

# Continuum Absorption at Frequency Tip

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# Continuum Absorption

- Occurs in *inhomogeneous* plasma where Alfvén velocity is a function of position ( $\mathbf{x}$ ).
- When external driving frequency matches with local shear Alfvén frequency, phase mixing occurs around the resonance and wave energy is dissipated to the plasma.
- The energy absorption rate is calculated by introducing an artificial growth rate to mimic the effects of viscous dissipation. (e.g., Chen and Hasegawa 1974)

# Estimation of Absorption Rate

$$\frac{d}{dx} \left( \rho [\tilde{\omega}^2 - \omega_A^2(x)] \frac{d\xi_x}{dx} \right) - k_{\perp}^2 \rho [\tilde{\omega}^2 - \omega_A^2(x)] \xi_x = 0$$

$$\tilde{\omega} = \omega + i\gamma$$

$$\frac{d}{dx} \left( \rho [\tilde{\omega}^2 - \omega_A^2(x)] \frac{d\xi_x}{dx} \right) = 0$$

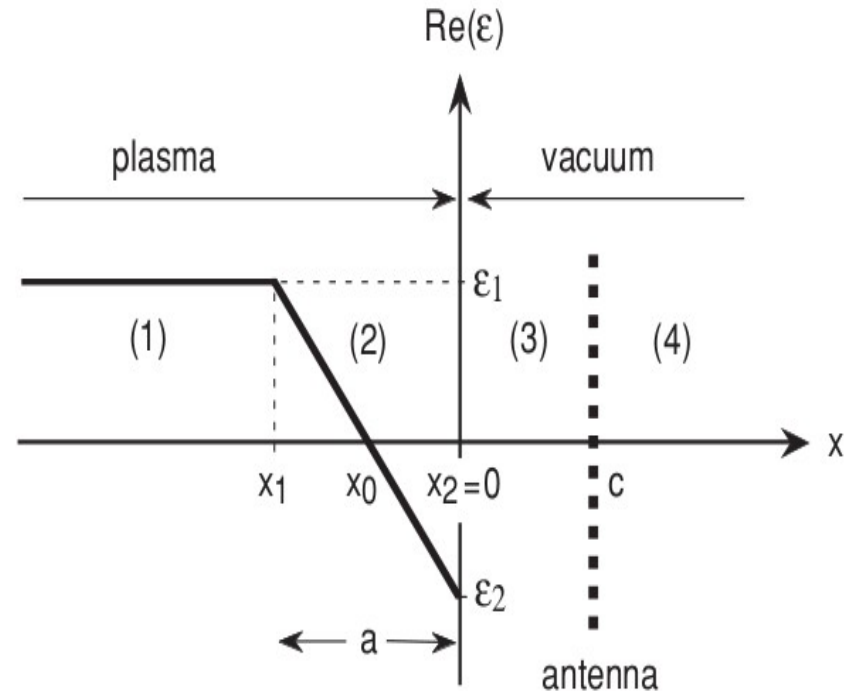
$$\rho [\tilde{\omega}^2 - \omega_A^2(x)] \frac{d\xi_x}{dx} = C$$

$$\frac{d\xi_x}{dx} = \frac{C}{\tilde{\omega}^2 - \omega_A^2(x)} = \frac{C}{\omega^2 - \omega_A^2(x) + 2i\omega\gamma} = \frac{C}{\frac{d\omega_A^2(x)}{dx}(x-x_0) + 2i\omega\gamma}$$

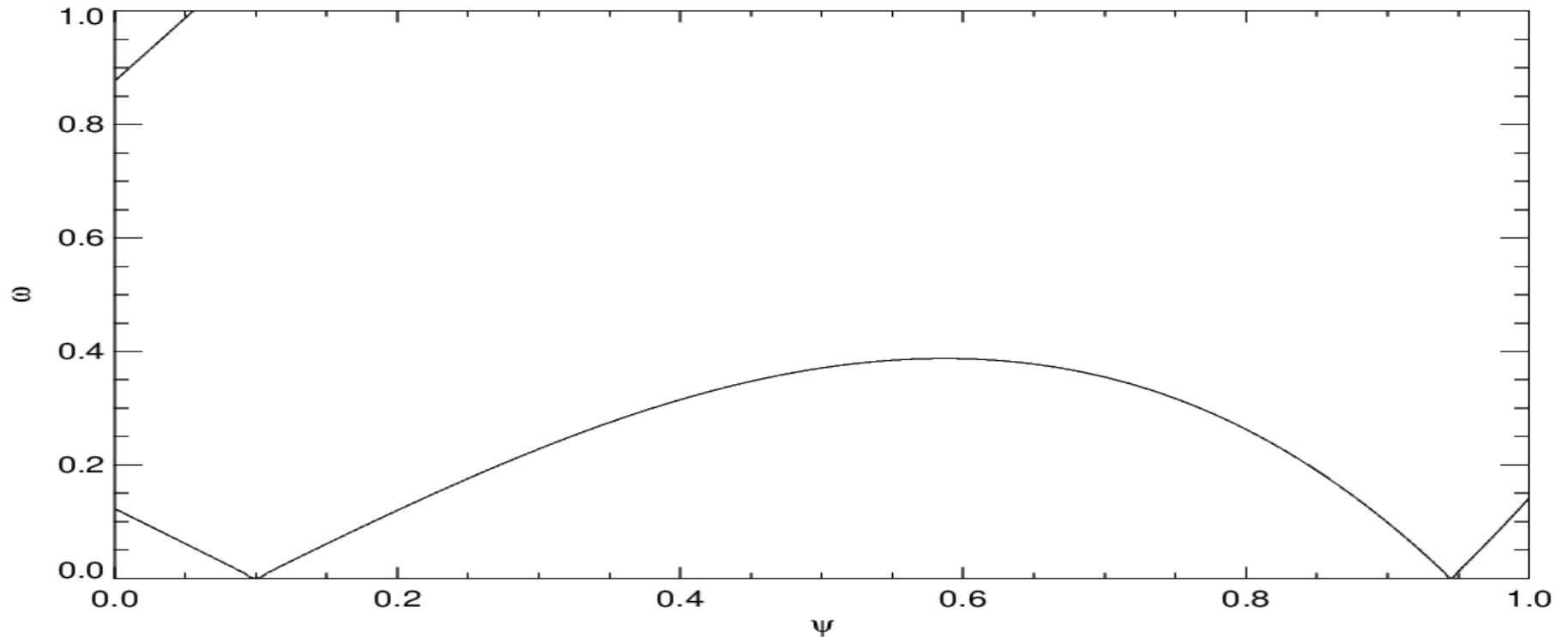
$$Q = - \int \rho \gamma v^2 dV \approx - \int \rho \gamma v_s^2 dV \approx - \int \rho \gamma \left( \frac{d\xi_x}{dx} \right)^2 dV \approx \int \rho \gamma \omega^2 \left( \frac{d\xi_x}{dx} \right)^2 dV$$

