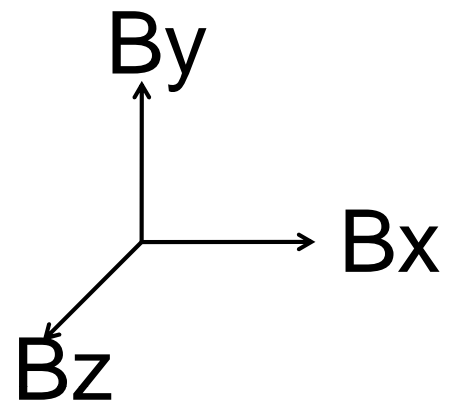
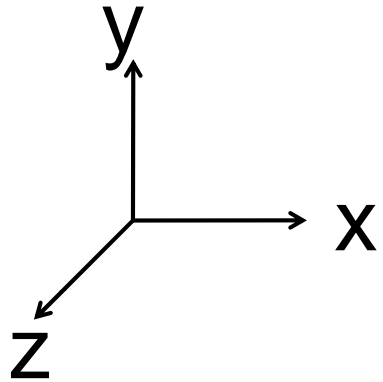
Does this cell contain
a null of the field?

Divide cube faces
into triangular
simplexes

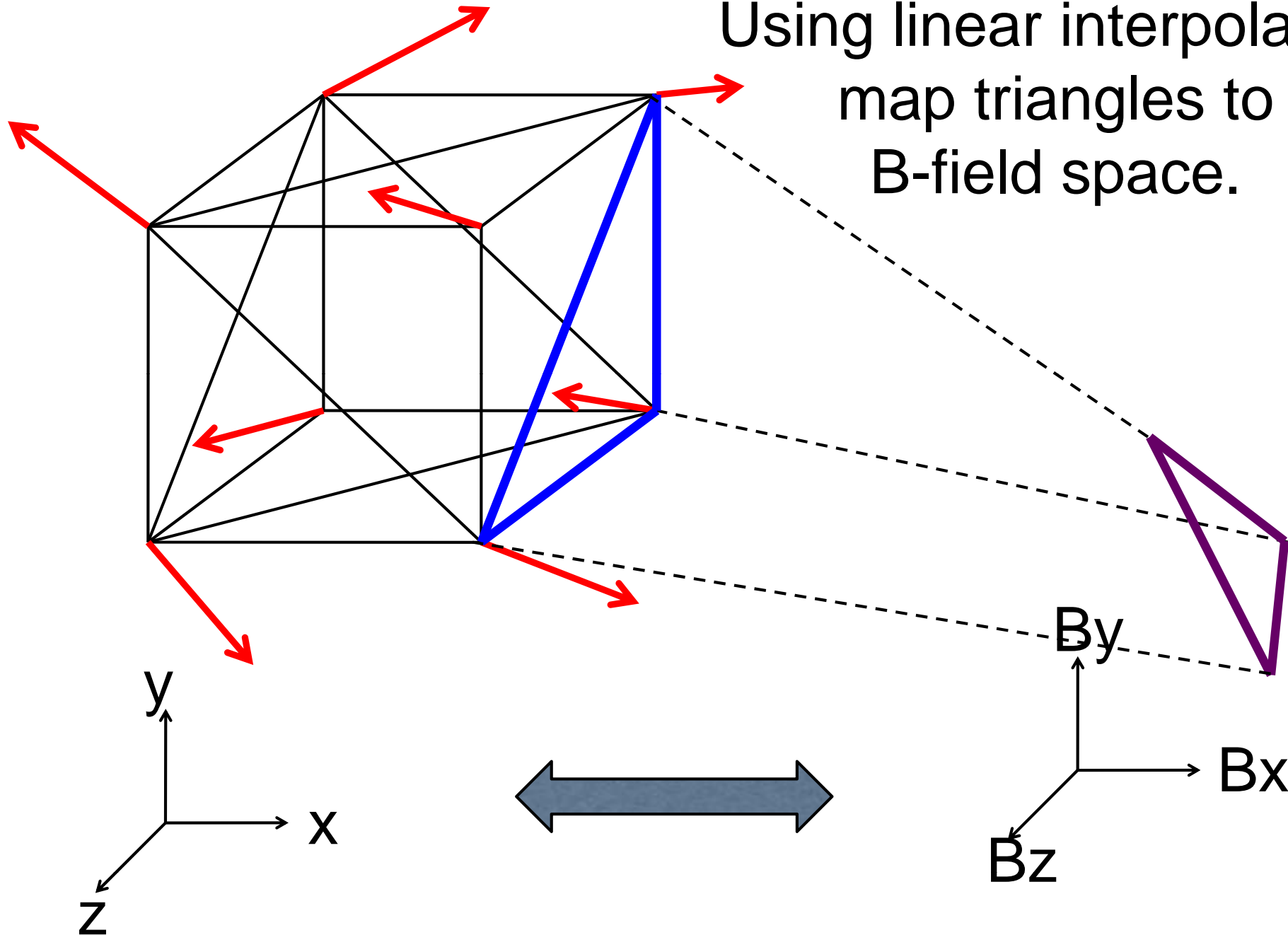


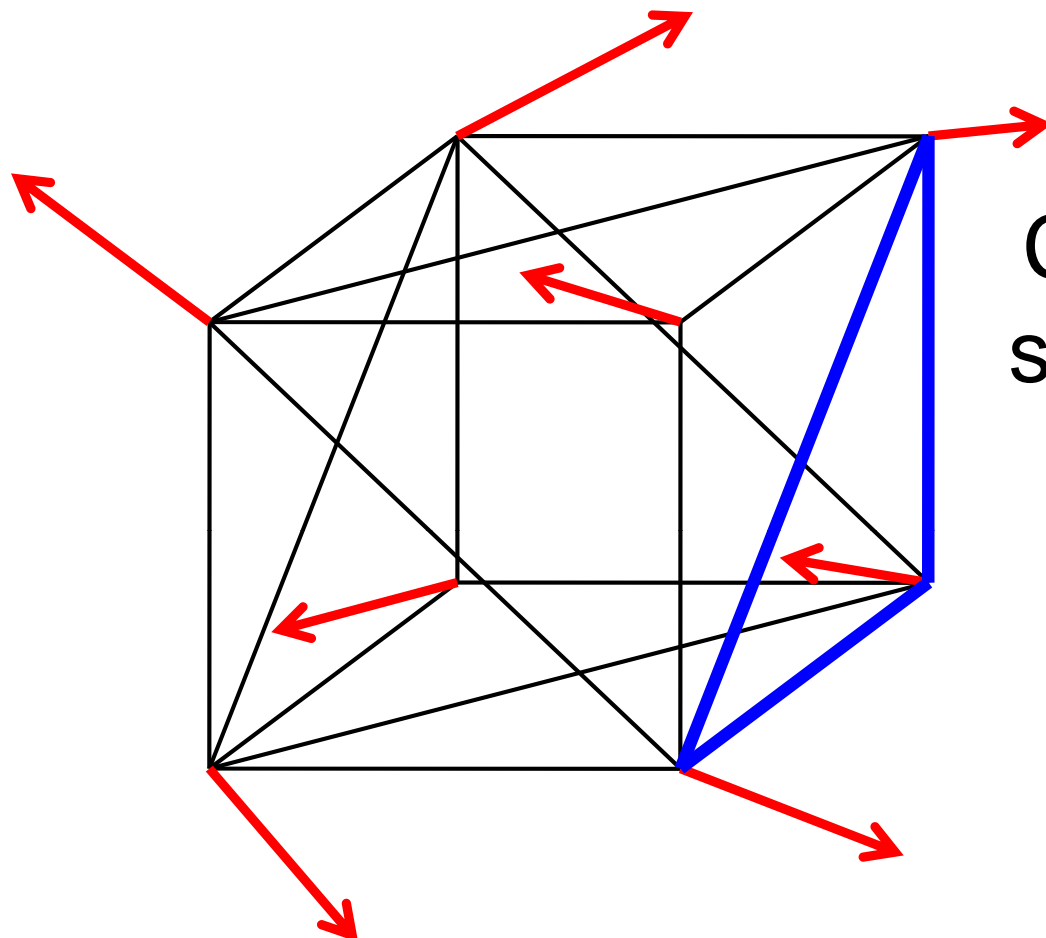


Suppose a mapping
from configuration
space to a B-field
space:

