

# Measurement and theoretical modelling of the damping rate of medium-n TAEs in JET

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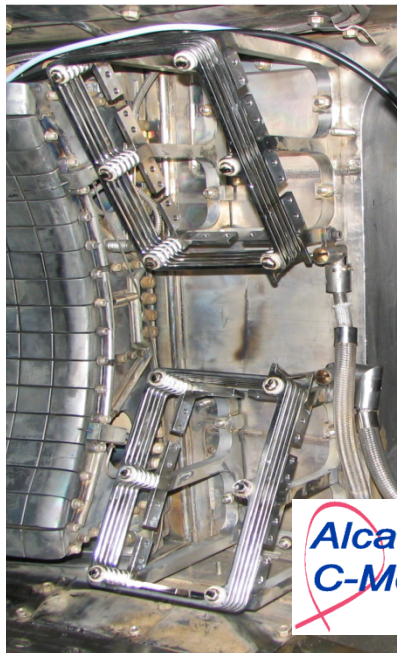
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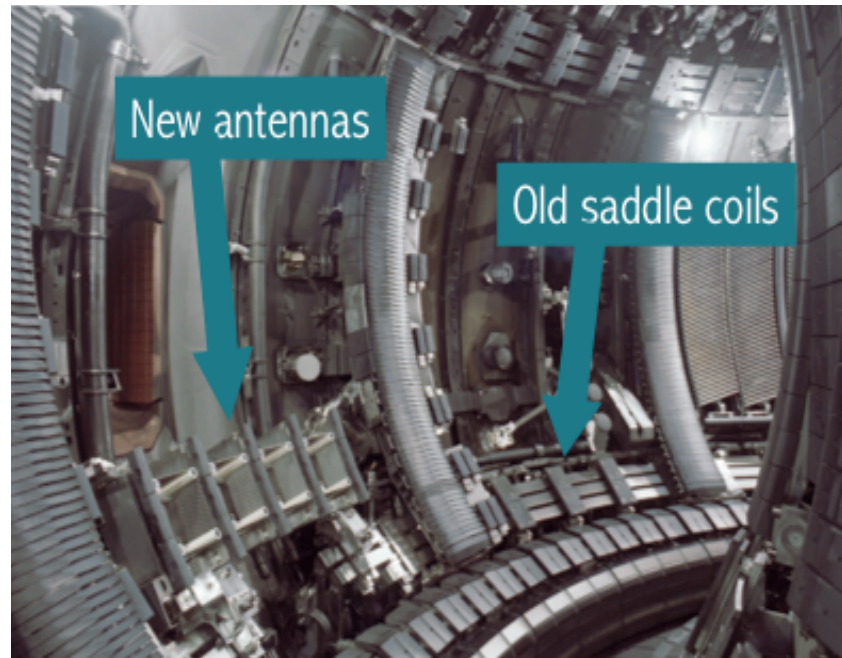
# Alfvén Eigenmode Active Diagnostics

Aim: address physics of mode damping, identify modes most prone to instability in different burning plasma scenarios, and parameters to control stability