$$= \int \frac{\rho \gamma \omega^2 |C|^2}{\left| \frac{d\omega_A^2(x)}{dx}(x-x_0) + 2i\omega\gamma \right|^2} dV = \int \frac{\rho \gamma \omega^2 |C|^2}{\left( \frac{d\omega_A^2(x)}{dx} \right)^2 (x-x_0)^2 + 4\omega^2 \gamma^2} dV \propto \left| \frac{d\omega_A^2(x)}{dx} \right|^{-1}$$

Q is independent of  $\gamma$ , proportional to  $\left| \frac{d\omega_A^2(x)}{dx} \right|^{-1}$



# Slope of the Continuum is usually not a constant...

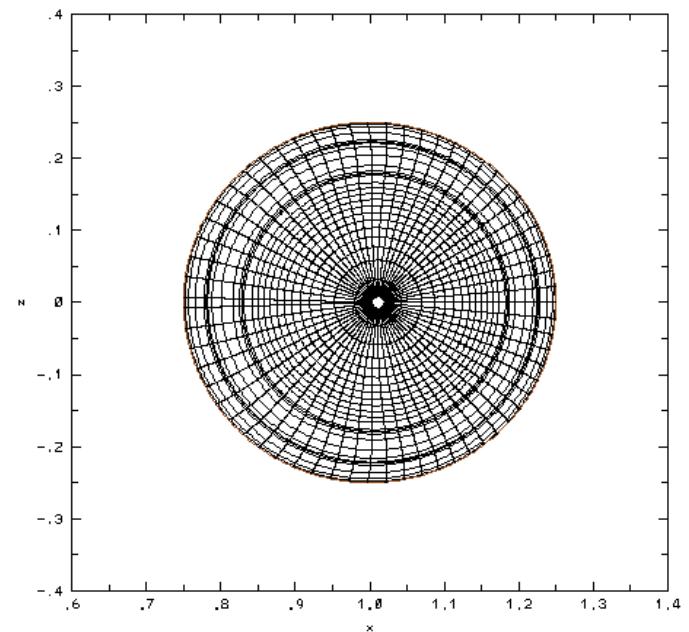
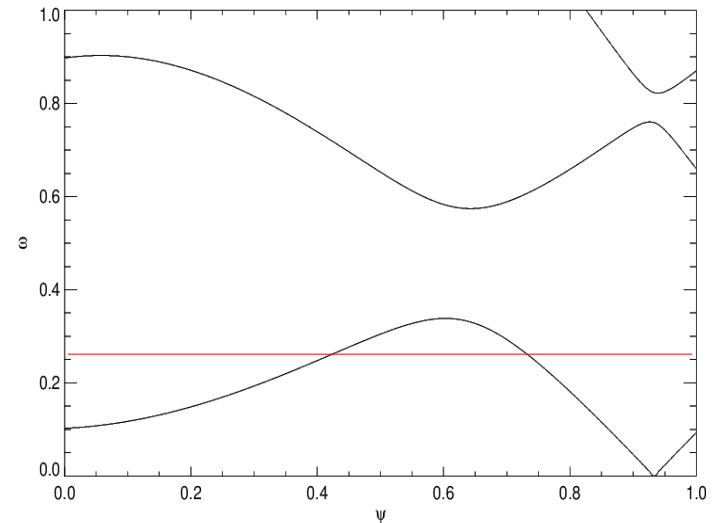


$$\left| \frac{d\omega_A^2(x)}{dx} \right|^{-1} \rightarrow \infty$$

What is the absorption rate as the slope approaches zero??