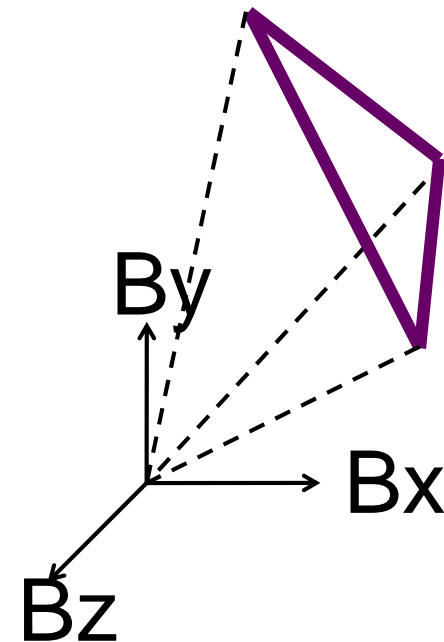
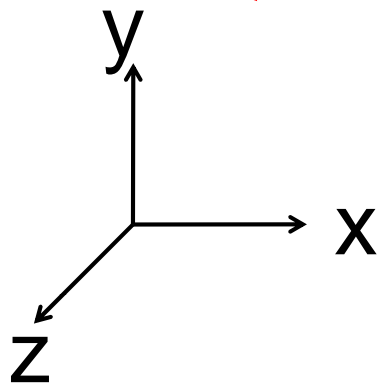


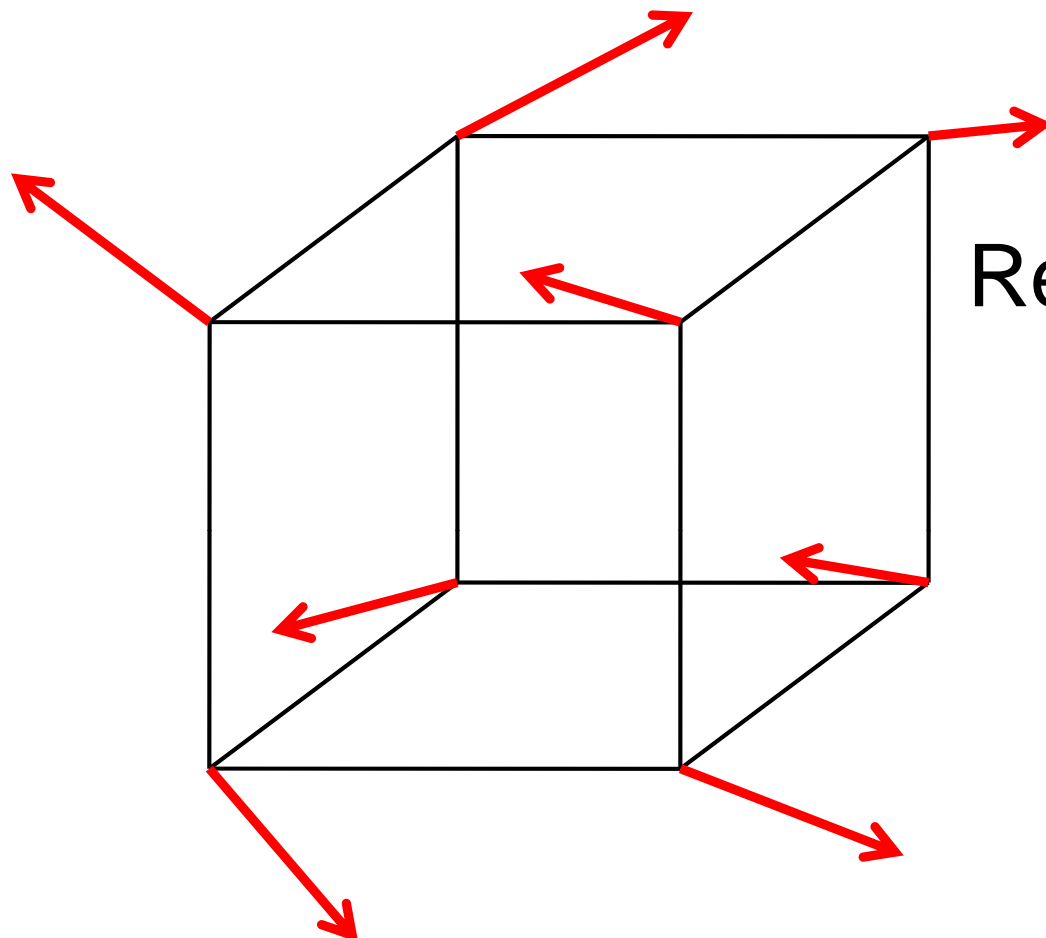
Using linear interpolation
map triangles to
B-field space.