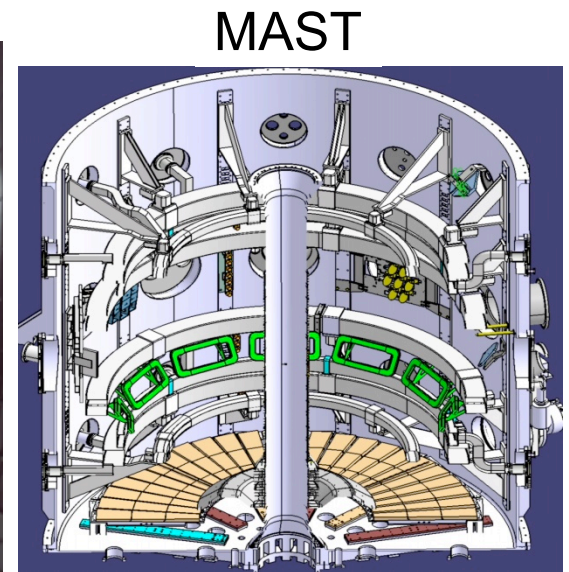


Alcator  
C-Mod

high field & density,  $T_e \sim T_i$

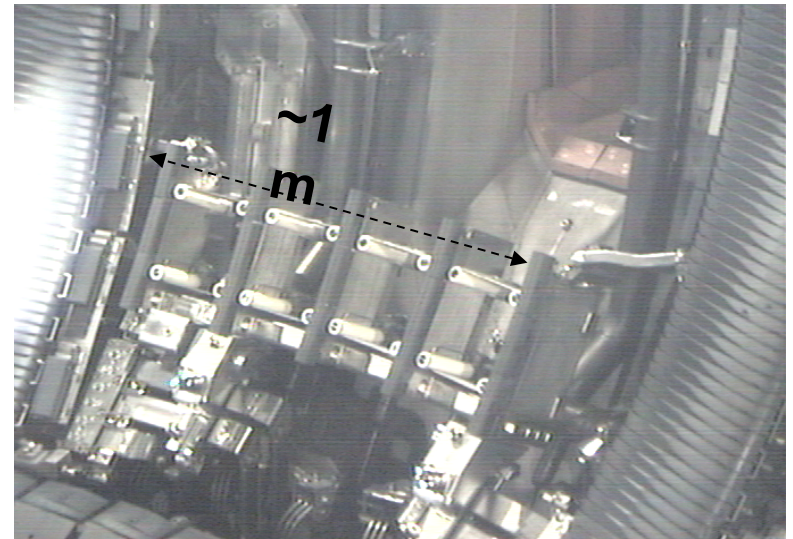


ITER-relevance for size and shape scaling, scenarios

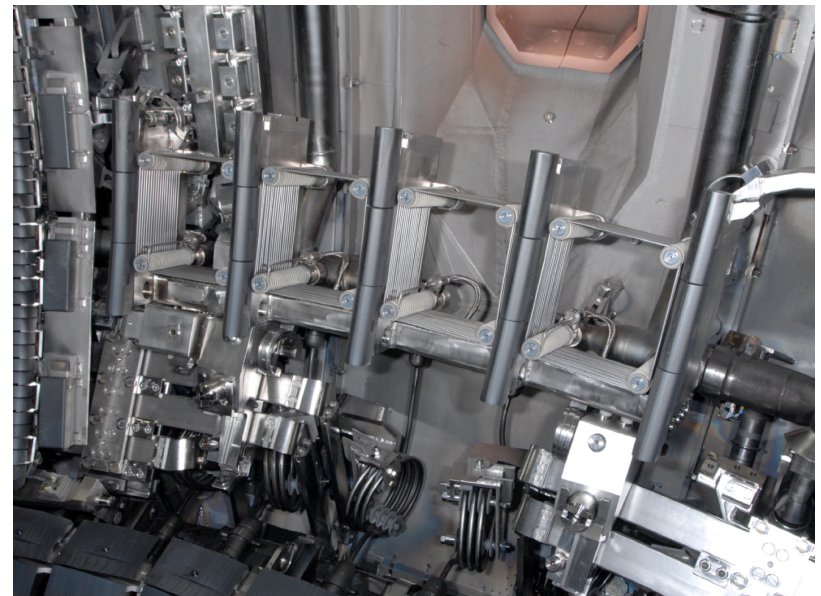


MAST  
tight aspect ratio,  
broad range of  $\beta$

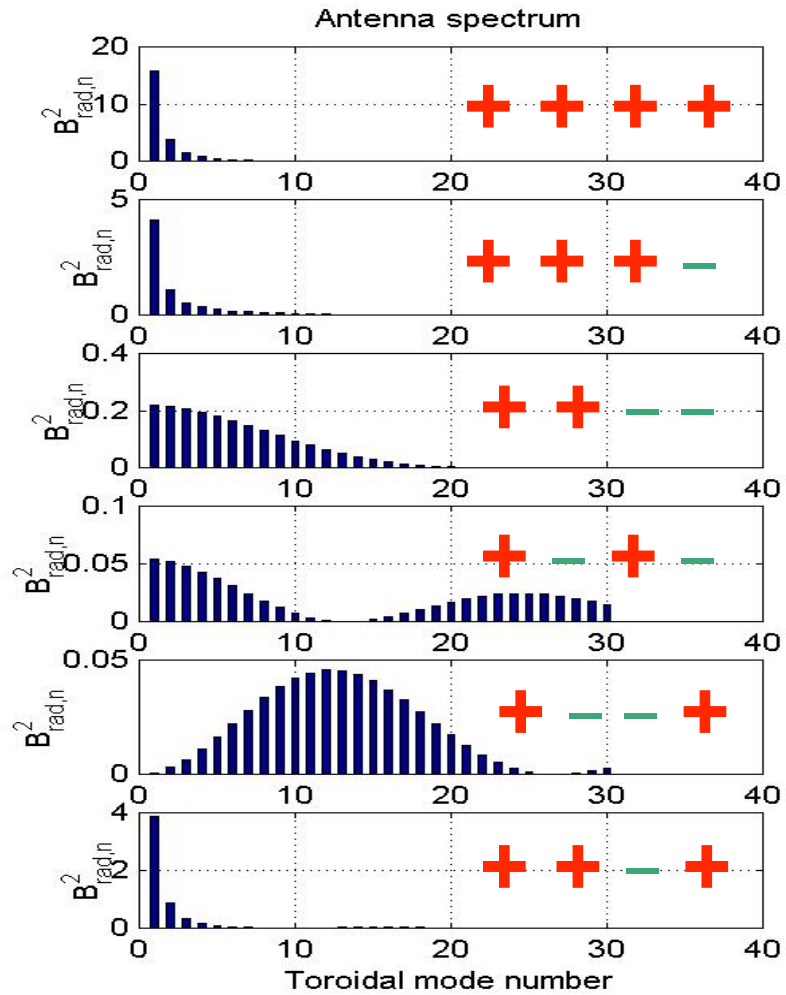
# JET AE antennas



Oct 8  
(2005)



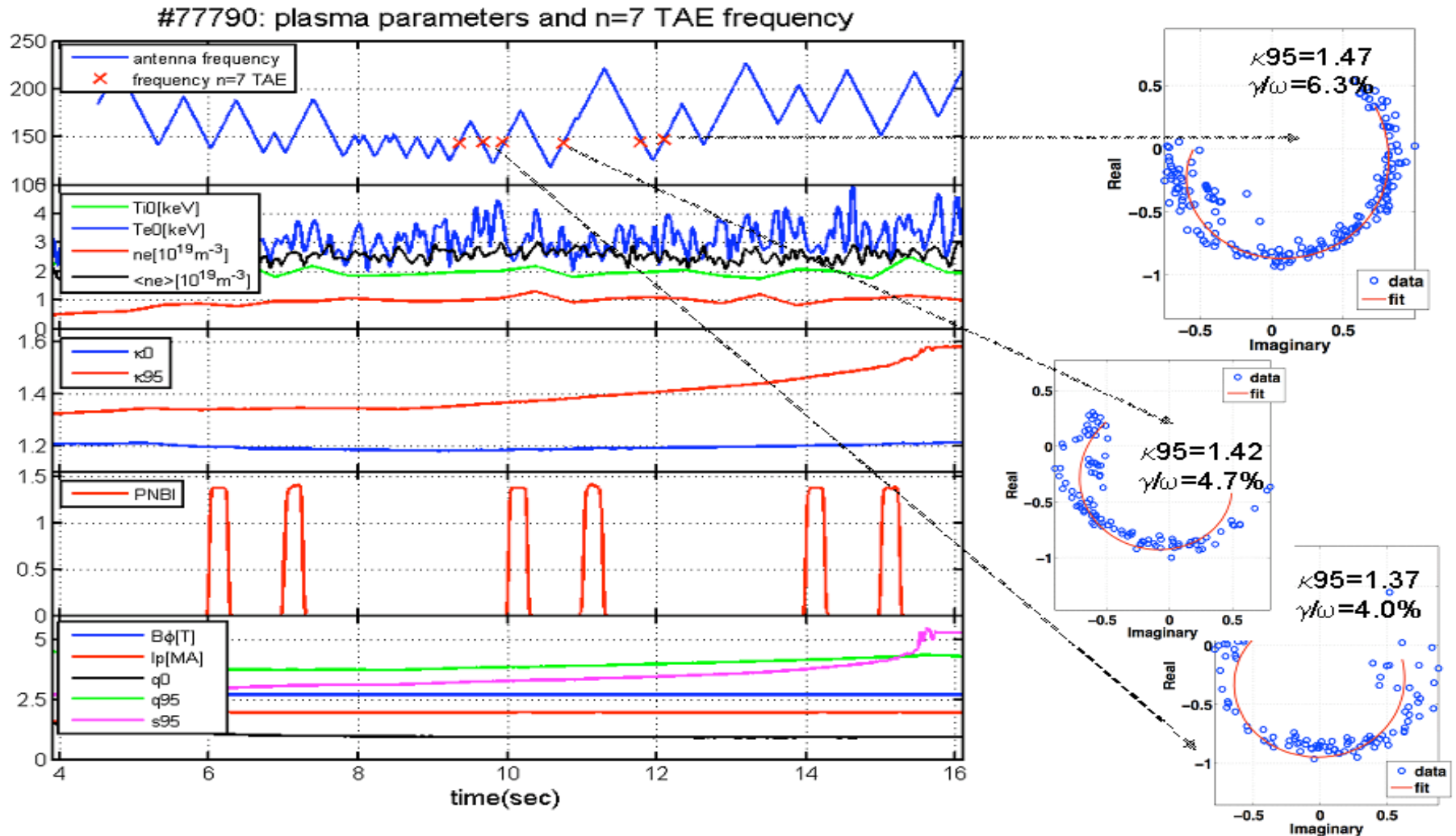
Oct 4  
(2007)



Unique capability: real-time tracking to follow mode evolution as plasma parameters change



# Ex. of $\gamma/\omega$ measurements for n=7 TAE



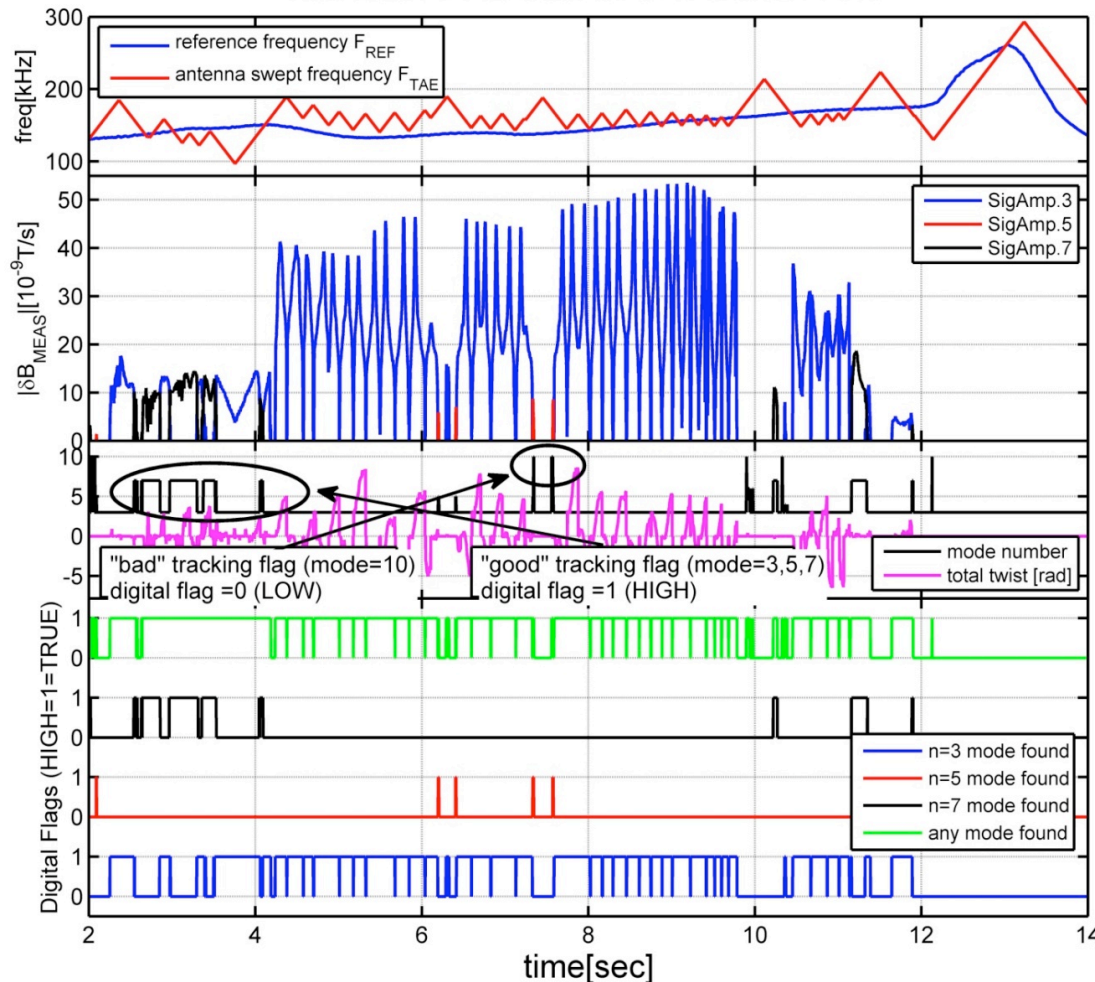
Note: in the absence of fast ion drive,  $\gamma/\omega$  is the mode damping rate