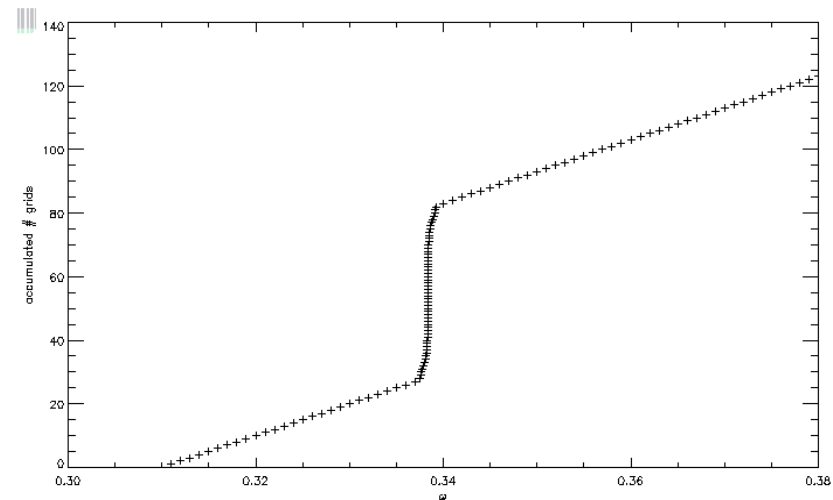
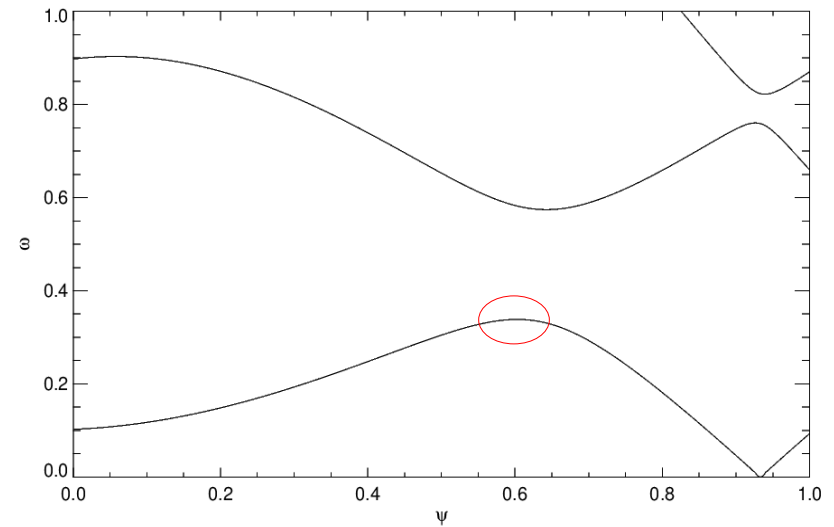
# Numerical Investigation (I)

- We use a Grad-Shafranov Solver TOQ to generate the numerical equilibria. The input parameters are poloidal current flux and pressure profile.
- Our ideal MHD code AEGIS is adaptive, thus, high density of grids are placed in the vicinity of continuum.