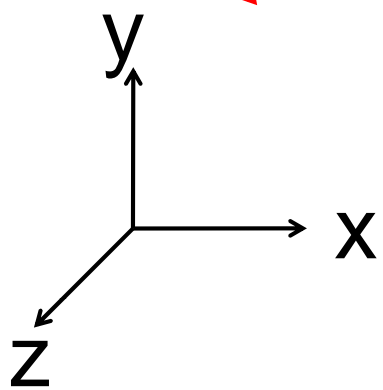


Calculate solid angle subtended by triangle at origin of B-space

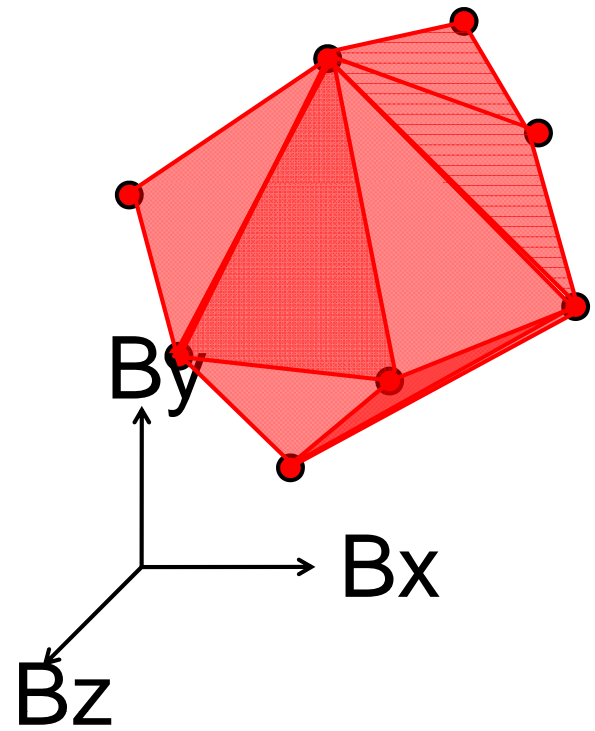