# Outline

- Real-time mode decomposition and tracking
- Measurements of TAE damping rate
  - Database and trends
    - Parameter ranges in which measurements are possible
    - Parametric dependencies for dominant damping mechanism(s)
  - Individual discharge analysis
    - Essential for detailed comparisons with theory
    - Ex.: modeling of  $n=3$  TAEs using LEMan, CASTOR, TAEFL
- Outlook: Active diagnostic upgrade

# Real-time detection and tracking of individual n-components

real time AELM data for JET shot #77417



- Antenna spectrum contains several frequency-degenerate modes
- Discrimination of n's done in real-time using sparse representation method (SparSpec)
- Computation within 850 $\mu$ s

## Example

- Antenna configuration to drive odd modes ( $3 < |n| < 11$ )
- n=3 mode dominates
- Weaker n=5, 7 modes

Post-pulse analysis (no time limitations for calculation) reveals entire set of modes  
 For modes also measured in real-time the difference in  $f$ ,  $\gamma/\omega$  and  $n$  are within 20%

# AE damping database

## Range of plasma parameters

plasma parameter	min	max
$I_p$ (MA)	0.93	2.58
$B_0$ (T)	1.87	3.37
$q_0$	0.72	2.70
$q_{95}$	2.63	6.56
$s_{95}$	2.21	4.73
$\kappa_{95}$	1.25	1.71
$\delta_u$	-0.02	0.34
$\delta_l$	-0.01	0.33
ROG (m)	0.01	0.12
$\beta_N$	0.15	0.54
$\beta_t$ (%)	0.06	0.46
$\beta_p$	0.13	0.25
$n_{e0}$ ( $10^{19} \text{ m}^{-3}$ )	$\sim 1$	$\sim 3$
$T_{e0}$ (keV)	$\sim 1$	$\sim 4$

~10000 TAE damping rate measurements in ohmic plasmas

# AE damping database

## Looking for correlations

plasma quantity	n						
	1	2	3	4	5	6	7
$q_0$	-0.05	-0.00	-0.30	-0.17	-0.05	0.09	-0.42
$q_{95}$	-0.05	-0.03	-0.18	-0.06	0.11	0.04	<b>-0.57</b>
$q_{95} - q_0$	-0.05	-0.05	-0.09	0.04	0.20	0.02	<b>-0.52</b>
$q_{95}/q_0$	0.14	-0.04	0.24	0.20	0.16	-0.21	0.11
$s_{95}$	<b>0.60</b>	0.22	0.31	0.29	0.09	-0.28	0.02
$\kappa_{95}$	<b>0.49</b>	0.19	0.22	0.25	-0.13	-0.26	-0.18
$\delta_u$	0.36	0.12	0.10	0.21	0.07	-0.39	-0.05
$\delta_l$	0.30	0.22	0.04	0.29	0.20	-0.29	-0.13
$I_p$	-0.40	-0.04	-0.00	0.12	-0.24	-0.14	-0.24
$\beta_N$	0.22	0.04	0.32	0.13	0.01	0.04	0.38
$\beta_t$	0.28	0.07	0.38	0.17	-0.09	0.02	<b>0.50</b>
$\beta_p$	0.09	-0.07	0.04	0.09	0.21	-0.03	-0.09
$B_0$	<b>-0.69</b>	-0.09	-0.34	-0.07	-0.24	-0.06	<b>-0.51</b>
ROG	-0.01	-0.07	0.03	0.21	-0.05	0.05	0.12
$T_{e0}$	-0.20	-0.01	-0.13	0.09	-0.38	0.01	-0.10
$n_{e0}$	-0.16	-0.09	-0.02	-0.07	-0.23	-0.27	-0.29
$\lambda$	<b>0.75</b>	0.18	0.23	0.29	-0.05	0.05	0.05

Correlation coefficient

$$r = \frac{\sum_{i=1}^N (x_i - \bar{x})(y_i - \bar{y})}{\sqrt{\sum_{i=1}^N (x_i - \bar{x})^2 \sum_{i=1}^N (y_i - \bar{y})^2}}$$



# AE damping database

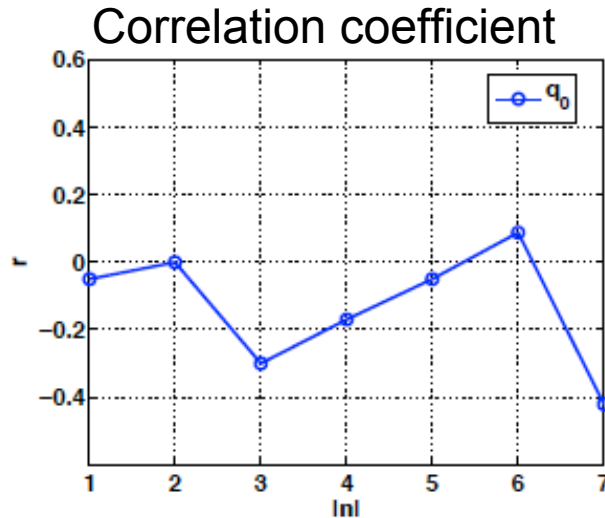
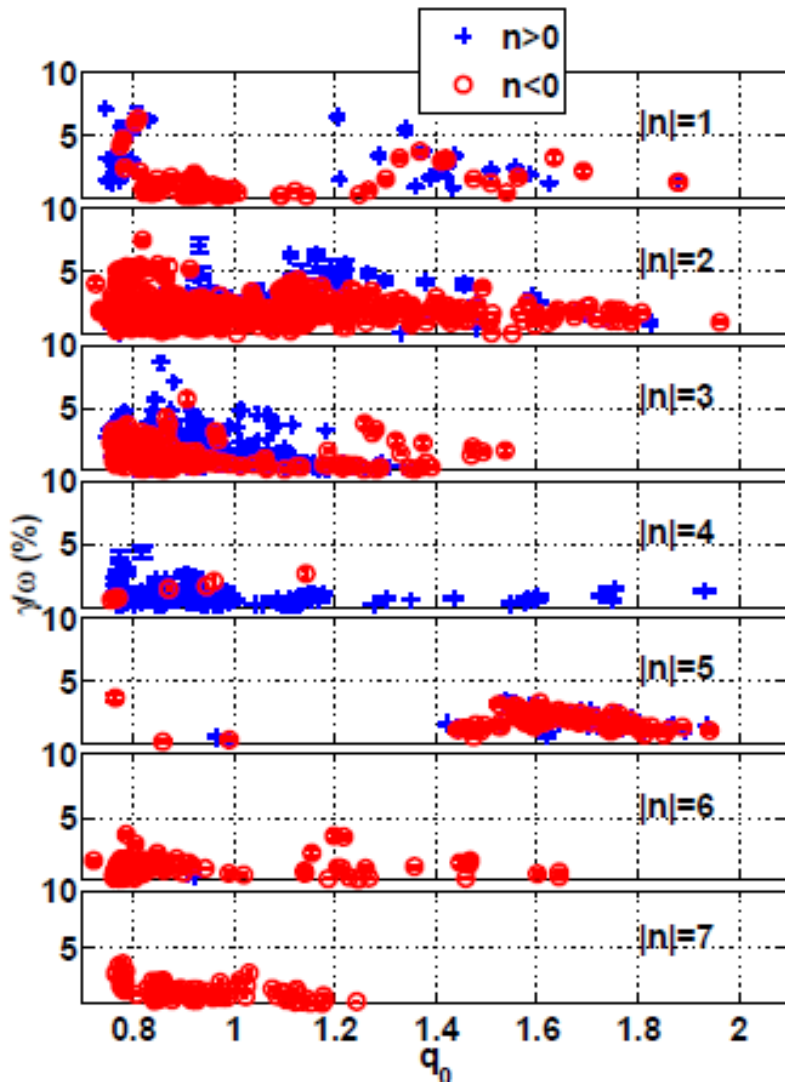
## Looking for correlations