# Numerical Investigation (II)

- A poloidal current is driven at  $x=0.2$  with a tunable frequency.
- We investigate the absorption as a function of frequency near to the frequency tip.



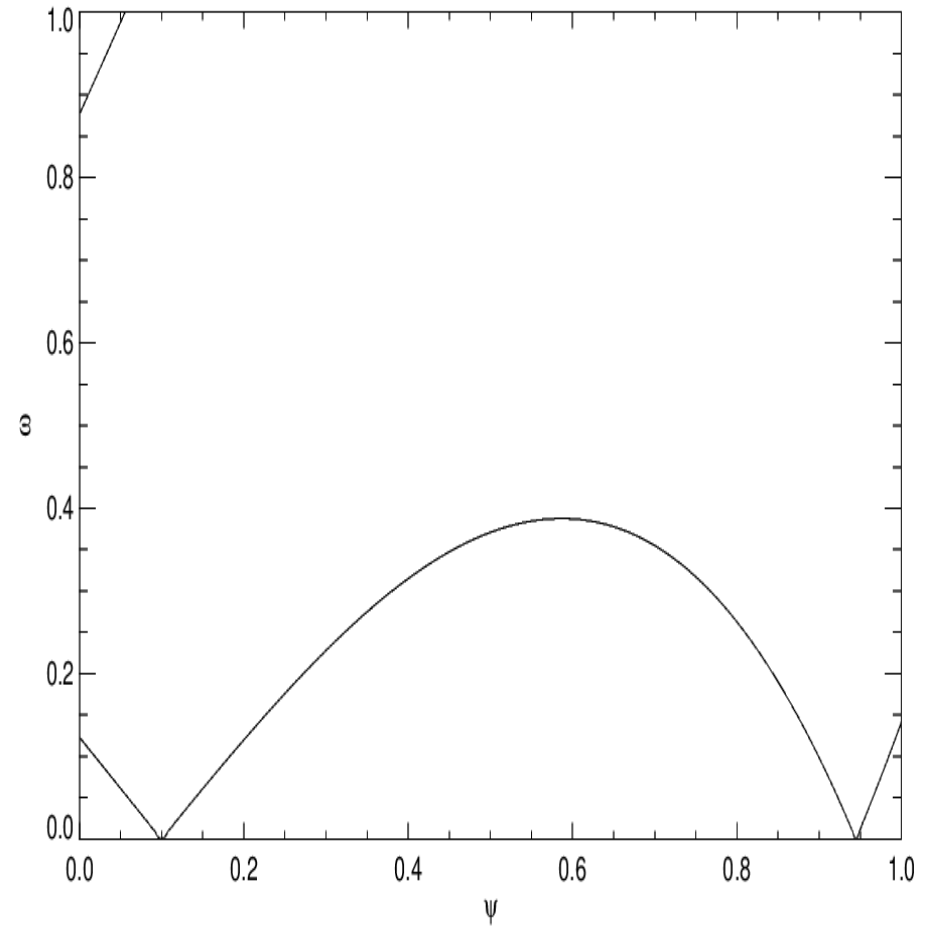
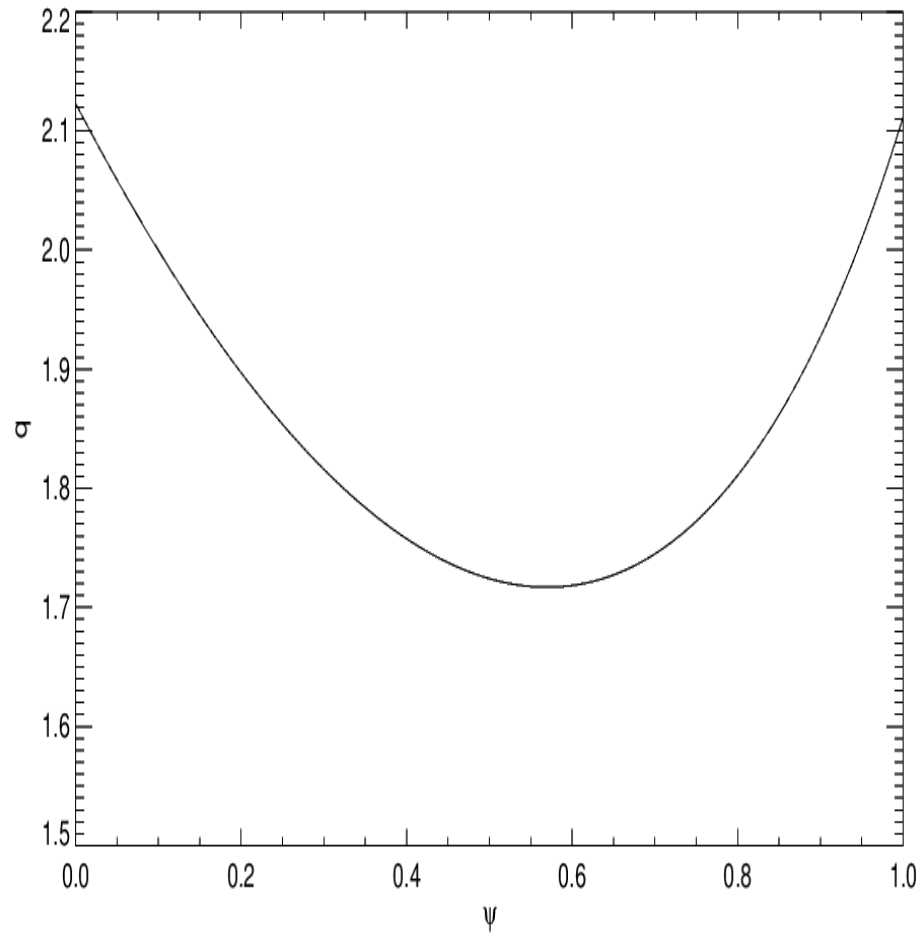
# Estimation of Absorption Rate at Continuum Tip

