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ELM Control in Tokamak Plasmas

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Stored energy vs 'scaling law'











> Introduction

- What is the Edge Localized Mode (ELM)?
- Theory of ELMs
- Why is ELM control urgent for ITER?
- Methods applied for Type-I ELM control/suppression
- > ELM control/suppression with magnetic perturbations
 - Application
 - Physics mechanism
- Combination of different ELM control methods

> Summary



What is the Edge Localized Mode (ELM)?







Theory of ELMs (I)



- ELMs are not well understood yet
- Ideal MHD modes driven by the steep current and pressure gradients at the edge transport barrier are regarded as the most likely candidates to explain their origin
- From stability calculations performed on the basis of experimental data three types of ideal MHD instabilities can be expected at the transport barrier:
 - ✓ kink-/peeling-modes
 - ✓ ballooning modes



H. Zohm, PPCF 38 (1996).

P.B. Snyder *et al,* Nucl. Fusion (2004)

 $\checkmark\,$ coupled peeling-ballooning modes



Theory of ELMs (II)







Theory of ELMs (III)





(*b*) A schematic showing the variation of pedestal stability boundaries with discharge shaping.

(c) Model of three types of ELM cycle.

P.B. Snyder *et al,* Nucl. Fusion **44** (2004) 320



Time (s)		n _e (10¹º m−³)
	Type-I	Type-III
$f_{\rm ELM}$ (Hz), typically	\sim 5-80	\sim 50-500
With increasing P_{SOL}	$f_{\rm ELM}$ increases	$f_{\rm ELM}$ decreases
n_e - T_e operational space	near edge pressure limit	slightly above L-H power
		threshold / at high n_e
$\Delta W_{ m ELM}/W_{ m ped}$	\sim 2-20 %	(<5%)

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Mixed Type-I and II ELM H-mode plasmas





J. Stober, et al., Nuclear Fusion, 45,1213 (2005)

Mixed Type I and II ELM H-mode has been observed in high δ and high density plasmas in JET

Pedestal Physics and ELM Behavious



Standard ELM-free H-mode plasmas

No ELMs, low edge transport \rightarrow good energy and particle confinement but also an impurity exhaust problem (not stationary)

Type-III ELMs plasmas

Relaxation oscillations with a high repetition frequency, sufficient particle exhaust and tolerable transient heat loads (rather high overall energy transport, leading to a degradation of the energy confinement of the plasma)

Type-II ELMs plasmas

Relaxation oscillations with a high repetition frequency, sufficient particle exhaust and tolerable transient heat loads. In contrast to type-III ELMs, they also provide good energy confinement. (a narrow operational window, and it is still unclear whether type-II ELMs will be possible to achieve in a burning fusion plasma)

Type-I ELMs H-mode plasmas

More or less strong relaxation oscillations with a low repetition frequency and have sufficiently low edge transport \rightarrow good compromise between high confinement and sufficient particle exhaust (unacceptably high transient heat loads expected in the divertor of a burning fusion plasma)





ELM Simulations on QSPA (0.1-0.6 ms, 30° to surface)

<0.4 MJ/m² Negligible erosion

0.4-1.0 MJ/m² (JET<1.0MJ/m²) Edge melting and surface cracking

1.0-1.6 MJ/m² Surface melting, bridge formation and droplet ejection





Why is ELM control urgent for ITER?







Using best estimates for divertor wetted area and inout asymmetry, one finds

 $\Delta W_{ELM} = Q_{ELM} \times S_{in} \times (1 + P_{out}/P_{in}) = 0.5 \text{ MJ/m}^2 \times 1.3 \\ m^2 \times 1.5 \sim 1 \text{ MJ}$

This requires a decrease in the 'natural' ELM size by a factor of ~ 20

ELM mitigation is required for a steady state operation of ITER!





Active control of Type-I ELM with acceptable confinement

degradation

- Radiating divertors (Impurity gas puffing)
- Magnetic triggering ("vertical kicks")
- Pellets pacing making
- Edge ergodization / external edge resonant magnetic perturbation (RMP) fields







P. MONIER-GARBET et al., Nucl. Fusion, 45, 1404 (2005)

Radiative dissipation of ELM energy is less than 20% (outer target) and less than 25% (inner target)



ELM size reduction by pellet injection JÜLICH

Type-I ELM frequency can be increased by injection of small deuterium pellets, provided that pellet freq. > 1.5 natural ELM freq. (results from AUG)

- Can the effects of plasma fuelling and ELM pacing be decoupled?
 - Can ELM pacing be demonstrated at N_GW ~ 0.75?



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Non-linear MHD simulations of pellets injected in the H-mode pedestal



JOREK



G T A Huysmans, PPCF 51 (2009)

• A strong pressure develops in the high density plasmoid, in this case the maximum pressure is ~5 times the pressure on axis.

• There is a strong initial growth of the low-*n* modes followed by a growth phase of the higher-*n* modes ballooning like modes.

 The coupled toroidal harmonics lead to one single helical perturbation centred on the field line of the original pellet position.

Simulations of pellets injected in the H-mode pedestal show that pellet perturbation can drive the plasma unstable to ballooning modes.