plasma quantity	n						
	1	2	3	4	5	6	7
$q_0$	-0.05	-0.00	-0.30	-0.17	-0.05	0.09	-0.42
$q_{95}$	-0.05	-0.03	-0.18	-0.06	0.11	0.04	-0.57
$q_{95} - q_0$	-0.05	-0.05	-0.09	0.04	0.20	0.02	-0.52
$q_{95}/q_0$	0.14	-0.04	0.24	0.20	0.16	-0.21	0.11
$s_{95}$	<b>0.60</b>	0.22	0.31	0.29	0.09	-0.28	0.02
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$T_{e0}$	-0.20	-0.01	-0.13	0.09	-0.38	0.01	-0.10
$n_{e0}$	-0.16	-0.09	-0.02	-0.07	-0.23	-0.27	-0.29
$\lambda$	<b>0.75</b>	0.18	0.23	0.29	-0.05	0.05	0.05

Correlation coefficient

$$r = \frac{\sum_{i=1}^N (x_i - \bar{x})(y_i - \bar{y})}{\sqrt{\sum_{i=1}^N (x_i - \bar{x})^2 \sum_{i=1}^N (y_i - \bar{y})^2}}$$

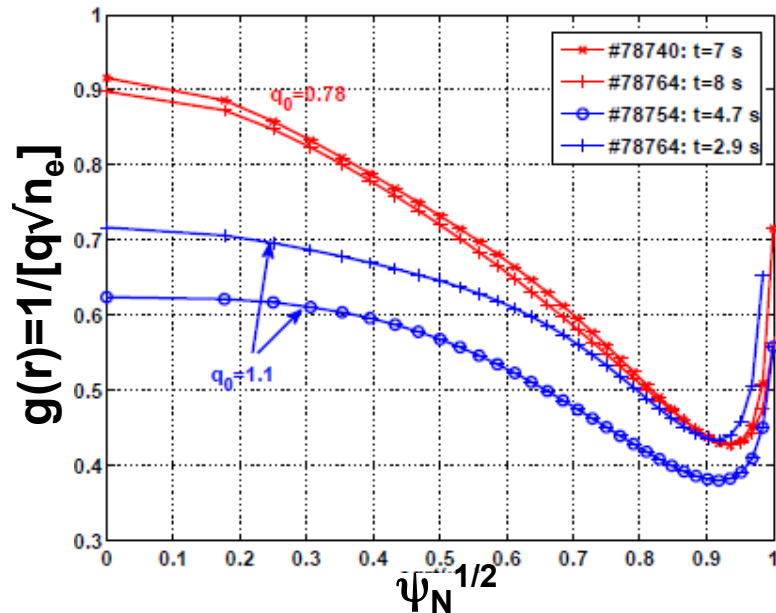
# $\gamma/\omega$ for $1 < |n| < 7$ TAEs vs. $q_0$



- No clear general trend in  $\gamma/\omega=f(q_0)$
- *Try selection of shots with similar  $q_{95}$  and varying  $q_0$  during  $q$ -profile relaxation*

# $\gamma/\omega$ for $1 < |n| < 7$ TAEs vs. $q_0$

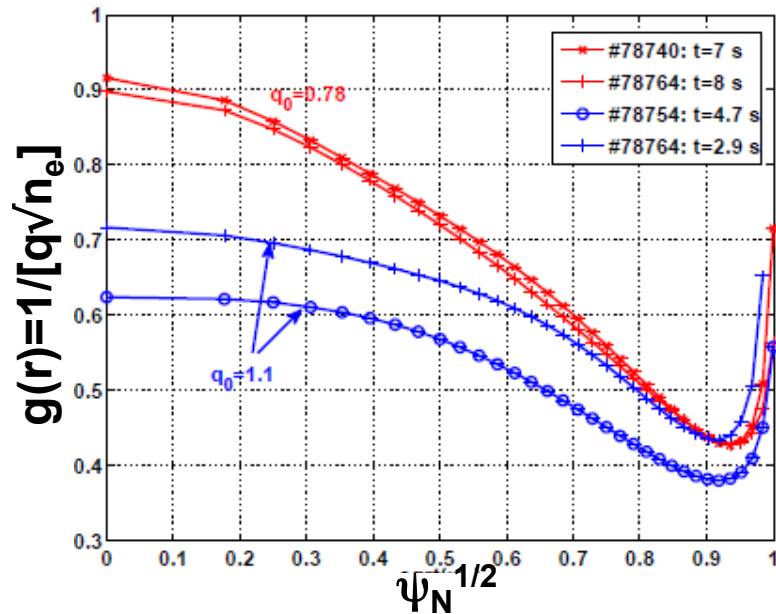
## Selection of shots with fixed $q_{95}$



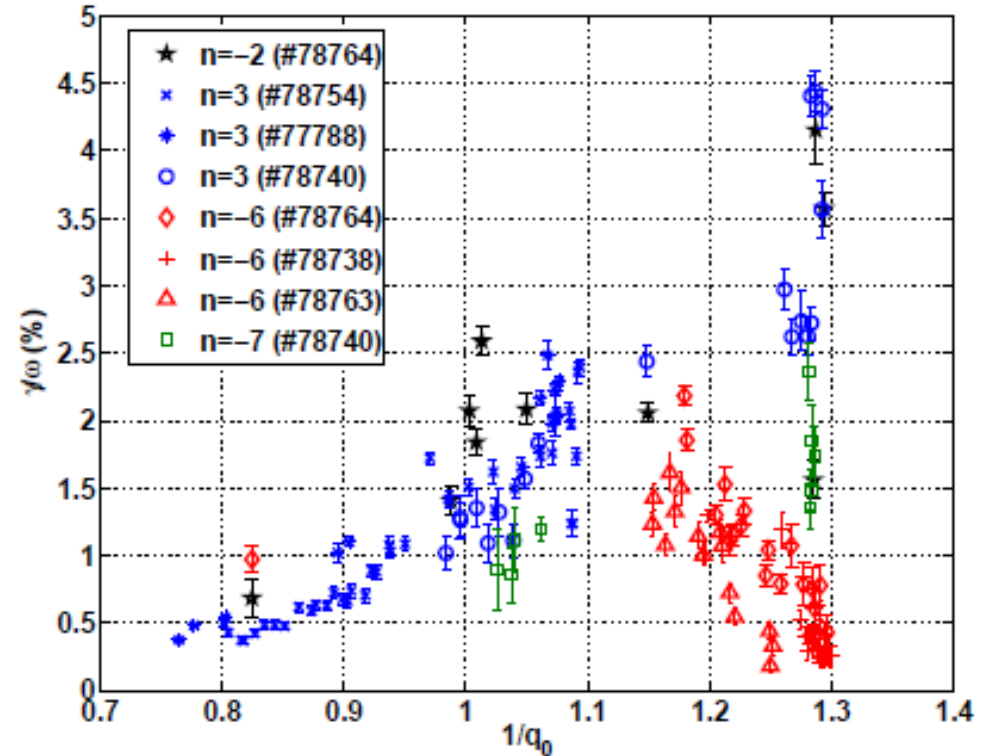
- For same  $q_{95}$ , Alfvén continuum gaps get less and less aligned as q-profile relaxes ( $q_0$  decreases)
- This effect is quantified by  $1/[q\sqrt{n_e}(r)]$
- Estimate of continuum damping