$$Q = \int \frac{\rho\gamma\omega^2|C|^2}{\left|\frac{d\omega_A^2(x)}{dx}(x-x_0) + 2i\omega\gamma\right|^2} dV = \int \frac{\rho\gamma\omega^2|C|^2}{\left(\frac{d\omega_A^2(x)}{dx}\right)^2 (x-x_0)^2 + 4\omega^2\gamma^2} dV \propto \left|\frac{d\omega_A^2(x)}{dx}\right|^{-1}$$



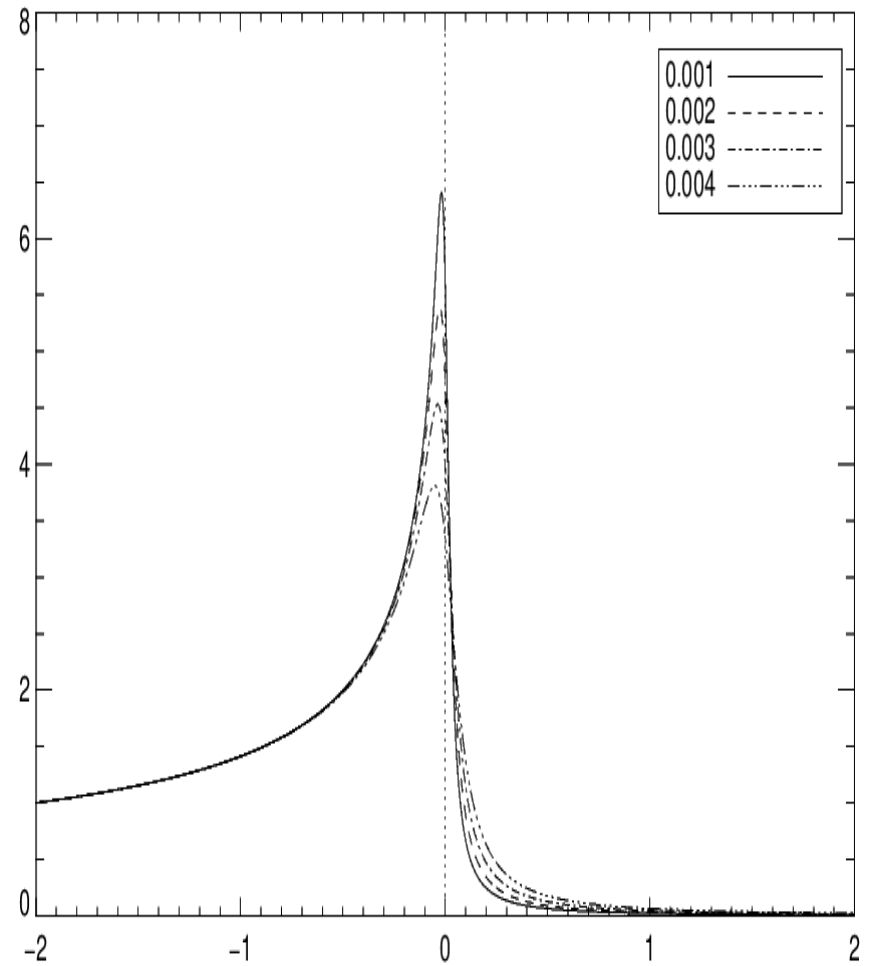
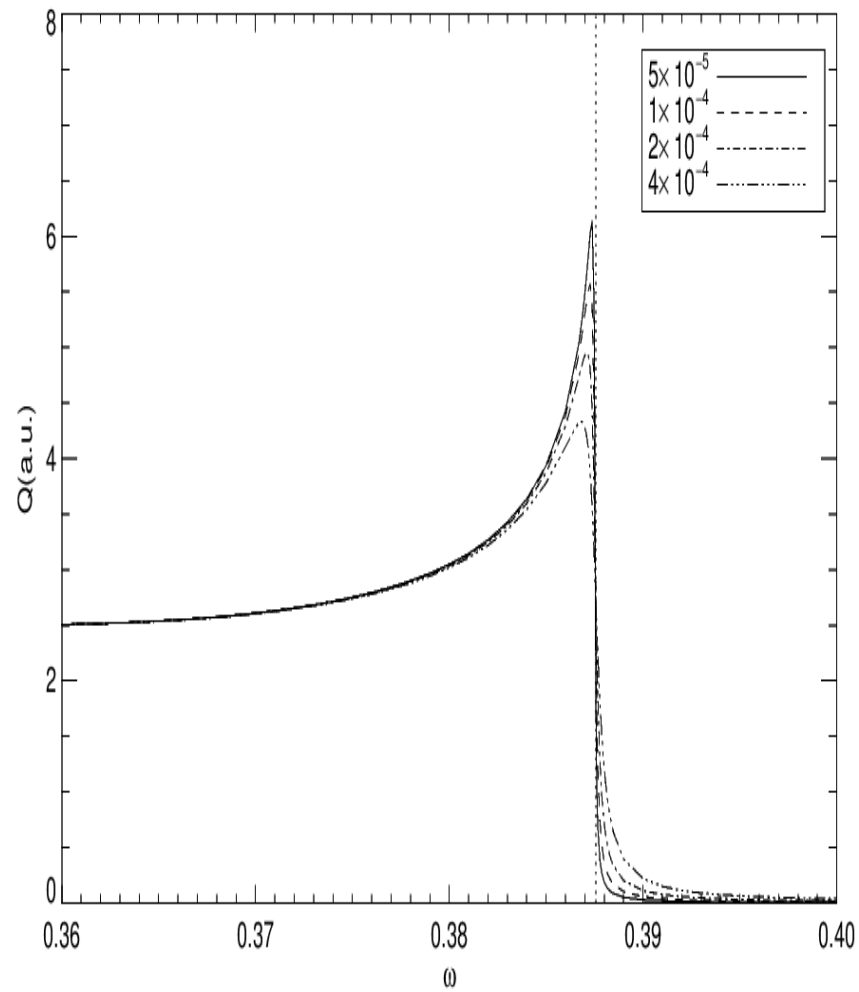
$$Q = \int \frac{\rho\gamma\omega^2|C|^2}{\left|\frac{1}{2}\frac{d^2\omega_A^2(x)}{dx^2}(x-x_0)^2 + 2i\omega\gamma\right|^2} dV = \int \frac{\rho\gamma\omega^2|C|^2}{\frac{1}{4}\left(\frac{d^2\omega_A^2(x)}{dx^2}\right)^2 (x-x_0)^4 + 4\omega^2\gamma^2} dV \propto \left|\frac{d^2\omega_A^2(x)}{dx^2}\right|^{-\frac{1}{2}} \gamma^{-\frac{1}{2}}$$