When doing pacing, due to the macroscopic pellet size this causes some fuelling,

 \rightarrow Additional convective losses reducing confinement



P. Lang, 16th ITPA PEP meeting 2009







Field penetration process

Mode excitation

Ergodisation

Rotation screening effect

3D equilibrium

NTV torque

Applications

Mapping Intrinsic field errors

RWM control

NTM control

Locked mode control

ELM control

Runaway electron control

Influence of RMP on sawtooth







- Depending on the relative phasing of the currents in individual coils, either *n*=1 or *n*=2 fields can be generated
- \succ $I_{\text{Coil}} \leq 3 \text{ kA x 16 turns}$ (*n* = 1 and 2)
- ➢ R ~ 6 m; Size ~ 6 m * 6 m
- $> B_r$ at wall ~ 0.25 mT/kAt

Y.Liang et al., PPCF 2007



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Type-I ELM control/suppression with RMP



ELM suppression window on DIII-D





- ELM suppression achieved in a narrow q₉₅ window on DIII-D with an *n*=3 field induced by the I-coils.
- \checkmark q₉₅ ELM suppression window can be enlarged slightly with a mixed *n*=1 and *n*=3 fileds.











Toroidal evolution of strike point



Connection length 79791 @ phi=0.00° -12010 -1308 Ę turns -1406 \mathbb{N} poloidal Height -1504 -1602 s*/ -1700 220 260 280 240 300 Major radius R [cm] Footprint on Tile 5: 79791 @ 63.400s сu 10 116 limiter 8 115 6 S 4 ģ 0 2 100 200 300 toroidal angle [deg]

•Field line tracing in vacuum approximation (superposition of equilibrium and perturbation field)

•No screening of RMP by poloidal rotation

•Ergodic field lines form lopes which generate multiple strike points on the divertor

•Strike point splitting depends on toroidal position

•Footprint represents N=2 symmetry of perturbation field

D. Harting, JET science meeting 2010

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O. Schmitz, PPCF (2008)

I. Joseph JNM, 2007

Splitting of the inner strike-point has been observed during ELM suppression with an n = 3 field on DIII-D.



Influence of magnetic perturbation on the Edge Electric field and rotation





With an n = 3 field applied,

edge $Er \rightarrow$ more positive;

K. Burrell, PPCF 47, B37, 2005

spin-up plasma rotation in co-current direction,

A large enhancement of the electron losses rather than ions by reason of the edge ergodisation.



Criterion for ELM suppression with RMPs





✓ Chrikov parameter number larger than 1 in the edge layer (sqrt(ψ) >0.925).





ELM frequency and temperature drop during ELM follow perturbation field amplitude (above threshold)

Heat and particle fluxes onto the divertor U JÜLICH



JET #69555

 $q_{95} = 4.4, \delta = 0.45;$

 $P_{\text{NBI}} = 9.5 \text{ MW}, n_e l = 1.3 (10^{20} \text{m}^{-2}),$

$$I_{\rm EFCC}$$
 = 32 kAt

✓ Reduction of ELM peak heat✓ No much effect on the inter-ELM heat flux

Outer Strike Line (Measured by embedded Langmuir probes) S. Jachmich, et al., EPS 2007

Influence of n = 1 field on profiles



EFCC *n* = 1; 135 degree; I_p = 1.6 MA; B_t = 1.84 T; q_{95} ~ 4.0; δ ~ 0.3 With n = 1 field Without n = 1 field #67951 5 8 (b) (a) 4 6 T_e(keV) T_i(keV) 3 2 t=17.120s t=17.125s t=17.725s t=17.725s 0 08 (d) (C) 5 ω_{tor}^{4} (10⁴rad/s) 6 n<mark>e(10¹⁹m⁻³)</mark> 5 4 2 =17.125s t=17.131s 0 t=17.725s t=17.631s -2 0 3.4 3.6 **R(m)** 3.2 3.8 3.2 3.8 ^{3.4} **R(m)** ^{3.6} field

Y.Liang et al., PRL 2007

✓ Electron and ion temperatures are increased during ELM mitigation phase

✓ Electron density decreases in the centre and at the edge due to pump-out effect

✓ Plasma braking observed during application of n = 1



Influence of n=1 field on confinement





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The minimum perturbation field amplitude for ELM mitigation increased but remained always below the *n*=1 locked mode threshold.



- \checkmark Drop of density at the plasma core and edge when the RMP field was applied.
- There is a threshold of density pump-out, However, it is different to the threshold of ELM control.
- Depends on the target plasmas
 - □ No clear density pump-out in L-mode, and type-III H mode plasmas
 - □ Less density pump-out in discharges with a less pump efficiency.
- ✓ No change of particle confinement in plasma core; (JET, TEXTOR).



Density pump-out effect can be compensated by either gas fuelling or pellet injection
 However, no recovery of energy confinement has been observed
 Y. Liang, 19th ITC (2009)

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> DIII-D results show not only to slow the plasma rotation, but also to accelerate the plasma, depending on the initial rotation.