# $\gamma/\omega$ for $1 < |n| < 7$ TAEs vs. $q_0$

## Selection of shots with fixed $q_{95}$



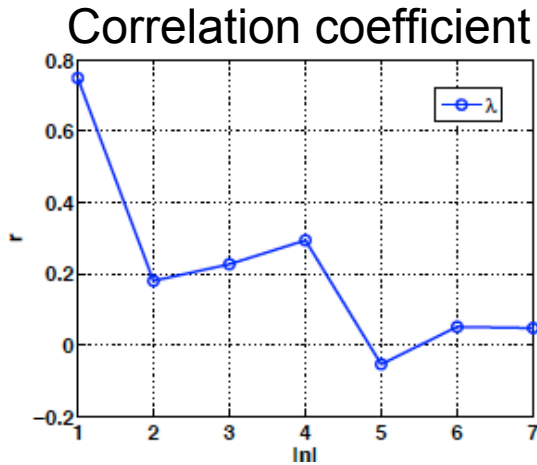
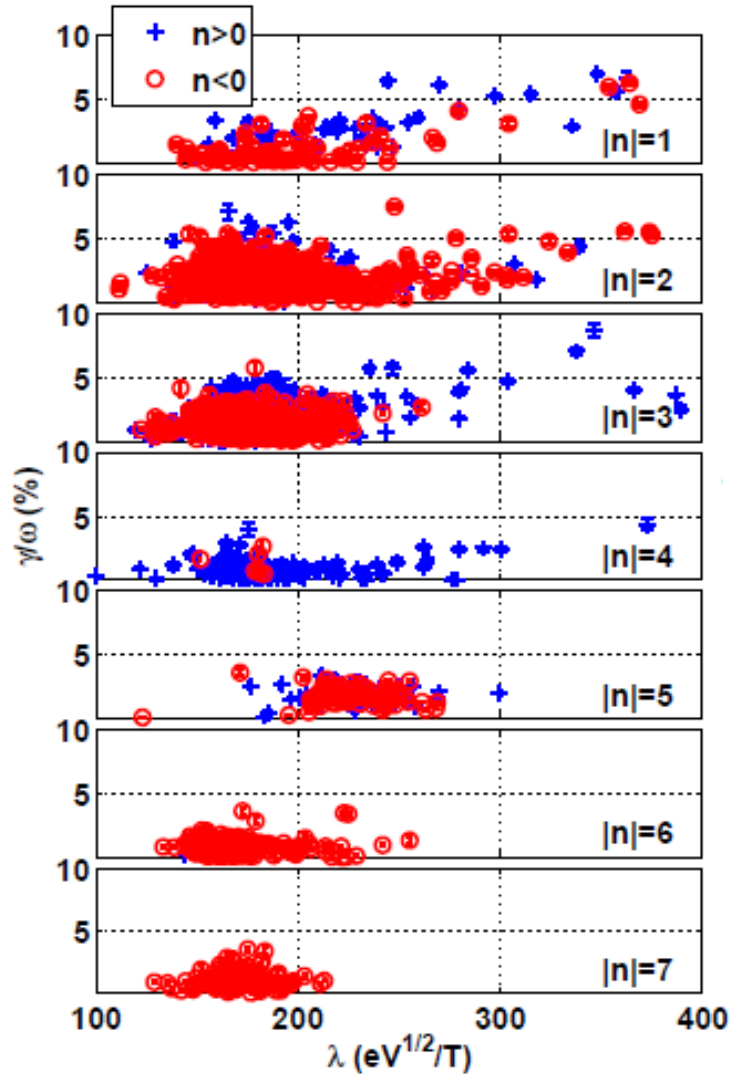
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- This effect is quantified by  $1/[q\sqrt{n_e}(r)]$
- Estimate of continuum damping



- Evidence for difference between n-ranges: continuum and radiative damping less effective as n increases



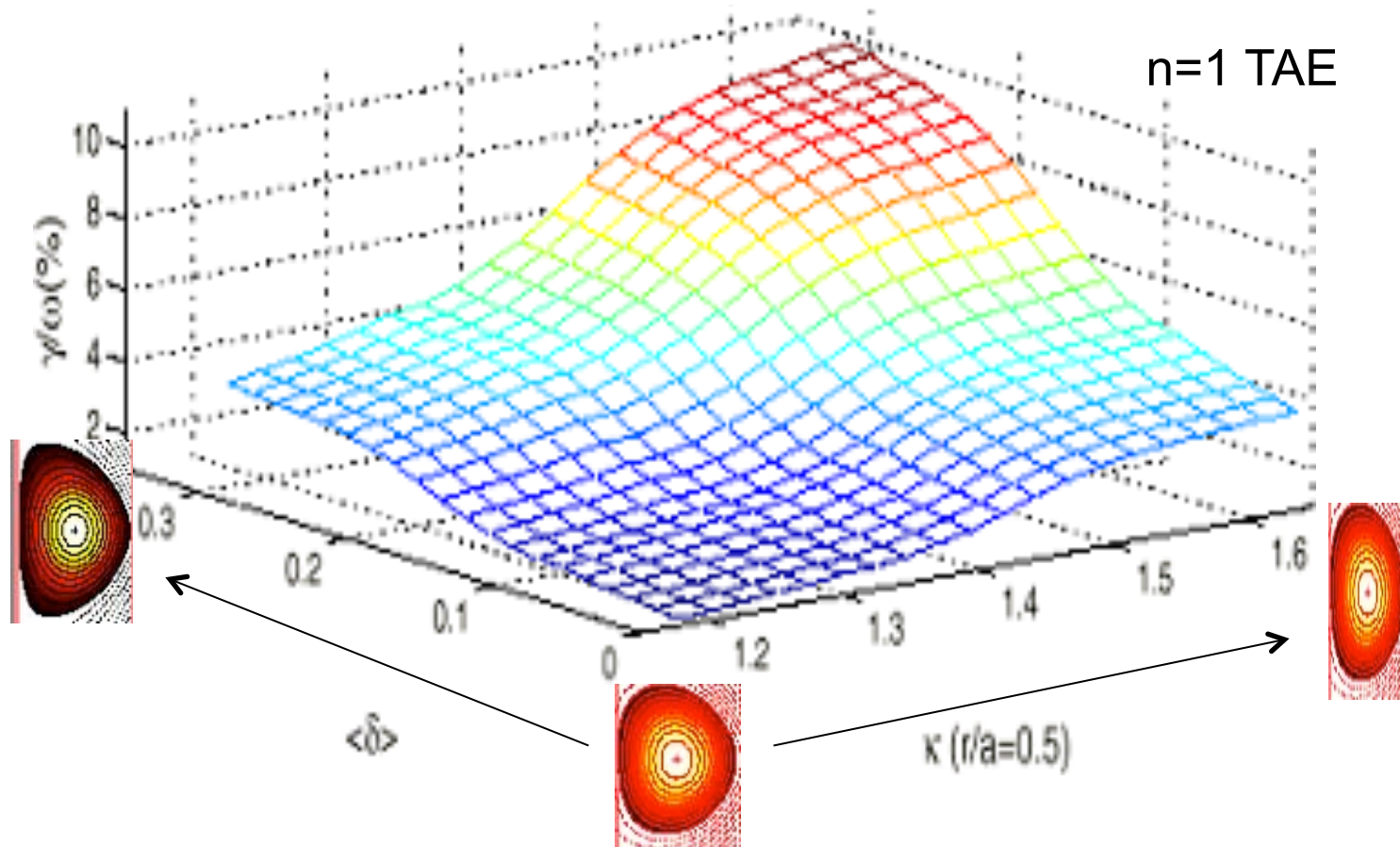
# $\gamma/\omega$ for $1 < |n| < 7$ TAEs vs. $\lambda$