# Reverse Shear Tip (I)

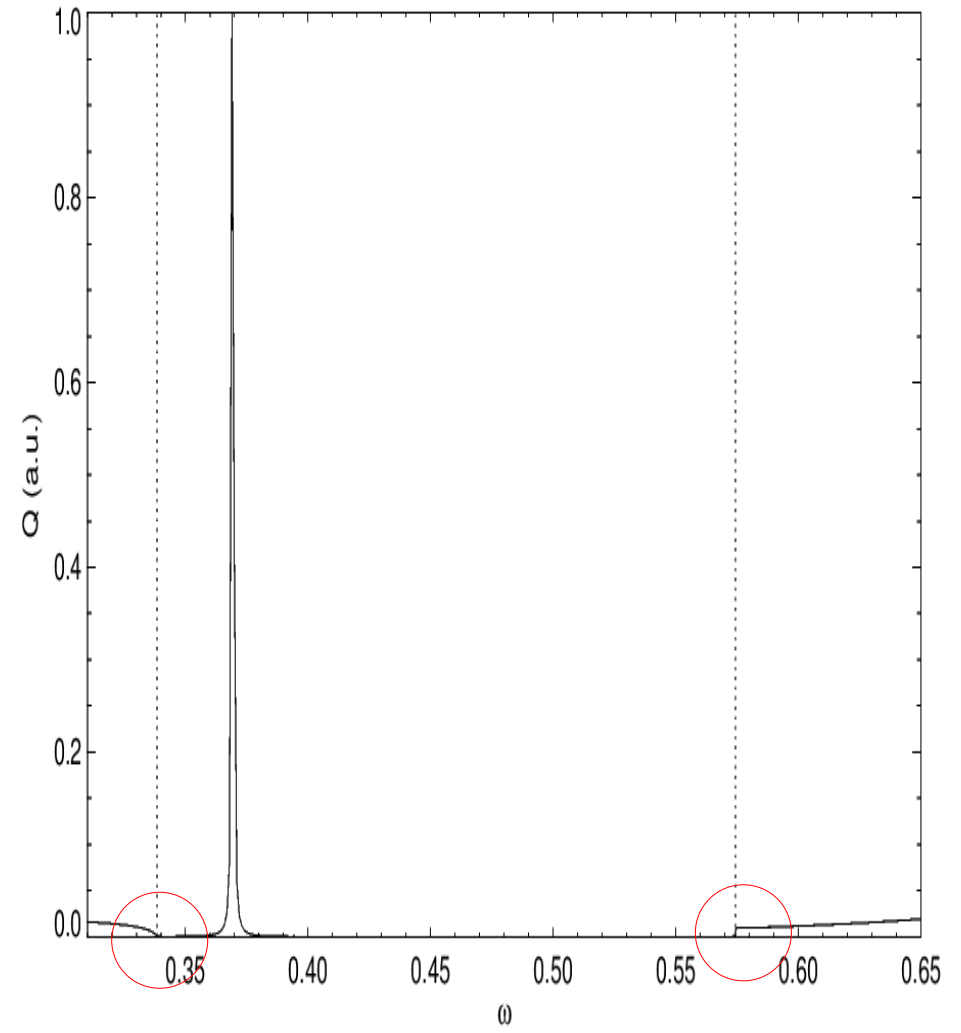
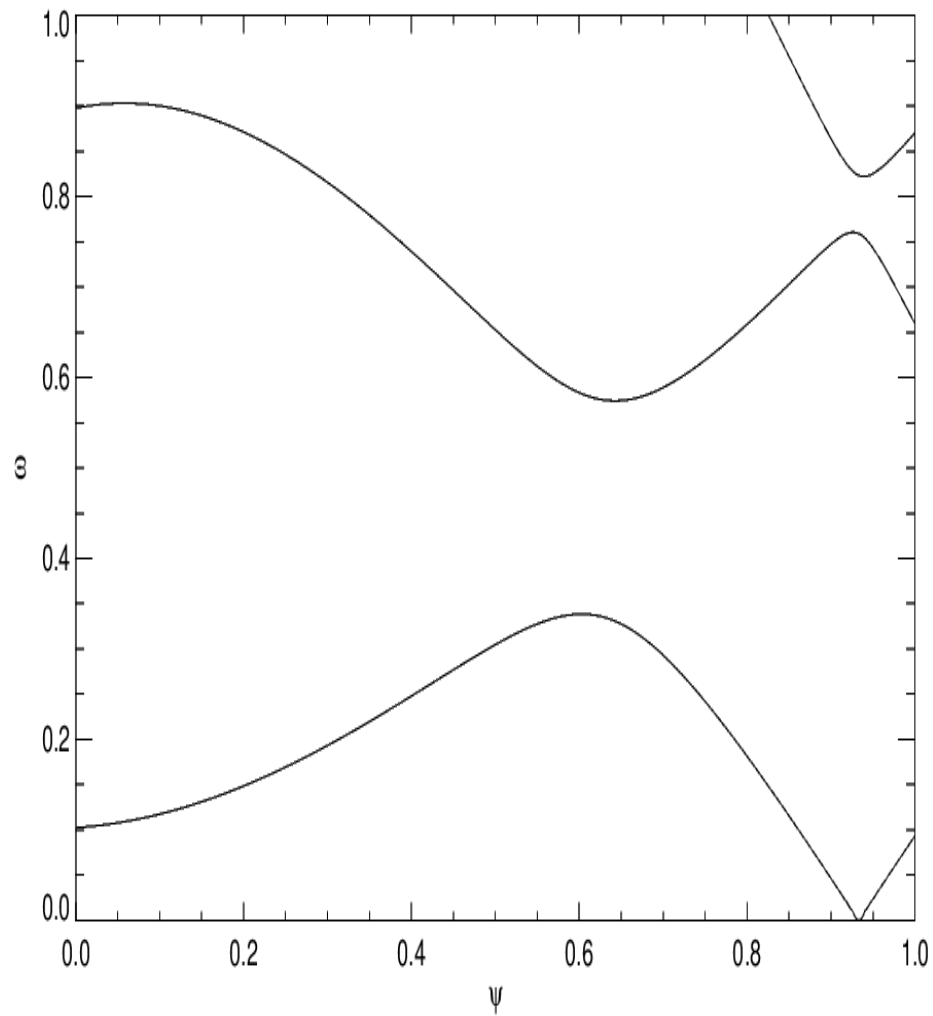




# Reverse Shear Tip (II)



# TAE Frequency Gap



# Estimation of Continuum Absorption at TAE Tip

- The tip is formed due to the toroidal coupling of two adjacent poloidal harmonics. Both harmonics need to be taken into account.

$$v_s \approx i\omega \frac{\begin{bmatrix} z-g & -1 \\ 1 & z+g \end{bmatrix} \begin{bmatrix} C_m \\ C_{m-1} \end{bmatrix}}{1-g^2+z^2}$$

- We applied the TAE tip model (Rosenbluth, Berk et al.; Breizman and Sharapov) to the problem.