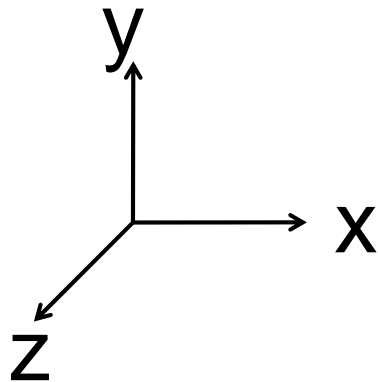
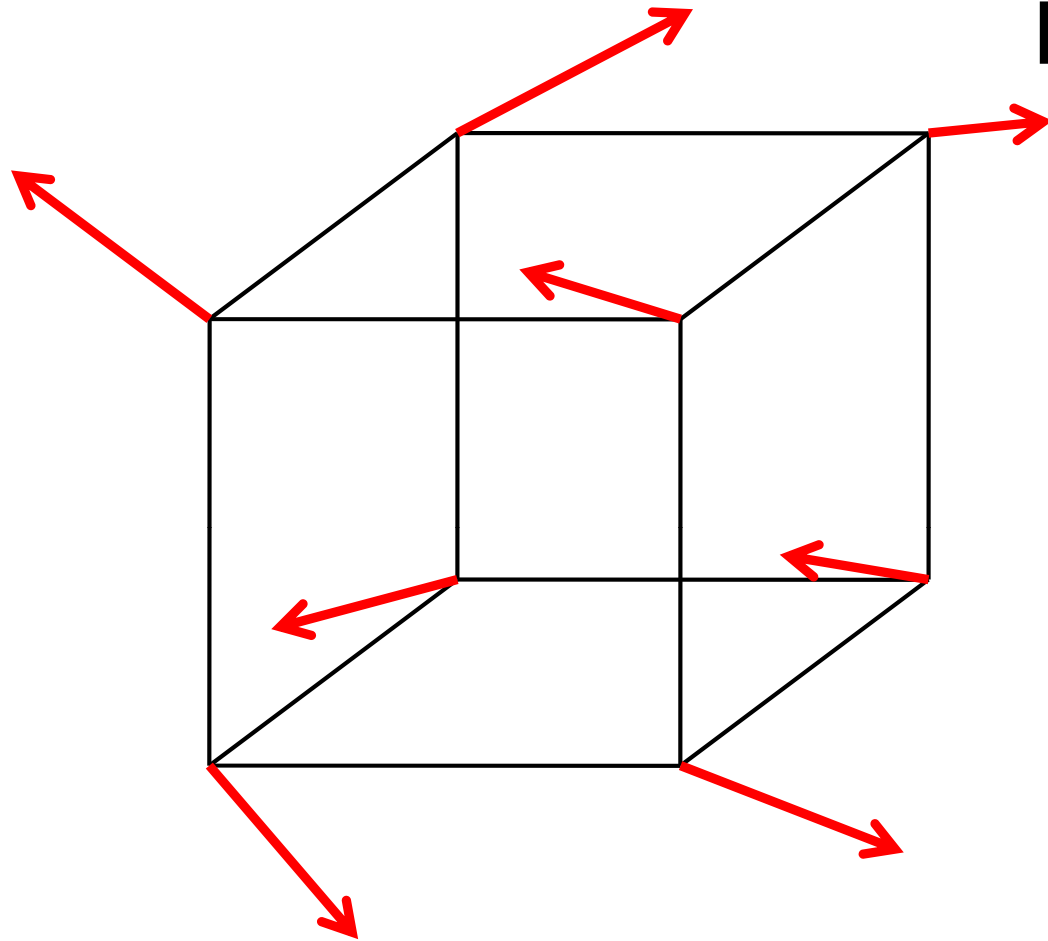
Repeat for all the triangles



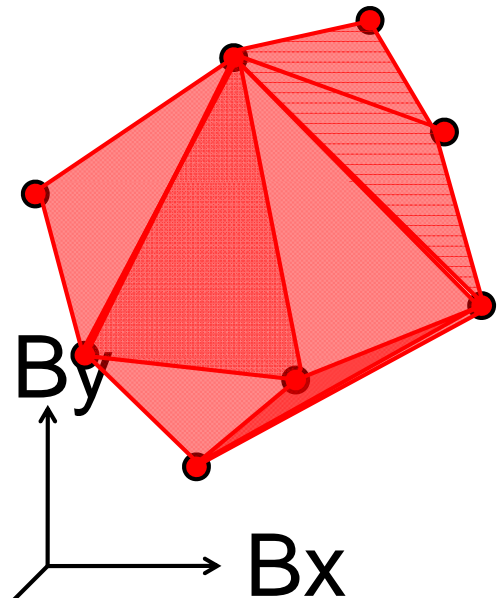
(x, y, z) (B_x, B_y, B_z)