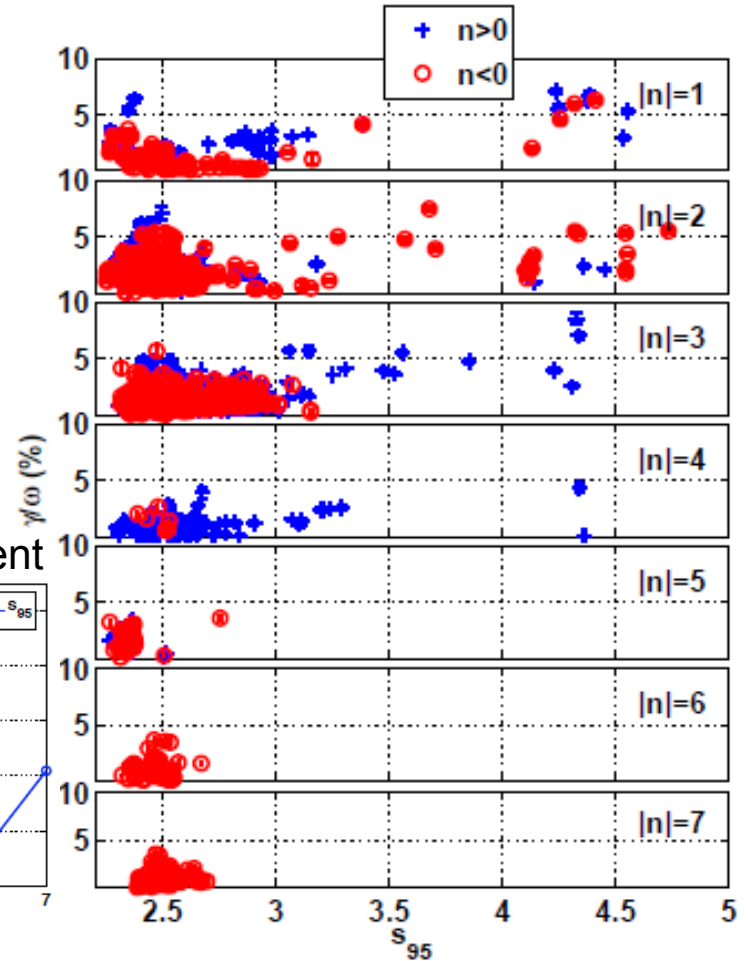
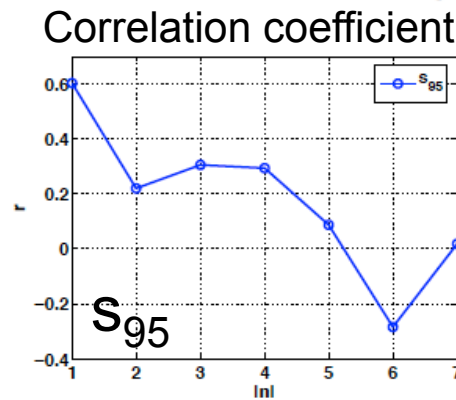
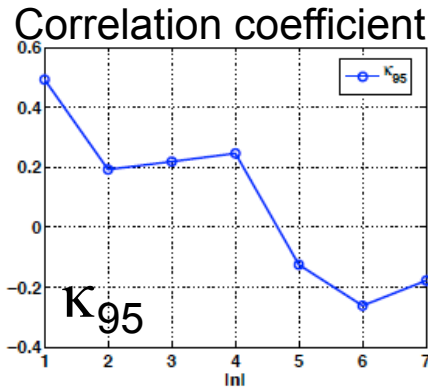
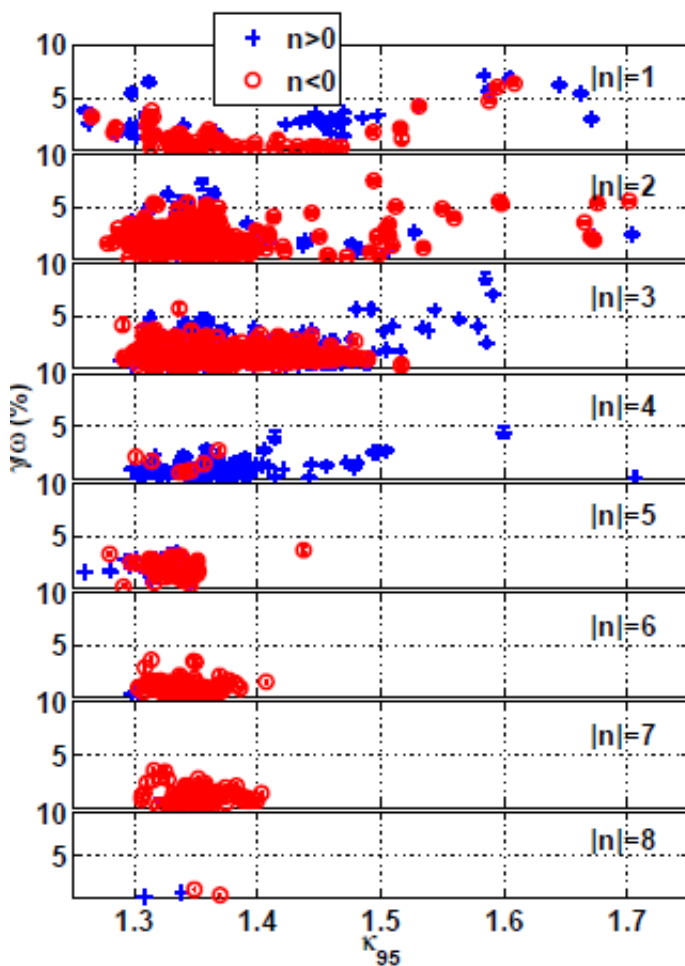
- Non-ideal parameter  $\lambda \propto q_{95} s_{95} \sqrt{T_{e0}} / B_0$  provides estimate of radiative damping
- Clear trend for  $|n|=1$ :  $\gamma/\omega_{|n|=1} \propto \lambda$
- Similar trend for  $2 \leq |n| \leq 4$ , only for large  $\lambda$
- No clear trend for  $|n| \geq 5$

# Edge shape effect on TAE damping

Old, saddle coil driven  $n=1$  measurements indicated a clear trend

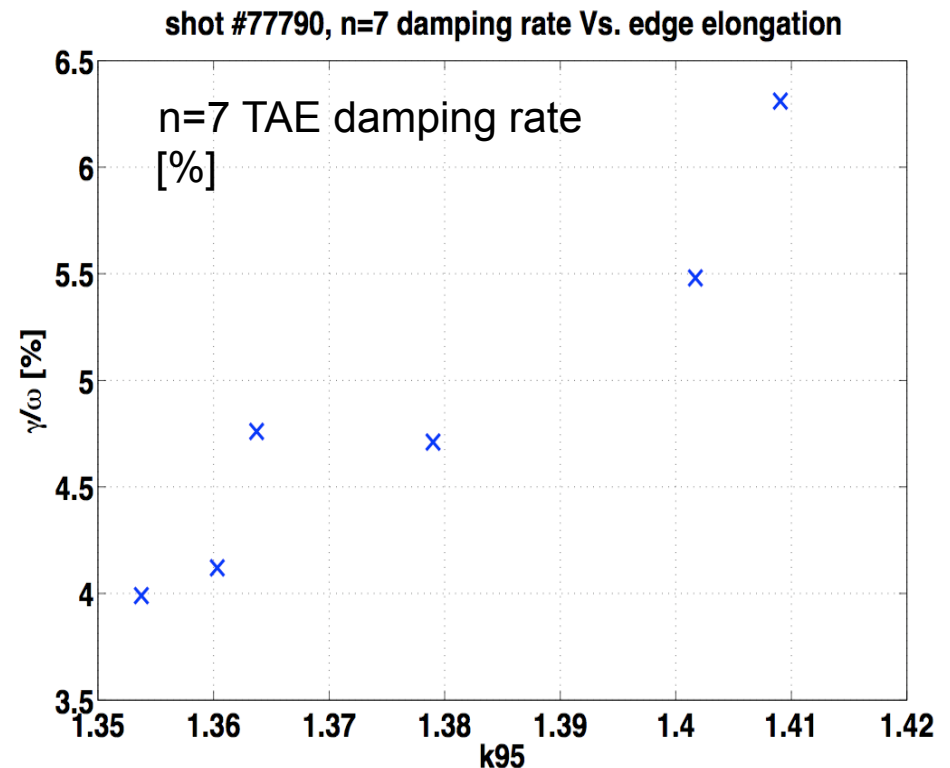
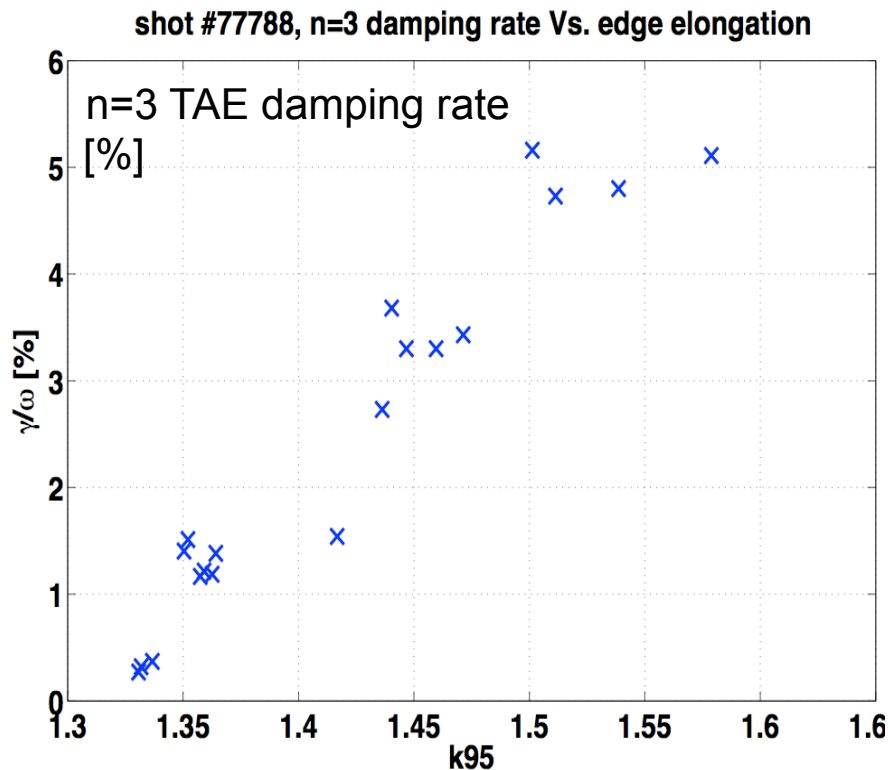