- It can be shown that as  $|g| \rightarrow 1$

$$Q \sim \sqrt{(1-g^2)} \text{ where}$$

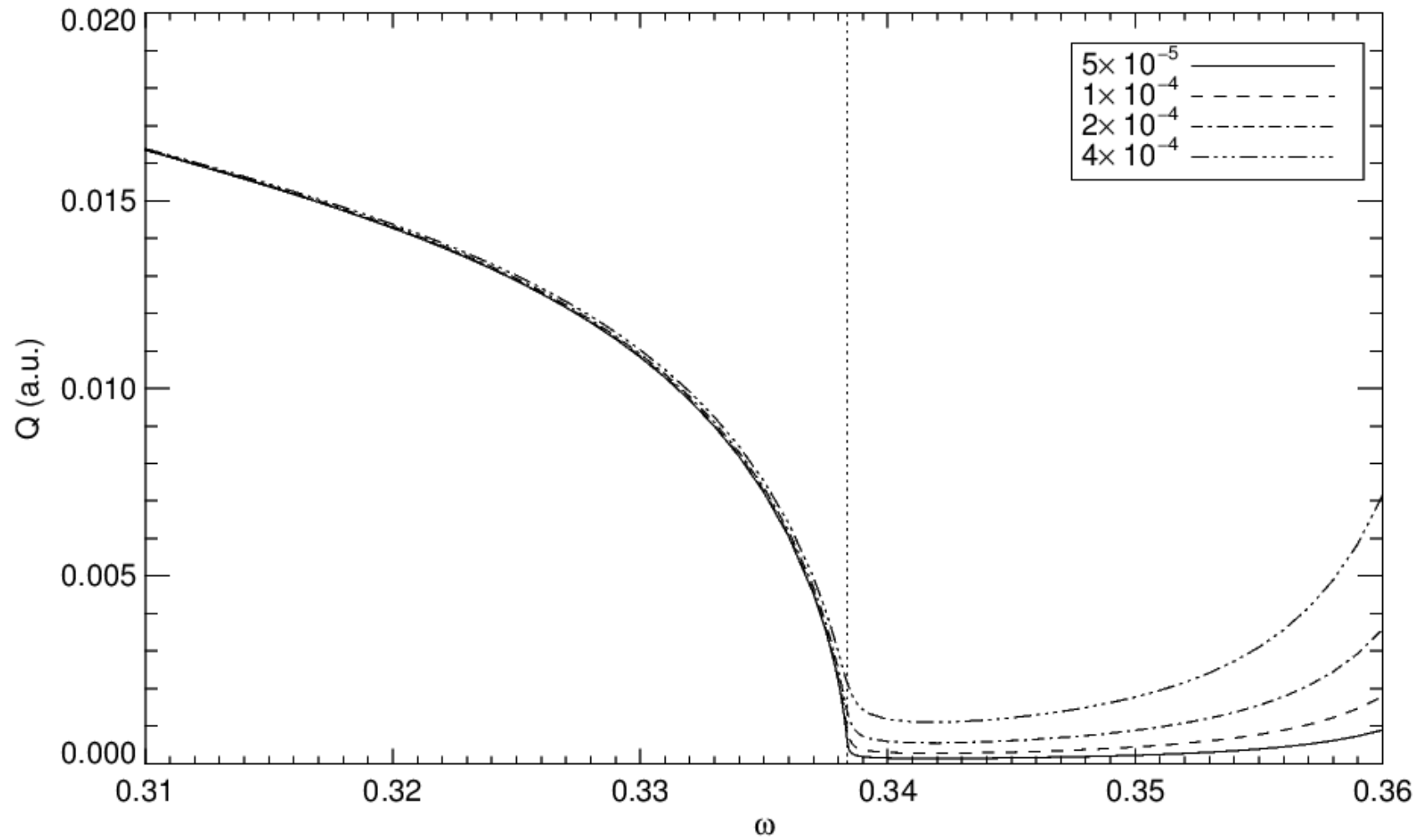
$$g = \frac{1}{\varepsilon} \left( \frac{4q^2 R^2 \tilde{\omega}^2}{v_A^2} - 1 \right)$$

- Absorption shall vanish at the tip!

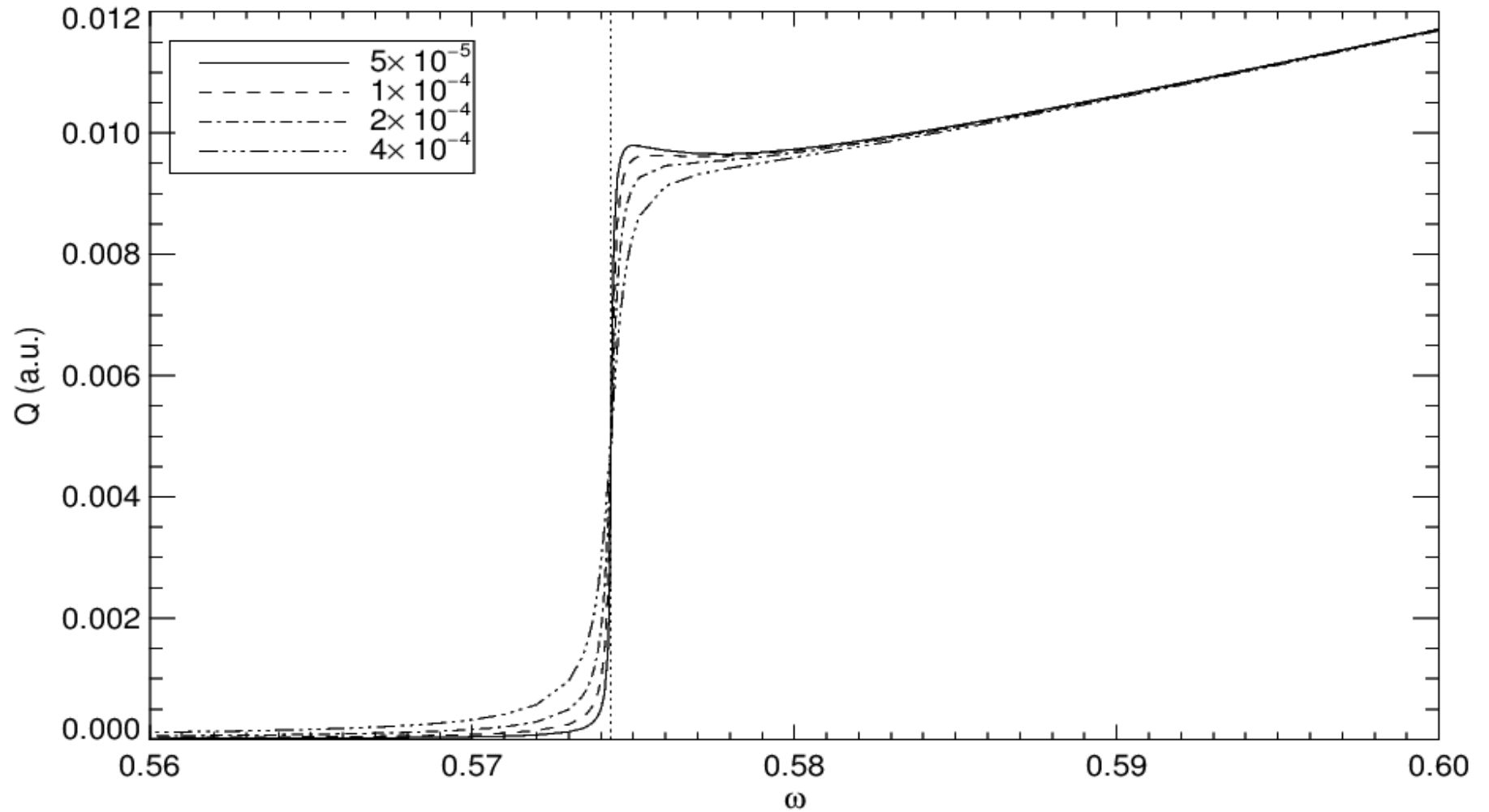
$$\begin{bmatrix} C_m \\ C_{m-1} \end{bmatrix} \approx \frac{\pi}{\sqrt{(1-g^2)}\pi^2 + \frac{g\pi^3 S}{2}} \begin{bmatrix} g\tilde{J}_m + \tilde{J}_{m-1} \\ -\tilde{J}_m - g\tilde{J}_{m-1} \end{bmatrix}$$