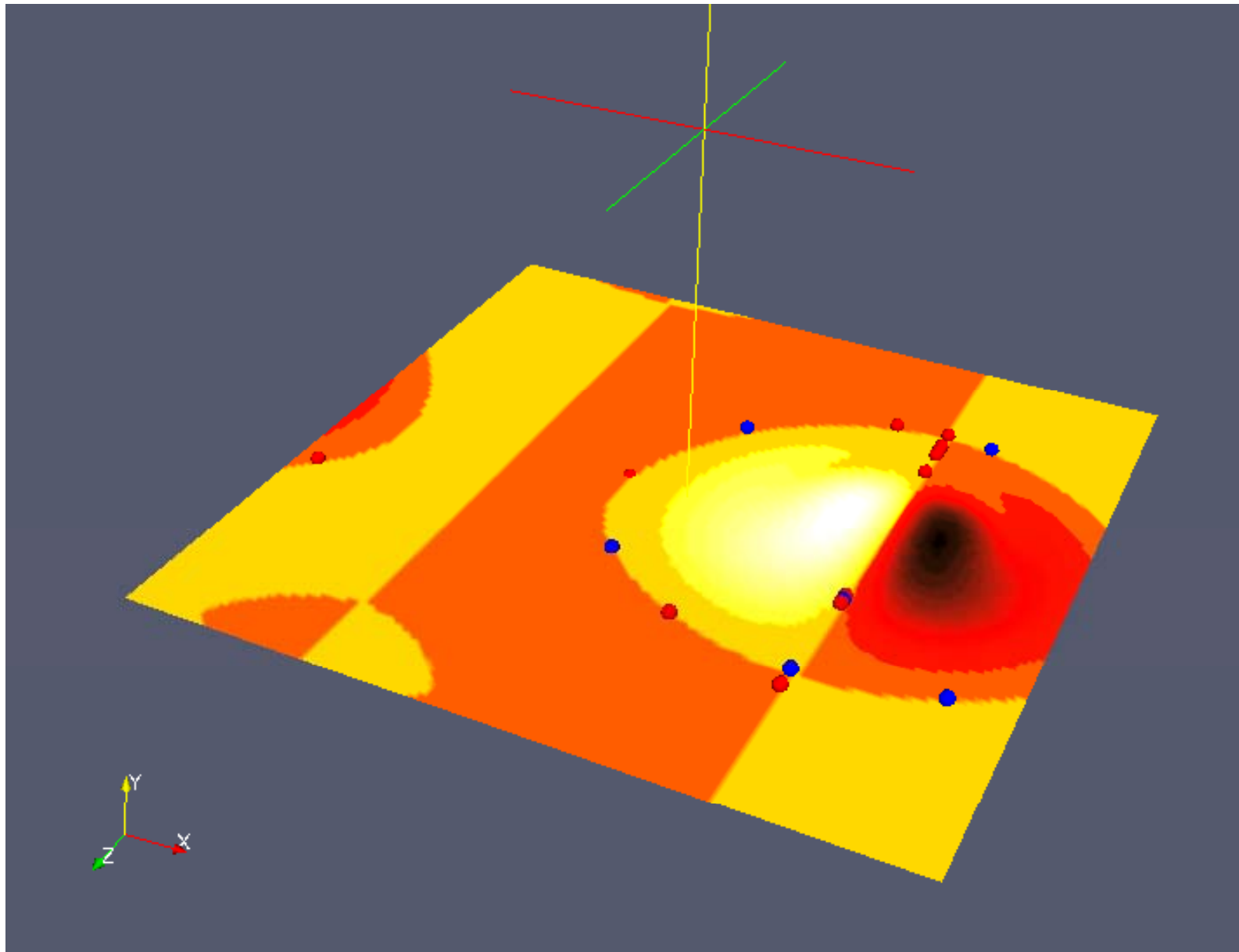
Does surface in B-space
enclose the origin?



(x, y, z) (B_x, B_y, B_z)



“String of pearls” pattern of nulls



Conclusions

- GEM-Challenge-like Scaling of R. Rate similar in 2D and 3D for the parameters in this study
- 3D system presents less back pressure because plasma can flow out in many directions direction
- Spread and drift of Reconnected flux along z makes comparison to 2D tricky
- Structure of 3D nulls is complex & has no direct 2D analog