# $\gamma/\omega$ for $1 < |n| < 7$ TAEs vs. edge shape



- No data points for which  $\gamma/\omega < 7\%$  for large values of  $\kappa_{95}$  and  $s_{95}$
- Clear trend only visible in general for low- $n$
- To assess effect for medium- $n$  a single shot approach is necessary

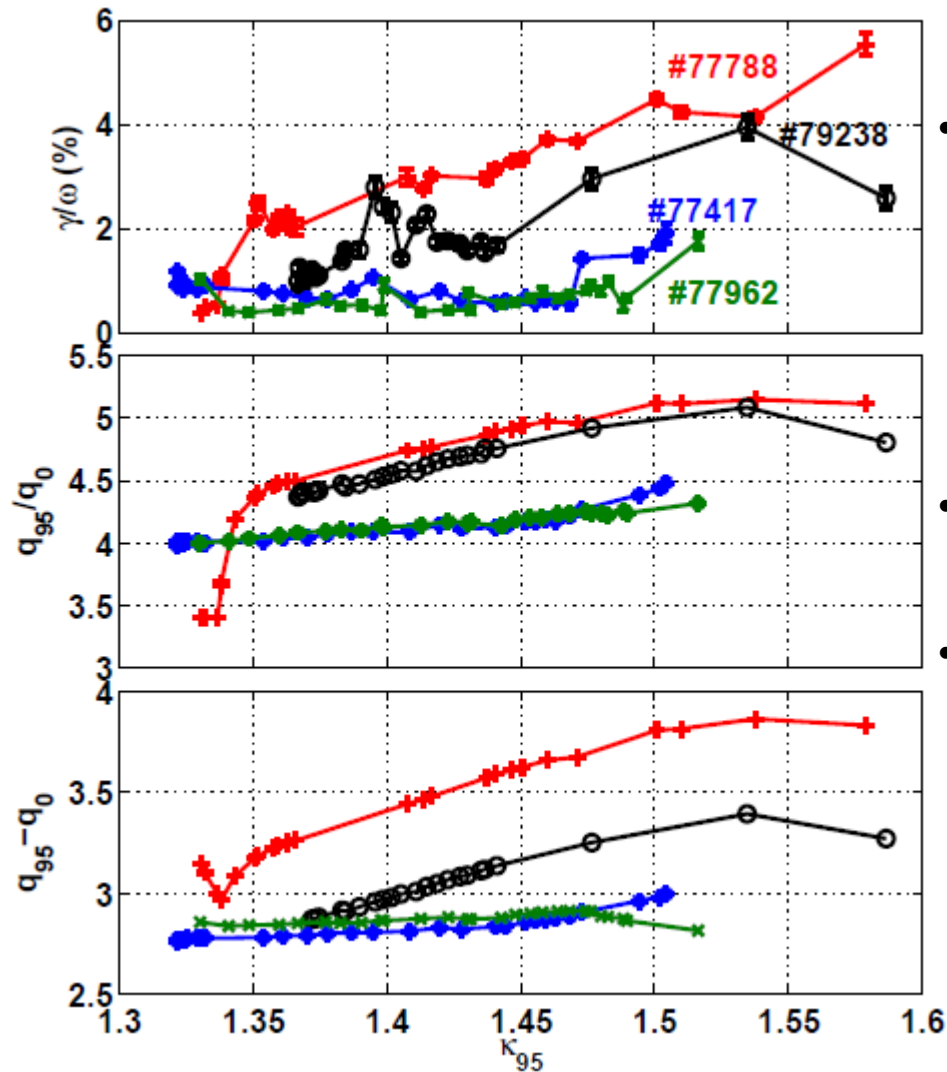
# $\gamma/\omega$ for $|n|=3, 7$ TAEs vs. edge shape Single shot approach



- In single shot, single-n measurements, edge elongation leads to increase in damping
- Why is this trend not visible in database?
- *Disentangle two dependences, e.g. q-profile and edge shape in specific experiment*

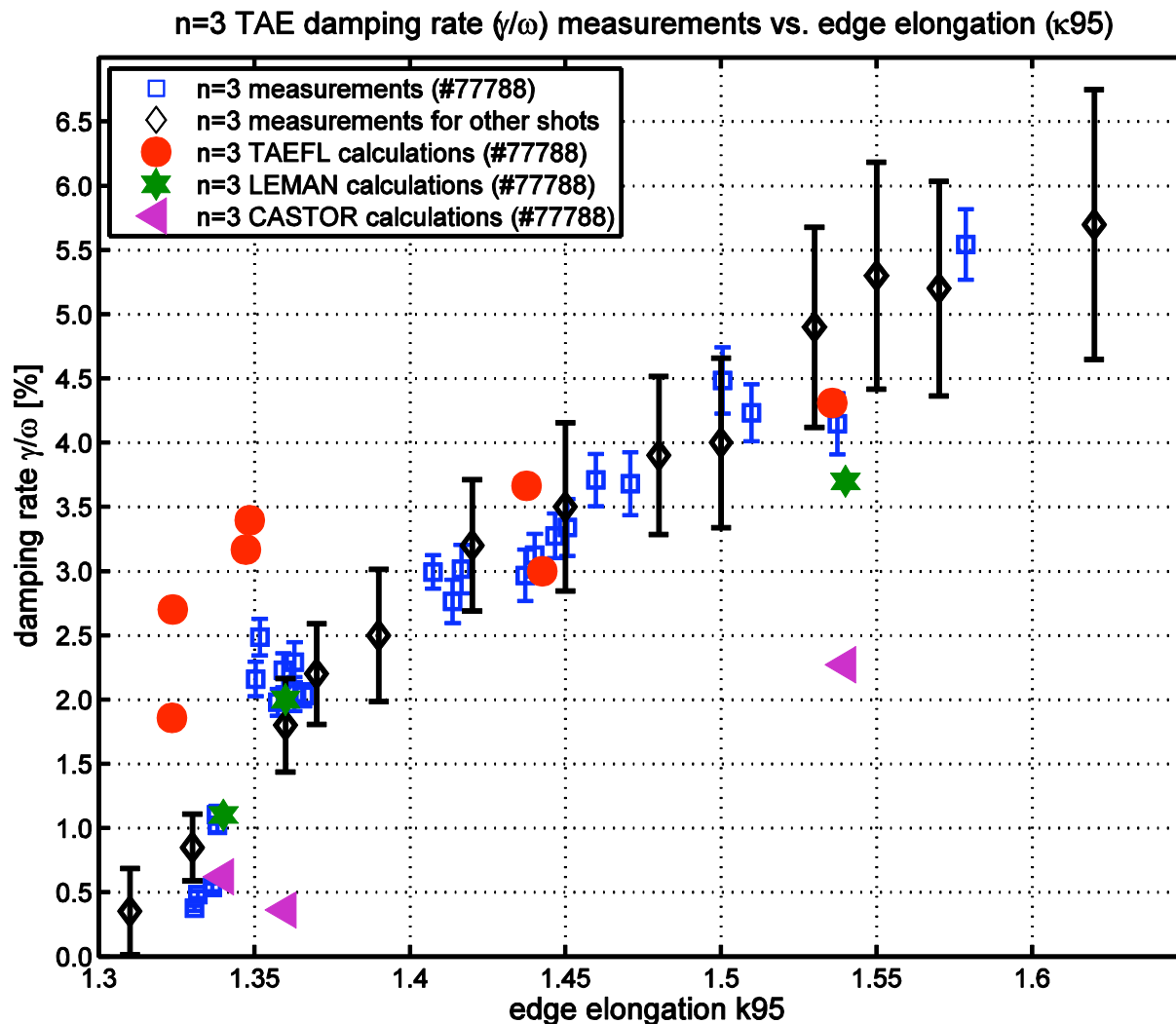


# Disentangling influence of edge shape and q-profile on n=3 TAE damping



- Differences in  $\gamma/\omega$  for same  $\kappa_{95}$  in different discharges seem related to 'span' of q-profile, quantified by  $q_{95}/q_0$  and  $q_{95}-q_0$  (proportional to number of poloidal harmonics)
- Damping increases with  $q_{95}/q_0$  and  $q_{95}-q_0$
- Trend consistent with electron Landau damping scaling  $(\gamma/\omega)_{\text{ELD}} \propto n^2(q_{95}-q_0)^2$