$$Q \sim \int \gamma \omega^2 \frac{[(1-g^2)\tilde{J}_m]^2 + [(1-g^2)\tilde{J}_{m-1}]^2}{[\sqrt{(1-g^2)}(\pi - \frac{\pi^3 S^2}{16} + \frac{g\pi^2 S}{2})](1-g^2+z^2)^2}$$

# Lower TAE Tip



# Upper TAE Tip

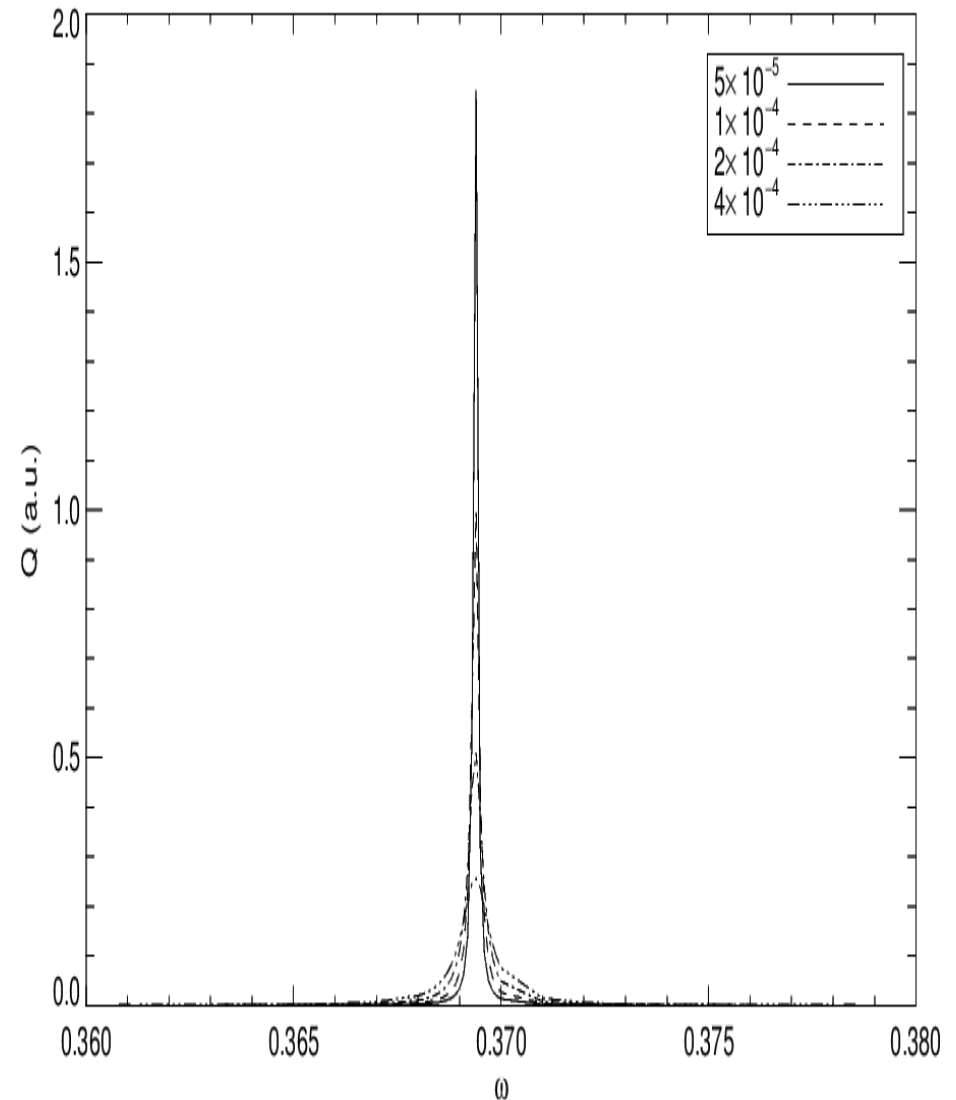


# TAE Frequency

- For eigenmodes,  $v \propto \frac{1}{\gamma}$ ,  
independent of position.

$$Q \sim \int \gamma v^2 dV \sim \gamma^{-1}$$

Numerical results show good agreement.



# Summary and Conclusions

- In the case of reverse shear tip and lower TAE tip, analytic results and numerical results are in adequate agreement with each other.
- The response of upper TAE tip cannot be obtained in the TAE tip model of Rosenbluth and Berk et al. / Breizman and Sharapov.
- The numerical results around TAE frequency are in agreement with analytic model.
- We are modifying the tip model to see why the discontinuous absorption at upper TAE tip is achieved.