> Similar plasma braking effect observed with n = 1 and n = 2 external fields on JET

Comparison between observed torque and NTV torque



Non-resonant magnetic braking Theory:



ELM suppression:

✓ RMP ELM suppression has been achieved in plasmas with ITER similar shapes and collisionalities on DIII-D

✓ Edge safety factor dependence of ELM suppression may limit the application for all ITER scenarios.

ELM Control:

Comparison the results between DIII-D and JET

DIII-D	JET
(<i>n</i> =3; i-coils)	(<i>n</i> =1, 2 EFCCs)

What are the same observations?

- ✓ Density pump-out
- ✓ Drop pedestal pressure and pressure gradient
- ✓ Plasma rotation braking

What are the different observations?

- ✓ ELM suppression
- ✓ A single narrow q_{95} window
- ✓ ELM control (frequency/size)
- ✓ A wide q_{95} window

 $I_{\rm p}$ = 2MA; $B_{\rm t}$ = 1.85T q_{95} = 3.1; low δ

 $P_{NBI} / \longrightarrow f_{ELM} /$

The power dependence of the ELM frequency is similar to normal type-I ELMs. However, the mitigated ELMs with n = 1field have a higher frequency and smaller in size.

Y. Liang et al., NF, 2010

> Pedestal n_e is reduced by ~20% while the edge T_e is increased. ∇p_e is ~20% smaller.

With n = 1 perturbation field the operational point moves from intermediate-n peeling-ballooning (wide mode) boundary to low-n peeling (narrow mode) boundary.

Resonance effect in ELM frequency vs q_{95}

 \succ ELM control with n = 1field is very sensitive to the edge safety factor.

Small change of q_{95} from 4.5 to 4.8 results in an increase of f_{ELM} by a factor of 2-3 and a drop of $n_{\rm e}l$ by 15% while almost no difference is observed without n = 1 field.

Plasma rotation braking from the n = 1 field does not depend on q_{95} .

Y Liang Submitted to PRL (2010)

- > Multiple resonances in f_{ELM} vs q_{95} have been observed with n = 1 and 2 fields
- Possible explanation in terms of ideal peeling mode model by Gimblett et al [C G Gimblett et al., PRL, 96, 035006-1-4(2006)] currently being investigated

The mechanism of edge ergodisation, can not explain the multi-resonance effect observed with the low *n* fields on JET.

What is the physics mechanism of ELM suppression with magnetic perturbations?

DIII-D *n*=3 Even parity

B.) Spectrum

Optimisation of stochastic edge region

DPG 2010

What is the role of the magnetic perturbation spectrum?

DIII-D

Upper invessel coils only

Both Upper and lower In-vessel coils

External C coils

3D effect of perturbation fields on the plasma equilibrium

3D equilibrium code IPEC

Magnetic flux surfaces
 of the target plasma can
 be perturbed by each
 dominant error field.
 It suggests 3D effect
 need to be included in the
 stability analysis.

Jong-kyu Park, PRL 2007

Influence of magnetic perturbation on X-point

C. Wiegmann, et al, EPS2009, P1.132

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R[m]

RMP experiments

Combination of different ELM control methods

- RMP + vertical kicks
- RMP + pellet injection
- RMP + impurity gas puffing

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However, no recovery of energy confinement has been observed

Arriving pellet trigger ELMs (and do fuel, here it was welcome)

P. Lang, 16th ITPA PEP meeting 2009

Summary (I)

Active control of ELMs by resonant magnetic perturbation fields offers an attractive method for next-generation tokamaks, e.g. ITER.

- ✓ D-III D has shown that type-I ELMs are completely suppressed when n = 3 magnetic perturbations are applied.
- Increasing of ELM frequency or ELM triggering has been observed on JET, MAST and NSTX, but not DIII-D with mid-plane C-coils.
- ✓ Up to date, no complete ELM suppression was obtained on JET, MAST even with a Chirikov parameter larger than 1 at $\Psi_{pol}^{1/2}$ > 0.925 which is one of the important criterions for the design of ITER ELM suppression coil.
- Density pump-out effect with application of RMP from midplane coils has been observed on JET, MAST and NSTX, but not DIII-D with mid-plane C-coils. It can be compensated by either gas fuelling or pellets injection. However, no recovery of energy confinement has been observed.
- Plasma response (screening and 3D equilibrium) helps for understanding the mechanism of ELM suppression/control with magnetic perturbations

✓ Radiating divertors (type-III ELM), successful ELM control and full Hmode confinement have still to be demonstrated.

✓ Magnetic triggering ("vertical kicks") need in-vessel coils. Promising technique for ILW on JET, in which case the ELM size need only be reduced by ~ 2-3 times

 ✓ Pellet pacing can typically achieve a factor of two reduction in the energy per ELM – this is not enough. Also, for ITER the reliability of a pellet system, for a safety application, has to be questioned.

 ✓ External magnetic perturbation Very promising results up to now and further development needed in the future. Joint experiments (DIII-D, MAST, TEXTOR, AUG, …) will help to understand physics

✓ ITER may need combination of different ELM control methods

Thanks for your attention!

Fusion ...

... on Earth

We are on the way