# n=3 TAE damping: modeling shape effect



- **CASTOR**: fluid
    - Large discrepancies with data
  - **TAEFL**: gyro-fluid
    - $\gamma/\omega$  extrapolated back from marginal stability threshold for fast-ion driven modes
    - Good agreement with data, important to test validity of modeling approach
  - **LEMan**: gyro-kinetic
    - Good agreement with data
    - Electron Landau damping of mode converted kinetic AW
    - *Note: good agreement was also found in comparisons with gyro-kinetic code LIGKA*
- IAEA-TCM-EP, Austin, 07-10 September 2011

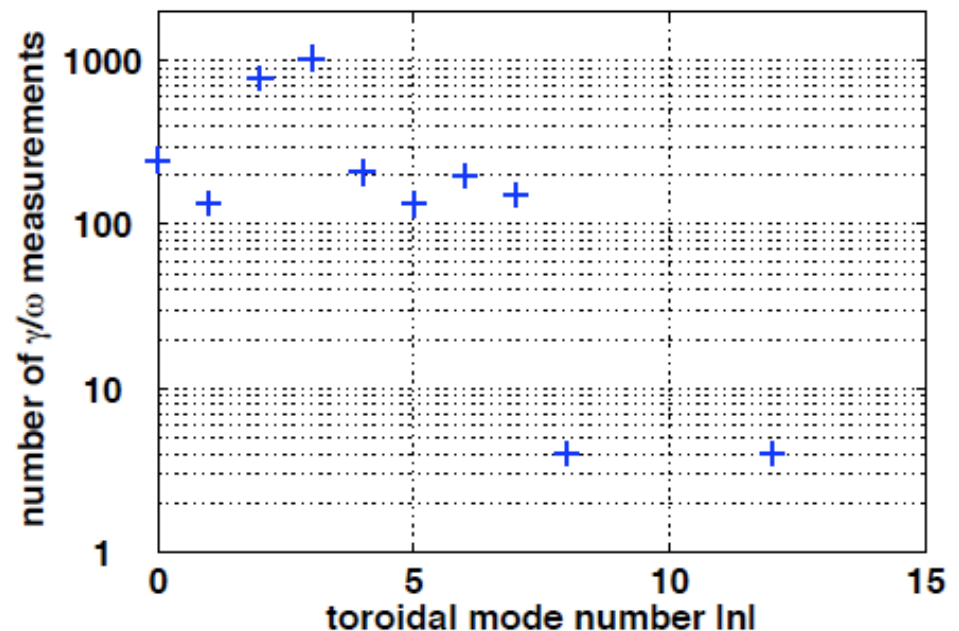
# Summary

- JET AE exciter drives and detects real-time selected spectrum of medium-n Aes
- Damping measurements database suggests some global trends for damping rates
- Individual discharge and n analysis is needed to disentangle complex dependences of  $\gamma/\omega$  on various parameters and profiles
- Agreement with numerical codes based on gyro-kinetic or gyro-fluid models has reached quantitative level

# Limitations of present AE system

- Coupling
  - Even with optimal matching and coils from both antennas, with one amplifier ( $I_{\max} \sim 5A$ ), core amplitudes are very small
    - $\sim 0.05mG/1A$  in the plasma core for modes with  $n \sim 7-15$
  - Tracking is difficult in the presence of noise (e.g. during strong additional heating) and if LCFS is distant from antennas
  - Mode identification requires sophisticated n-detection algorithms
- Mode selection
  - Wide spectrum; plasma preferentially selects low-n's

n distribution of  
damping rate  
measurements





# Upgrade to 8 independent amplifiers

- Maximize antenna currents within feed-through limits (25A)
- Better definition of antenna spectrum
  - Gain in single mode excitation by ~ factor of 4
  - More balanced distribution of currents without transformer coupling
- Arbitrary phasing
  - Definition of sign of  $n$  (traveling wave): identification of fast ion contribution
  - Simultaneous excitation of selected modes (different  $n$ 's or frequencies)
    - Stronger constraint on theory simulations for ITER extrapolation
    - Multi-point diagnostic applications
      - $q(r_1, r_2, \dots), A_{\text{EFF}}(r_1, r_2, \dots), \omega_{\text{TOR}}(r_1, r_2, \dots)$
- Test of burn control ideas
  - Real time control of fast ion stability by mode tracking



# the Sparse Representation Method

SparSpec minimizes the L1-norm penalized criterion:

$$J(x) = \frac{1}{2} \|y - Wx\|^2 + \lambda \sum_{k=-K}^K |x_k|_{L1}$$

$y$ : vector of data taken at time  $t_k$  [ $\equiv$  position  $\phi_k$ ]

$W$ : spectral window  $\exp(i2\pi t_k f_n)$  [ $\equiv \exp(i2\pi \phi_k n)$ ]

$x$ : vector of (I,Q) signals for frequencies  $f_n$

$\lambda$ : parameter fixed to obtain a satisfactory sparse solution  $\rightarrow$  **penalty criterion for invoking more modes to find adequate solution**

$\lambda$  can be fixed a-priori from known noise variance

- the sparse signal representation method is ideally suited for mode number analysis in fusion plasmas:
  - specifically designed for un-evenly distribution of sensors
  - allowable mode numbers are discretized:  $|n| = 0, \pm 1, \pm 2, \pm 3 \dots$
  - large (n,m)-range, number of modes not assumed a priori
  - amplitude and phase equally important for fitting algorithm
  - no need for a-posteriori thresholding to discriminate between solutions as  $\lambda$ -penalty determined a-priori from knowledge of noise variance
  - very efficient, very fast convergence  $\rightarrow$  ideal for real-time applications
  - **now implemented and fully validated in JET real-time and post-pulse mode tracking algorithm for stable Alfvén Eigenmodes**
  - **accuracy  $\rightarrow$  need correct interpretation of the spectral window**

# $\gamma/\omega$ measurements for medium-n AEs: data available for theory comparisons

- database compiled of  $\gamma/\omega(n)$  as function of plasma parameters and configurations for individual mode numbers
  - in excess of 10'000 individual  $\gamma/\omega(n)$  measurements already analyzed
  - in excess of 60 individual discharges already analyzed
  - various dedicated scans in plasma parameters have been run:
    - elongation scan during ohmic phase,  $1.25 < \kappa_{95} < 1.65$  without ICRF
      - add ICRF with PRF=2MW and PRF=3MW, different phasing (dipole and +/-90)
      - add PRF modulations 2MW +1MW/300ms, different phasing (dipole and +/-90)
    - ohmic Bfield/ne scan, change RF deposition profile and edge continuum
      - add PRF with power ramp-up to 4.5MW, different phasing (dipole and +/-90)
- damping rate as function of plasma isotope composition and ion Larmor radius
- damping rate for medium-n ( $n=3-7$ ) TAEs at ICRF power switch off with constant plasma parameters
  - direct measurement of MeV-ions drive to the modes?
- effect of ripple in the magnetic field medium-n ( $n=3-7$ ) TAEs with/out fast ions:
  - fast ion losses (resonant NBI ions with  $V_{||} \sim V_A/3$ ), affecting drive for the modes?
  - change density scale length at plasma edge, affecting the continuum damping?